Lecture 18

20.110/5.60/2.772

Clausius Clapeyron equation

Condensed phase equilibria

CLAUSIUS-CLAPERYON EQUATION

We derived the Clapeyron Equation last lecture:

$$\left(\frac{dp}{dT}\right)_{coexist} = \left[\frac{\overline{S}_{\beta} - \overline{S}_{\alpha}}{\overline{V}_{\beta} - \overline{V}_{\alpha}}\right] = \left(\frac{\Delta \overline{S}}{\Delta \overline{V}}\right)_{\alpha \to \beta}$$

This is exact!

Let's examine this for equilibrium between solid-liquid and solid-solid phases. For example, fusion:

$$\left(\frac{dp}{dT}\right)_{coexist} = \frac{\Delta S_{\rm fus}}{\Delta V_{\rm fus}} \qquad \text{use} \qquad \Delta H_{\rm fus} = T\Delta S_{\rm fus}$$

We can make some assumptions about these types of phase equilibria.

$$\int_{p_1}^{p_2} dp = \int_{T_m}^{T_m'} \frac{\Delta H_{\text{fus}}}{\Delta V_{\text{fus}}} \cdot \frac{dT}{T} \qquad \text{assume } \Delta H_{\text{fus}}, \Delta V_{\text{fus}} \text{ are independent of p,T}$$

$$p_2 - p_1 = \frac{\Delta H_{fus}}{\Delta V_{fus}} \ln \left(\frac{T_m'}{T_m} \right) \qquad T_m = \text{melting temp } @p_2, T_m = \text{melting temp } @p_1$$

 $T_m - T_m \sim 0$ (expect small change in T_m)

$$\ln\left(\frac{T_{m}}{T_{m}}\right) = \ln\left(\frac{T_{m} + T_{m}' - T_{m}}{T_{m}}\right) = \ln\left(1 + \frac{T_{m}' - T_{m}}{T_{m}}\right)$$

Since fraction is small, then

$$\ln\left(\frac{T_{m}'}{T_{m}}\right) \approx \left(\frac{T_{m}' - T_{m}}{T_{m}}\right) = \frac{\Delta T}{T}$$

Then

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 $\Delta p \cong \frac{\Delta H_{\text{fus}}}{\Delta V_{\text{fus}}} \cdot \frac{\Delta T_m}{T_m}$

gives you ΔT_m , the melting point increase corresponding to a Δp increase

also says: the coexistence line is approximately linear

• Equilibrium between gas/liquid or gas /solid (gas/condensed phase) substance A:

A(I)= A(g)orA(s) = A(g)
$$\Delta H_{vap}$$
 ΔH_{sub}

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$$\left(\frac{dp}{dT}\right) = \frac{\Delta \overline{S}}{\Delta \overline{V}} = \frac{\Delta \overline{H}}{T\left(\overline{V}_{g} - \overline{V}_{c}\right)} \qquad \qquad V_{c} = V_{\text{solid}} \text{ or } V_{\text{liquid}}$$

- a) We can approximate
- $\overline{V}_g \overline{V}_c \approx \overline{V}_g$ because V_g>>V_c

b) We can assume the gas is ideal

$$\overline{V}_g = \frac{RT}{p}$$

Clausius Clapeyron Equation

$$\left(\frac{dp}{dT}\right) \approx \frac{\Delta \overline{H}}{T(\overline{V}_{g})} = \frac{\Delta \overline{H} \cdot p}{RT^{2}}$$
$$\left(\frac{d\ln p}{dT}\right) \approx \frac{\Delta \overline{H}}{RT^{2}}$$

to get here we made approximations, but very good far from Tc.

Relates the vapor pressure of a liquid to ΔH_{vap} or vapor pressure of a solid to ΔH_{sub} .

c) We can make yet another approximation: assume ΔH independent of T. Then can integrate:

$$\int_{p_0}^p d\ln p = \int_{T_0}^T \frac{\Delta H}{RT^2} dT$$

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$$\ln\left(\frac{p}{p_0}\right) \approx -\frac{\Delta \overline{H}}{R} \left(\frac{1}{T} - \frac{1}{T_0}\right)$$

This equation allows us to get ΔH_{sub} from p

