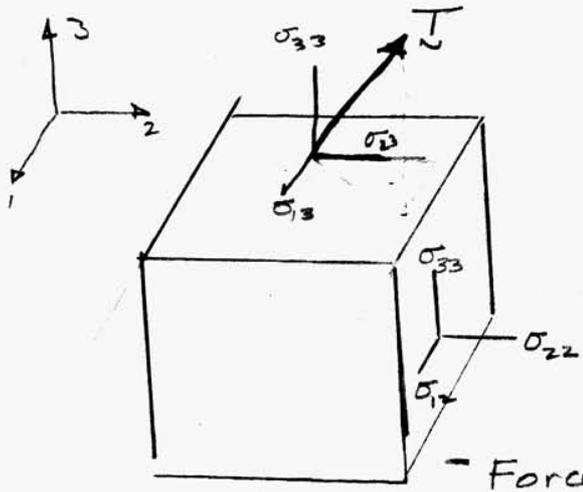


Stress



Traction \vec{T} on material may depend on area

Body forces \vec{g} depend on volume neglect in this discussion

- Forces inside body which are a reaction to traction are stresses
- Homogeneous if forces acting on a surface of fixed shape and orientation do not depend on position in body.

Then traction on each face of unit-cube can be characterized.

For body to be in static equilibrium, Traction on (+2 face) = (-2 face) and opposite in dir otherwise net imbalance causes acceleration governed by Newton's Law.

Define stress on each face as traction / area

$\Rightarrow \vec{T} = \sigma \vec{A}$ i.e. stress tensor relates area normal \rightarrow traction

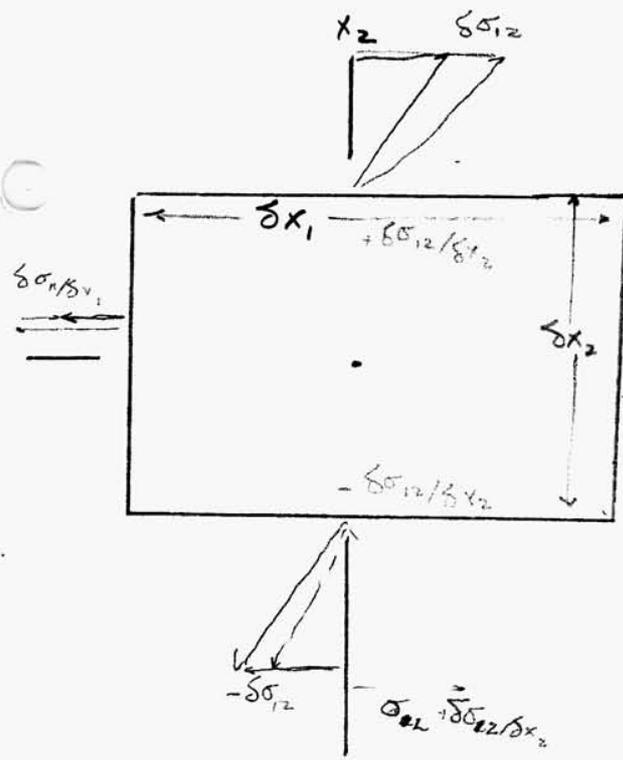
$$\vec{T}_i = \sigma_{ij} A_j$$

$\sigma_{ij} \triangleq T_i$ on j face

- normal components, σ_{ii} + if traction
- shear components $\sigma_{ij}, i \neq j$ + if in positive dir on " face
- + if in negative dir on " face



Egn. Equilibrium



On x_1 face in x_1 direct

left $\sigma_{11} \delta x_2 \delta x_3 + \frac{\partial \sigma_{11}}{\partial x_1} \frac{\delta x_1}{2} \delta x_2 \delta x_3$

right $-\left[\sigma_{11} \delta x_2 \delta x_3 - \frac{\partial \sigma_{11}}{\partial x_1} \frac{\delta x_1}{2} \delta x_2 \delta x_3 \right]$

result. $\frac{\partial \sigma_{11}}{\partial x_1} (\delta x_1 \delta x_2 \delta x_3)$

on x_2 in x_1 direction

top $\sigma_{12} (\delta x_1 \delta x_3) + \frac{\partial \sigma_{12}}{\partial x_2} \frac{\delta x_2}{2} (\delta x_1 \delta x_3)$

bottom $-\left[\sigma_{12} (\delta x_1 \delta x_3) + \frac{\partial \sigma_{12}}{\partial x_2} \left(-\frac{\delta x_2}{2} \right) (\delta x_1 \delta x_3) \right]$

$\frac{\partial \sigma_{12}}{\partial x_2} (\delta x_2 \delta x_1 \delta x_3)$

By Newton's law, sum of forces in x_1 direction equals mass/vol \times acceleration in that direction

$$\frac{\partial \sigma_{11}}{\partial x_1} + \frac{\partial \sigma_{12}}{\partial x_2} + \frac{\partial \sigma_{13}}{\partial x_3} (+ \rho g_1) = \rho \ddot{x}_1$$

In general,

$$\frac{\partial \sigma_{ij}}{\partial x_j} + \rho g_i = \rho \ddot{x}_i$$

If

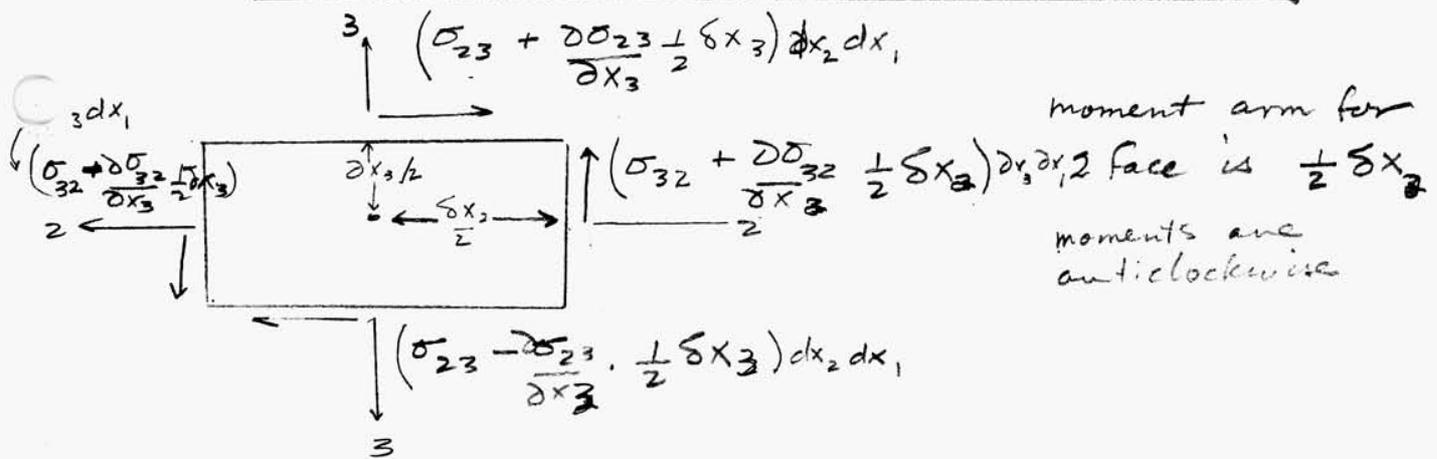
$\ddot{x}_i = 0$, i.e.

in static eq.

$$\frac{\partial \sigma_{ij}}{\partial x_j} + \rho g_i = 0$$

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Shear moments and symmetry



$$2\sigma_{32} \frac{dx_3}{2} (dx_2 dx_1) - 2\sigma_{23} \frac{dx_3}{2} dx_2 dx_1 + G_1 dx_1 dx_2 dx_3 = I_1 \ddot{\theta}_1$$

but assume no body torques $G_1 = 0$

and note I_1 order of mag $\rho \delta x^5 \rightarrow 0$ as $dx \downarrow$

then $(\sigma_{32} - \sigma_{23}) dx_1 dx_2 dx_3 = 0$

$$\Rightarrow \sigma_{32} = \sigma_{23}$$

or $\sigma_{ij} = \sigma_{ji}$

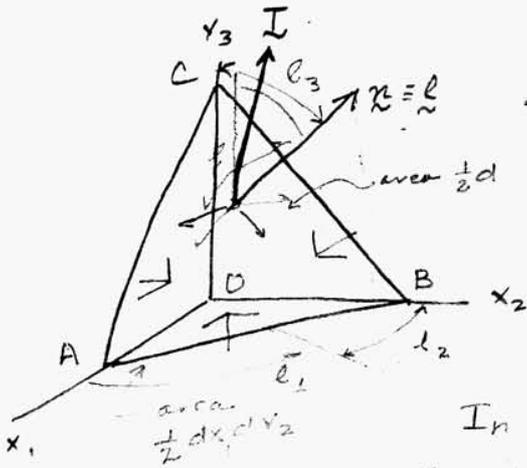
\Rightarrow Stress tensor is symmetric.

Proof that stress is a tensor

Cauchy tetrahedron

Total force on face with area ABC

$$\vec{T} = \vec{P} \cdot (\text{area } ABC)$$



In x_1 direction

$$P_1(\text{ABC}) = \sigma_{11}(\text{BOC}) + \sigma_{12}(\text{AOC}) + \sigma_{13}(\text{AOB})$$

$$P_1 = \sigma_{11} l_1 + \sigma_{12} l_2 + \sigma_{13} l_3$$

where $l_1 = \frac{\text{area } \text{BOC}}{\text{area } \text{ABC}} = \cos \angle x_1, \vec{n}$

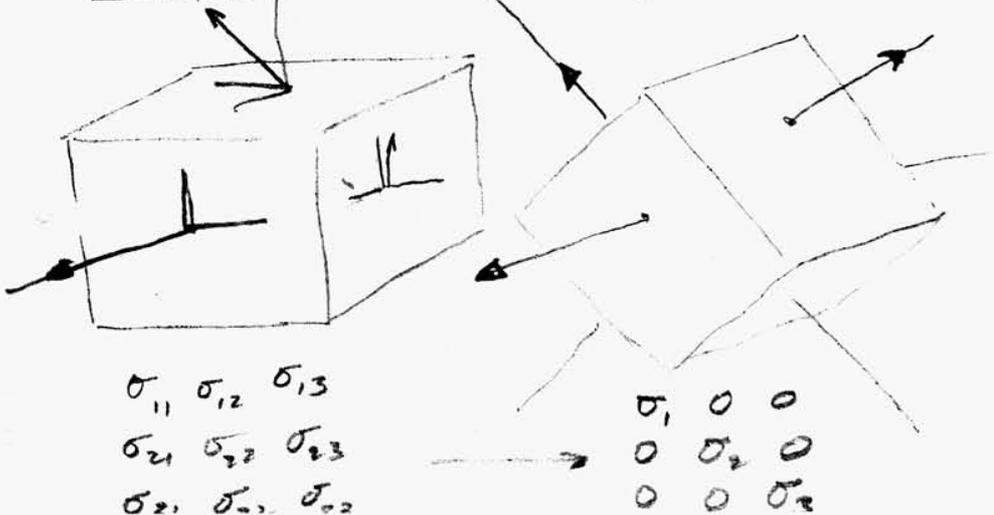
sim. $P_2 = \sigma_{21} l_1 + \sigma_{22} l_2 + \sigma_{23} l_3$

$P_3 = \sigma_{31} l_1 + \sigma_{32} l_2 + \sigma_{33} l_3$

or $P_i = \sigma_{ij} l_j$
 i.e. tensor transformation rule

⇒ Stress is a 2nd rank, symmetric tensor

⇒ Principal stresses, Quadratic



$$\begin{matrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{matrix}$$

$$\begin{matrix} \sigma_1 & 0 & 0 \\ 0 & \sigma_2 & 0 \\ 0 & 0 & \sigma_3 \end{matrix}$$

Special Forms of Stress

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i.) uniaxial stress

$$\begin{bmatrix} \sigma & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

ii.) biaxial stress

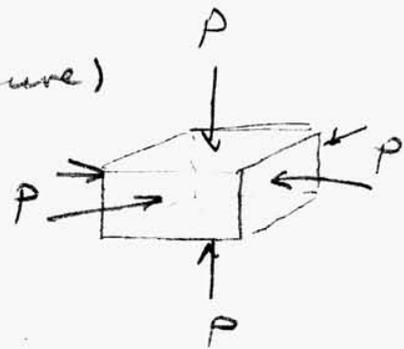
$$\begin{bmatrix} \sigma_1 & 0 & 0 \\ 0 & \sigma_2 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

iii.) triaxial stress

$$\begin{bmatrix} \sigma_1 & 0 & 0 \\ 0 & \sigma_2 & 0 \\ 0 & 0 & \sigma_3 \end{bmatrix}$$

iv.) hydrostatic stress (pressure)

$$\begin{bmatrix} -P & 0 & 0 \\ 0 & -P & 0 \\ 0 & 0 & -P \end{bmatrix} = -P\delta_{ij}$$



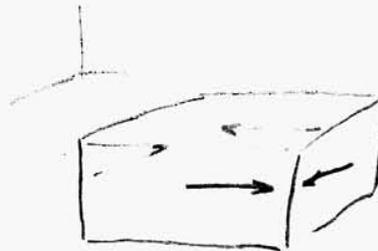
v.) Pure shear

$$\begin{bmatrix} -\sigma & 0 & 0 \\ 0 & \sigma & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

see above

vi.) Simple shear

$$\begin{bmatrix} 0 & \sigma & 0 \\ \sigma & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$



Rod stretch

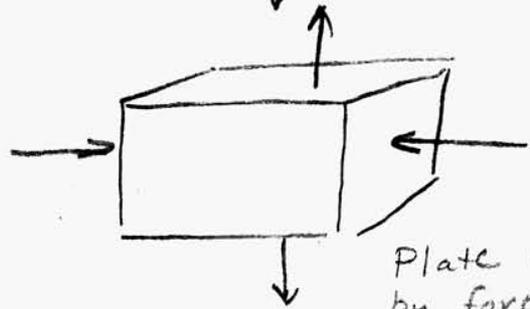
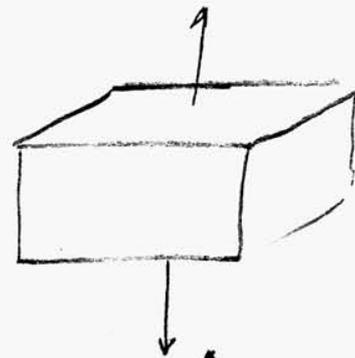


Plate loaded by forces ϵ_{comp} at end
 $\Delta x_1, \Delta x_2 \ll x_3$



Summary: Stress Tensor

1.) Traction on a plane with direction cosine \underline{l}

$$\underline{T} = \underline{\underline{\sigma}} \underline{l} \quad T_i = \sigma_{ij} l_j$$

2.) Eqs of motion

$$\frac{\partial \sigma_{ij}}{\partial x_i} + \rho g_j = m \ddot{x}_j$$

3.) Stress is symmetric

$$\sigma_{ij} = \sigma_{ji}$$

4.) Principal stress directions
values

$$\begin{matrix} \sigma_I & 0 & 0 \\ 0 & \sigma_{II} & 0 \\ 0 & 0 & \sigma_{III} \end{matrix}$$

5.) Stress quadric construction

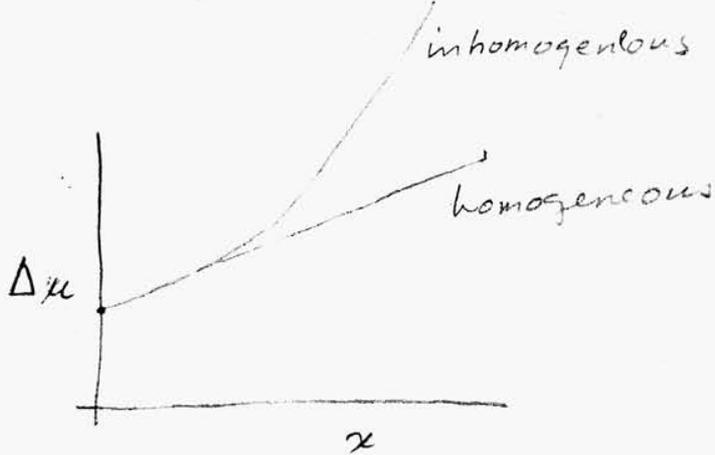
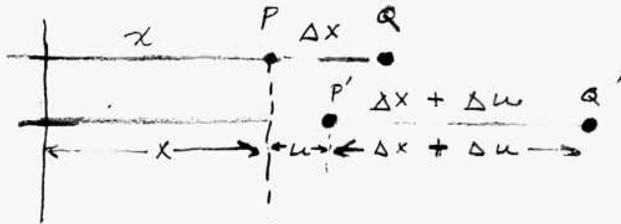
Stress, Strain, Elasticity:

1 of 8

Nye, Chap 5 & 6, pps 82-105.

1. One dimensional strain.

Relative displacements important



Strain at a point P is

$$\epsilon = \frac{\text{increase in length}}{\text{original length}} = \frac{P'Q' - PQ}{PQ} = \frac{\Delta u}{\Delta x}$$

(slope of curve above)

Goal: For a given deformation of body, Define a tensor that describes the change of direction and length of any vector in body.

Want to map vector in body (undef) into vector in body (def.)