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5.80 Small-Molecule Spectroscopy and Dynamics
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MASSACHUSETTS INSTITUTE OF TECHNOLOGY
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Problem Set #2 ANSWERS

1. (a) Construct the state $|L = 2, S = 1, J = 1, M_J = 0\rangle$ from the $|LM_L S M_S\rangle$ basis using the ladder operator plus orthogonality technique.

ANSWER:

$$|L = 2, S = 1, J = 1, M_J = 0\rangle$$

We know that $|L = 2, S = 1, J = 3, M_J = 3\rangle = |L = 2, S = 1, M_L = 2, M_S = 1\rangle$

$$\Rightarrow \mathbf{J}_- |2133\rangle = (\mathbf{L}_- + \mathbf{S}_-) |2121\rangle$$

$$\sqrt{6} |2132\rangle = 2 |2111\rangle + \sqrt{2} |2120\rangle$$

$$|2132\rangle = \frac{2}{\sqrt{6}} |2111\rangle + \frac{1}{\sqrt{3}} |2120\rangle$$

$$(L S J M_S) (L S M_L M_S)$$

$$|2122\rangle = a |2111\rangle + b |2120\rangle \text{ where } a^2 + b^2 = 1$$

$$\text{and } \langle 2132 | 2122 \rangle = 0 = \frac{2a}{\sqrt{6}} \langle 2111 | 2111 \rangle + \frac{b}{\sqrt{3}} \langle 2120 | 2120 \rangle = \frac{1}{\sqrt{16}} (2a + b\sqrt{2})$$

$$\Rightarrow \langle 2122 | = \frac{-1}{\sqrt{3}} \langle 2111 | + \sqrt{\frac{2}{3}} \langle 2120 |$$

$$\mathbf{J}_- |2122\rangle = (\mathbf{L}_- + \mathbf{S}_-) \left[\frac{-1}{\sqrt{3}} |2121\rangle + \sqrt{\frac{2}{3}} |2120\rangle \right]$$

$$\Rightarrow |2121\rangle = \frac{1}{\sqrt{3}} |212-1\rangle + \frac{1}{\sqrt{6}} |2110\rangle - \frac{1}{\sqrt{2}} |2101\rangle$$

$$\text{Then } \mathbf{J}_- |L = 2, S = 1, J = 3, M_J = 2\rangle = (\mathbf{L}_- + \mathbf{S}_-) \left[\frac{2}{\sqrt{6}} |2111\rangle + \frac{1}{\sqrt{3}} |2320\rangle \right]$$

$$\Rightarrow |2131\rangle = \frac{2}{\sqrt{10}} |2101\rangle + \frac{4}{\sqrt{30}} |2110\rangle + \frac{1}{\sqrt{15}} |212-1\rangle$$

$$\text{We know that } |L = 2, S = 1, J = 1, M_S = 1\rangle = a |2101\rangle + b |2110\rangle + c |212-1\rangle$$

$$\langle L = 2, S = 1, J = 1, M_J = 1 | L = 2, S = 1, J = 2, M_J = 1 \rangle = 0$$

$$\text{and } \langle L = 2, S = 1, J = 1, M_J = 1 | L = 2, S = 1, J = 3, M_J = 1 \rangle = 0$$

$$\left. \begin{aligned} \Rightarrow \frac{2a}{\sqrt{10}} + \frac{4b}{\sqrt{30}} + \frac{c}{\sqrt{15}} &= 0 \\ a^2 + b^2 + c^2 &= 1 \\ -\frac{a}{\sqrt{2}} + \frac{b}{\sqrt{6}} + \frac{c}{\sqrt{3}} &= 0 \end{aligned} \right\} \text{ so } \begin{aligned} a &= \frac{1}{\sqrt{10}} \\ b &= \sqrt{\frac{3}{10}} \\ c &= -\sqrt{\frac{3}{5}} \end{aligned}$$

$$\text{so } \langle L = 2, S = 1, J = 1, M_J = 1 | = -\frac{1}{\sqrt{10}} \langle 2101 | + \sqrt{\frac{3}{10}} \langle 2110 | - \sqrt{\frac{3}{5}} \langle 212-1 |$$

$$\text{then operating on both sides with } \mathbf{J}_- = \mathbf{S}_- + \mathbf{L}_- \text{ we get } |L = 2, S = 1, J = 1, M_J = 0\rangle = \sqrt{\frac{3}{10}} |21-11\rangle - \sqrt{\frac{4}{10}} |2100\rangle + \sqrt{\frac{3}{10}} |211-1\rangle$$

(b) Construct the states

$$|L = 2, S = 1, J = 1, M_J = 0\rangle \quad {}^3D_1$$

$$|L = 2, S = 2, J = 1, M_J = 0\rangle \quad {}^5D_1$$

$$|L = 5, S = 2, J = 3, M_J = 1\rangle \quad {}^5H_3$$

from the $|L M_L S M_S\rangle$ basis using Clebsch-Gordan coefficients. The 3D_1 function is the same as in Part (a) and is intended as a consistency check.

ANSWER:

$$|L = 2, S = 1, J = 1, M_J = 0\rangle$$

Since $S \equiv j_2 = 1$ look in table 2³ of Condon and Shortley,

also $J \equiv j = j_1 - 1$ look in column 3

$$\text{then } |L = 2, S = 1, J = 1, M_J = 0\rangle = \sqrt{\frac{(2-0)(2-0+1)}{2 \cdot 2(2 \cdot 2+1)}} |L = 2, S = 1, M_L = -1, M_S = 1\rangle$$

$$+ \sqrt{\frac{(2-0)(2+0)}{2(2 \cdot 2+1)}} |L = 2, S = 1, M_L = 0, M_S = 0\rangle + \sqrt{\frac{(2+0+1)(2+0)}{2 \cdot 2(2 \cdot 2+1)}} |L = 2, S = 1, M_L = 1, M_S = -1\rangle$$

$$|L = 2, S = 1, J = 1, M_J = 0\rangle = \sqrt{\frac{3}{10}} |L = 2, S = 1, M_L = -1, M_S = 1\rangle$$

$$+ \sqrt{\frac{2}{10}} |L = 2, S = 1, M_L = 0, M_S = 0\rangle + \sqrt{\frac{3}{10}} |L = 2, S = 1, M_L = 1, M_S = -1\rangle$$

$$|L = 2, S = 2, J = 1, M_J = 0\rangle = \frac{-2}{\sqrt{10}} |L = 2, S = 2, M_L = -2, M_S = 2\rangle$$

$$+ \frac{1}{\sqrt{10}} |L = 2, S = 2, M_L = -1, M_S = 1\rangle + 0 |L = 2, S = 2, M_L = 0, M_S = 0\rangle$$

$$- \frac{1}{\sqrt{10}} |L = 2, S = 2, M_L = 1, M_S = -1\rangle + \frac{2}{\sqrt{10}} |L = 2, S = 2, M_L = 2, M_S = -2\rangle$$

$$|L = 5, S = 2, J = 3, M_S = 1\rangle = |L = 5, S = 2, M_L = -1, M_S = 2\rangle$$

$$- |L = 5, S = 2, M_L = 0, M_S = 1\rangle - |L = 5, S = 2, M_L = 1, M_S = 0\rangle$$

$$- |L = 5, S = 2, M_L = 2, M_S = -1\rangle - |L = 5, S = 2, M_L = 3, M_S = -2\rangle$$

} these are not correct

2. We know that the spin-orbit Hamiltonian, $\mathbf{H}^{\text{SO}} = \mathbf{A}\mathbf{L} \cdot \mathbf{S}$, is diagonal in the $|L S J M_J\rangle$ basis but not in the $|L M_L S M_S\rangle$ basis.

(a) Construct the full nine by nine \mathbf{H}^{SO} matrix in the $|L = 1 M_L S = 1 M_S\rangle$ basis.

ANSWER:

The nine basis functions are:

	$ LS M_L M_S\rangle$
1⟩	$ 11 - 1 - 1\rangle$
2⟩	$ 1111\rangle$
3⟩	$ 1100\rangle$
4⟩	$ 111 - 1\rangle$
5⟩	$ 11 - 11\rangle$
6⟩	$ 110 - 1\rangle$
7⟩	$ 1101\rangle$
8⟩	$ 11110\rangle$
9⟩	$ 11 - 10\rangle$

The diagonal matrix elements: $\langle i|\mathbf{A}\mathbf{L} \cdot \mathbf{S}|i\rangle = \langle i|\mathbf{A}(\mathbf{L}_z\mathbf{S}_z + \frac{1}{2}[\mathbf{L}_+\mathbf{S}_- + \mathbf{L}_-\mathbf{S}_+])|i\rangle = A \langle i|\mathbf{L}_z\mathbf{S}_z|i\rangle$

So

$$H_{11} = H_{22} = A$$

$$H_{33} = H_{66} = H_{77} = H_{88} = H_{99} = 0$$

$$H_{44} = H_{55} = -A$$

The operators $\mathbf{L}_\pm\mathbf{S}_\mp$ connect states with $\Delta M_L = \mp 1$ and $\Delta M_S = \pm 1$.

So

$$H_{35} = H_{53} = A$$

$$H_{34} = H_{43} = A$$

$$H_{69} = H_{96} = A$$

$$H_{78} + H_{87} = A$$

All the rest are zero!

(b) Construct the

$$|L = 1, S = 1, J = 2, M_J = 0\rangle \quad {}^3P_2$$

$$|L = 1, S = 1, J = 1, M_J = 0\rangle \quad {}^3P_1$$

and $|L = 1, S = 1, J = 0, M_J = 0\rangle \quad {}^3P_0$

functions in the $|L M_L S M_S\rangle$ basis.

ANSWER:

$$|L = 1, S = 1, J = 2, M_J = 0\rangle = \frac{1}{\sqrt{6}} |L = 1, S = 1, M_L = -1, M_S = 1\rangle$$

$$+ \frac{2}{\sqrt{6}} |L = 1, S = 1, M_L = 0, M_S = 0\rangle + \frac{1}{\sqrt{6}} |L = 1, S = 1, M_L = 1, M_S = -1\rangle$$

$$|L = 1, S = 1, J = 1, M_J = 0\rangle = -\frac{1}{\sqrt{2}} |L = 1, S = 1, M_L = -1, M_S = 1\rangle$$

$$+ \frac{1}{\sqrt{2}} |L = 1, S = 1, M_L = 1, M_S = -1\rangle$$

$$|L = 1, S = 1, J = 0, M_J = 0\rangle = \frac{1}{\sqrt{3}} |L = 1, S = 1, M_L = -1, M_S = 1\rangle$$

$$- \frac{1}{\sqrt{3}} |L = 1, S = 1, M_L = 0, M_S = 0\rangle + \frac{1}{\sqrt{3}} |L = 1, S = 1, M_L = 1, M_S = -1\rangle$$

(c) Show that the matrix elements

$$\begin{aligned} & \langle L = 1, S = 1, J = 2, M_J = 0 | \mathbf{H}^{\text{SO}} | 1, 1, 2, 0 \rangle \\ & \langle 1, 1, 2, 0 | \mathbf{H}^{\text{SO}} | 1, 1, 1, 0 \rangle \\ & \langle 1, 1, 2, 0 | \mathbf{H}^{\text{SO}} | 1, 1, 0, 0 \rangle \\ & \langle 1, 1, 1, 0 | \mathbf{H}^{\text{SO}} | 1, 1, 1, 0 \rangle \\ & \langle 1, 1, 1, 0 | \mathbf{H}^{\text{SO}} | 1, 1, 0, 0 \rangle \\ & \langle 1, 1, 0, 0 | \mathbf{H}^{\text{SO}} | 1, 1, 0, 0 \rangle \end{aligned}$$

expressed in terms of the $|L M_L S M_S\rangle$ basis in part (b) have the values expected from $\mathbf{L} \cdot \mathbf{S} = 1/2(\mathbf{J}^2 - \mathbf{L}^2 - \mathbf{S}^2)$ evaluated in the $|L S J M_J\rangle$ basis.

ANSWER:

$$\begin{aligned} \langle 1120 | \mathbf{AL} \cdot \mathbf{S} | 1120 \rangle &= \frac{1}{2} A \langle 1120 | \mathbf{J}^2 - \mathbf{L}^2 - \mathbf{S}^2 | 1120 \rangle = \frac{1}{2} A [2(2+1) - 1(1+1) - 1(1+1)] = \frac{1}{2} A [2] = A \\ \langle 1120 | &= \frac{1}{\sqrt{6}} \langle 11 - 11 | + \frac{2}{\sqrt{6}} \langle 1100 | + \frac{1}{\sqrt{6}} \langle 111 - 1 | \\ \Rightarrow \langle 1120 | \mathbf{AL} \cdot \mathbf{S} | 1120 \rangle &= \frac{1}{6} \langle 11 - 11 | \mathbf{AL}_z \mathbf{S}_z | 11 - 11 \rangle + \frac{2}{3} \langle 1100 | \mathbf{AL}_z \mathbf{S}_z | 1100 \rangle + \frac{1}{6} \langle 111 - 1 | \mathbf{AL}_z \mathbf{S}_z | 111 - 1 \rangle + \\ &+ \frac{2}{6} \langle 11 - 11 | \frac{1}{2} \mathbf{AL}_+ \mathbf{S}_+ | 1100 \rangle + \frac{2}{6} \langle 1100 | \frac{1}{2} \mathbf{AL}_+ \mathbf{S}_+ | 11 - 11 \rangle + \frac{2}{6} \langle 111 - 1 | \frac{1}{2} \mathbf{AL}_+ \mathbf{S}_+ | 1100 \rangle + \\ &+ \frac{2}{6} \langle 1100 | \frac{1}{2} \mathbf{AL}_- \mathbf{S}_- | 111 - 1 \rangle = A \left(-\frac{1}{6} - \frac{1}{6} + \frac{2}{6} + \frac{2}{6} + \frac{2}{6} + \frac{2}{6} \right) = A \\ \langle 1120 | \mathbf{H}^{\text{SO}} | 1110 \rangle &= 0 \\ \langle 1120 | &= \frac{1}{\sqrt{6}} [\langle 11 - 11 | + 2 \langle 1100 | + \langle 111 - 1 |] \\ \langle 1110 | &= \frac{-1}{\sqrt{2}} [\langle 11 - 11 | - \langle 111 - 1 |] \\ \text{so } \langle 1120 | \mathbf{H}^{\text{SO}} | 1110 \rangle &= \frac{-1}{\sqrt{12}} \langle 11 - 11 | \mathbf{H}^{\text{SO}} | 11 - 11 \rangle + \frac{1}{\sqrt{12}} \langle 111 - 1 | \mathbf{H}^{\text{SO}} | 111 - 1 \rangle - \frac{2}{\sqrt{12}} \langle 1100 | \mathbf{H}^{\text{SO}} | 11 - 11 \rangle + \\ &+ \frac{2}{\sqrt{12}} \langle 1100 | \mathbf{H}^{\text{SO}} | 111 - 1 \rangle = 0 \\ \langle 1120 | \mathbf{H}^{\text{SO}} | 1100 \rangle &= 0 \\ \langle 1120 | &= \frac{1}{\sqrt{6}} [\langle 11 - 11 | + 2 \langle 1100 | + \langle 111 - 1 |] \\ \langle 1100 | &= \frac{1}{\sqrt{3}} [\langle 11 - 11 | - \langle 1100 | + \langle 111 - 1 |] \\ \text{so } \langle 1120 | \mathbf{H}^{\text{SO}} | 1100 \rangle &= \frac{1}{\sqrt{18}} \langle 11 - 11 | \mathbf{H}^{\text{SO}} | 11 - 11 \rangle + \frac{2}{\sqrt{18}} \langle 1100 | \mathbf{H}^{\text{SO}} | 1100 \rangle - \frac{1}{\sqrt{18}} \langle 111 - 1 | \mathbf{H}^{\text{SO}} | 111 - 1 \rangle + \\ &+ \frac{2}{\sqrt{18}} \langle 1100 | \mathbf{H}^{\text{SO}} | 11 - 11 \rangle + \frac{2}{\sqrt{18}} \langle 1100 | \mathbf{H}^{\text{SO}} | 111 - 1 \rangle - \frac{1}{\sqrt{18}} \langle 11 - 11 | \mathbf{H}^{\text{SO}} | 1100 \rangle - \frac{1}{\sqrt{18}} \langle 111 - 1 | \mathbf{H}^{\text{SO}} | 1100 \rangle = \\ &= \frac{A}{\sqrt{18}} (-1 - 1 + 2 + 2 - 1 - 1) = 0 \\ \langle 1110 | \mathbf{H}^{\text{SO}} | 1110 \rangle &= \frac{A}{2} [1(1+1) - 1(1+1) - 1(1+1)] = -A \\ \langle 1110 | &= \frac{-1}{\sqrt{2}} [\langle 11 - 11 | - \langle 111 - 1 |] \\ \text{So } \langle 1110 | \mathbf{H}^{\text{SO}} | 1110 \rangle &= \frac{A}{2} [\langle 11 - 11 | \mathbf{H}^{\text{SO}} | 11 - 11 \rangle + \langle 111 - 1 | \mathbf{H}^{\text{SO}} | 111 - 1 \rangle] = \frac{A}{2} [-1 - 1] = -A \\ \langle 1110 | \mathbf{H}^{\text{SO}} | 1100 \rangle &= 0 \\ |1110\rangle &= \frac{-1}{\sqrt{2}} [|11 - 11\rangle - |111 - 1\rangle] \\ |1100\rangle &= \frac{1}{\sqrt{3}} [|11 - 11\rangle - |1100\rangle + |111 - 1\rangle] \\ \text{So } \langle 1110 | \mathbf{H}^{\text{SO}} | 1100 \rangle &= \frac{-A}{\sqrt{6}} [\langle 11 - 11 | \mathbf{H}^{\text{SO}} | 11 - 11 \rangle - \langle 111 - 1 | \mathbf{H}^{\text{SO}} | 111 - 1 \rangle - \langle 11 - 11 | \mathbf{H}^{\text{SO}} | 1100 \rangle + \\ &+ \langle 111 - 1 | \mathbf{H}^{\text{SO}} | 1100 \rangle] = \frac{-A}{\sqrt{6}} (-1 + 1 - 1 + 1) = 0 \\ \langle 1100 | \mathbf{H}^{\text{SO}} | 1100 \rangle &= \frac{A}{2} (0(0+1) - 1(1+1) - 1(1+1)) = -2A \\ |1100\rangle &= \frac{1}{\sqrt{3}} [|11 - 11\rangle - |1100\rangle + |111 - 1\rangle] \\ \text{So } \langle 1100 | \mathbf{H}^{\text{SO}} | 1100 \rangle &= \frac{1}{3} [\langle 11 - 11 | \mathbf{H}^{\text{SO}} | 11 - 11 \rangle + \langle 111 - 1 | \mathbf{H}^{\text{SO}} | 111 - 1 \rangle - \langle 1100 | \mathbf{H}^{\text{SO}} | 11 - 11 \rangle - \\ &+ \langle 1100 | \mathbf{H}^{\text{SO}} | 111 - 1 \rangle - \langle 11 - 11 | \mathbf{H}^{\text{SO}} | 1100 \rangle - \langle 111 - 1 | \mathbf{H}^{\text{SO}} | 1100 \rangle] = \frac{A}{3} [-1 - 1 - 1 - 1 - 1 - 1] = -2A \end{aligned}$$

3. Calculate the energies for the hydrogenic systems H and Li^{2+} in the following states:

$$2^2P_{1/2} \text{ (means } n = 2, s = 1/2, \ell = 1, j = 1/2)$$

$$2^2P_{3/2}$$

$$3^2P_{1/2}$$

$$3^2P_{3/2}$$

$$3^2D_{3/2}$$

$$3^2D_{5/2}$$

Please express "energies" in cm^{-1} : $\sigma = \frac{E}{hc} \text{cm}^{-1}$ and locate the zero of energy at $n = \infty$.

ANSWER:

$$E_{n,\ell,s,j} = E_n^\circ + E_{\text{spin-orbit}}$$

$$E_n^\circ = \frac{-\mathfrak{R}Z^2}{n^2} \left(\frac{\mu}{m_e} \right) (\text{cm}^{-1})$$

\mathfrak{R} = Rydberg constant

Z = nuclear charge

$$\frac{\mu}{m_e} = \frac{m_e m_N}{m_e(m_e + m_N)}$$

and

$$E_{\text{spin-orbit}} = \left(\frac{5.90 \text{ cm}^{-1}}{2} \right) \frac{Z^4 [j(j+1) - \ell(\ell+1) - s(s+1)]}{n^3 \left(\ell + \frac{1}{2} \right) (\ell + 1) \ell}$$

For Hydrogen $Z = 1$, $\frac{\mu}{m_e} = 0.999456$, $\mathfrak{R} = 109737.42 \text{ cm}^{-1}$

For Li^{2+} $Z = 3$, $\frac{\mu}{m_e} = 0.9999218$.

Term	Hydrogen			Lithium		
	E_n°	E_{S-0}	E	E_n°	E_{S0}	E
$2^2P_{1/2}$	-27919.43	-.246	-27419.68	-246889.89	-19.91	-109798.75
$2^2P_{3/2}$	-27919.43	.123	-27419.31	-246889.89	9.96	-109718.88
$3^2P_{1/2}$	-12186.41	-.073	-12186.48	-109728.89	-5.9	-109734.79
$3^2P_{3/2}$	-12186.41	.036	-12186.37	-109728.89	2.9	-109725.99
$3^2D_{3/2}$	-12186.41	.022	-12186.39	-109728.89	-1.77	-109730.61
$3^2D_{5/2}$	-12186.41	.014	-12186.40	-109728.89	1.18	-109727.66

4. Consider the $(nd)^2$ configuration.

- (a) There are 10 distinct spin-orbitals associated with nd ; how many Pauli-allowed $(nd)^2$ Slater determinants can you form using two of these spin-orbitals?

ANSWER:

$(nd)^2$ available orbitals $2^+, 2^-, 1^+, 1^-, 0^+, 0^-, -1^+, -1^-, -2^+, -2^-$.

$M_L \backslash M_S$	1	0	-1
4	—	$(2^+, 2^-)$	—
3	$(2^+, 1^+)$	$(2^+, 1^-) (2^-, 1^+)$	$(2^-, 1^-)$
2	$(2^+, 0^+)$	$(2^-, 0^+) (2^+, 0^-) (1^+, 1^-)$	$(2^-, 0^-)$
1	$(1^+, 0^+) (2^+, -1^+)$	$(1^+, 0^-) (1^-, 0^+) (2^+, -1^-) (2^-, -1^+)$	$(1^-, 0^-) (2^-, 1^-)$
0	$(2^+, -2^+) (1^+, -1^+)$	$(2^+, -2^-) (2^-, -2^+) (1^+, -1^-) (1^-, -1^+) (0^+, 0^-)$	$(2^-, -2^-) (1^-, -1^-)$

\Rightarrow there are 45 $(nd)^2$ Slater Determinants.

- (b) What are the $L - S$ states associated with the $(nd)^2$ configuration? Does the sum of their degeneracies agree with the configurational degeneracy in part (a)?

ANSWER: $L - S$ states: $^1G, ^3F, ^3P, ^1D, ^1S$.

- (c) What is the lowest energy triplet state ($S = 1$) predicted by Hund's rules? Does Hund's rule predict the lowest energy singlet state?

ANSWER: Hund's rules say 3F will be lowest. Hund's rules don't really apply to other than the ground state so the lowest singlet state is not predicted.

- (d) Calculate the energies of all states (neglecting spin-orbit splitting) which arise from $(nd)^2$ in terms of the radial energy parameters F^0 , F^2 , and F^4 . [This is a long and difficult problem. The similar $(np)^2$ problem is worked out in detail in Condon and Shortley, pages 191-193, and in Tinkham, pages 177-178. The result for $(nd)^2$ is also given, without explanation and in slightly different notation, Condon and Shortley, page 202.] What relationship between F^2 and F^4 is required by Hund's rules?

ANSWER:

$$E(^1G) = \langle (2^+, 2^-) | \mathbf{H} | (2^+, 2^-) \rangle$$

$$E(^3F) = \langle (2^+, 1^+) | \mathbf{H} | (2^+, 1^+) \rangle$$

$$E(^3P) = \langle (1^+, 0^+) | \mathbf{H} | (1^+, 0^+) \rangle + \langle (2^+, -1^+) | \mathbf{H} | (2^+, -1^+) \rangle - E(^3F)$$

$$E(^1D) = \langle (2^-, 0^+) | \mathbf{H} | (2^-, 0^+) \rangle + \langle (2^+, 0^-) | \mathbf{H} | (2^+, 0^-) \rangle + \langle (1^+, 1^-) | \mathbf{H} | (1^+, 1^-) \rangle - E(^1G) - E(^3F)$$

$$E(^1S) = \langle (2^+, -2^-) | \mathbf{H} | (2^+, -2^-) \rangle + \langle (2^-, -2^+) | \mathbf{H} | (2^-, -2^+) \rangle + \langle (1^+, -1^-) | \mathbf{H} | (1^+, -1^-) \rangle \\ + \langle (1^-, -1^+) | \mathbf{H} | (1^-, -1^+) \rangle + \langle (0^+, 0^-) | \mathbf{H} | (0^+, 0^-) \rangle - E(^1G) - E(^3F) - E(^3P) - E(^1D)$$

$$\mathbf{H} = \underbrace{\sum_i \left(\frac{p_i^2}{2m} - \frac{Ze^2}{r} \right)}_{\substack{\text{Equal for all terms} \\ \text{since they come from} \\ \text{same configuration}}} + \underbrace{\sum_{r>j} \frac{1}{r_{ij}}}_{\substack{\text{we only need} \\ \text{to evaluate} \\ \text{these}}}$$

These 2-electron matrix elements are of the form

$$\sum_{a>b} \left\langle (a, b) \left| \frac{1}{r_{ab}} \right| (a, b) \right\rangle = \sum_{a>b} (\langle a, b | g | a, b \rangle - \langle a, b | g | b, a \rangle)$$

since the number of electrons is 2 the sum is just one term:

$$= \langle a, b | g | a, b \rangle - \langle a, b | g | b, a \rangle \equiv J(ab) - K(ab)$$

4D ANSWER, continued:

where

$$J(ab) = \sum_{k=0}^{\infty} a^k (\ell^a m_\ell^a, \ell^b m_\ell^b) F^k (n^a \ell^a, n^b \ell^b)$$

$$K(ab) = \left[\sum_{k=0}^{\infty} b^k (\ell^a m_\ell^a, \ell^b m_\ell^b) G^k (n^a \ell^a, n^b \ell^b) \right] \delta_{m_s^a m_s^b}$$

where $b^k = (c^k)^2$

so

$$E(^1G) = F^{\circ} + \frac{4}{49}F^2 + \frac{1}{441}F^4$$

$$E(^3F) = F^{\circ} - \frac{8}{49}F^2 - \frac{9}{441}F^4$$

$$E(^3P) = F^{\circ} + \frac{7}{49}F^2 - \frac{84}{441}F^4$$

$$E(^1D) = F^{\circ} - \frac{3}{49}F^2 + \frac{36}{441}F^4$$

$$E(^1S) = F^{\circ} + \frac{14}{49}F^2 + \frac{126}{441}F^4$$

Since Hund's rules say that $E(^3F) < E(^3P) \Rightarrow -\frac{8}{49}F^2 - \frac{9}{441}F^4 < \frac{7}{49}F^2 - \frac{84}{441}F^4 \quad \therefore \frac{F^2}{F^4} > 0.56$

5. If an atom is in a $(2p)^2 \ ^3P_0$ state, to which of the following states is an electric dipole transition allowed? Explain in each case.

(a) $2p3d \ ^3D_2$ **ANSWER:** Forbidden since $\Delta J = 2$.(b) $2s2p \ ^3P_1$ **ANSWER:** Allowed since $\Delta J = 1$, $\Delta \ell = -1$ & $\Delta L = 0$, $\Delta S = 0$.(c) $2s3s \ ^3S_1$ **ANSWER:** Forbidden in the absence of configuration interaction since it is a 2–electron transition.(d) $2s2p \ ^1P_1$ **ANSWER:** Forbidden since $\Delta S \neq 0$.