

Faraday's Law (Induced emf)

Reading - Shen and Kong - Ch. 16

Outline

- Magnetic Flux and Flux Linkage
- Inductance
- Stored Energy in the Magnetic Fields of an Inductor
- Faraday's Law and Induced Electromotive Force (*emf*)
- Examples of Faraday's Law

Magnetic Flux

Φ [Wb] (Webers)

due to macroscopic
& microscopic

Magnetic Flux Density

B [Wb/m²] = T (Teslas)

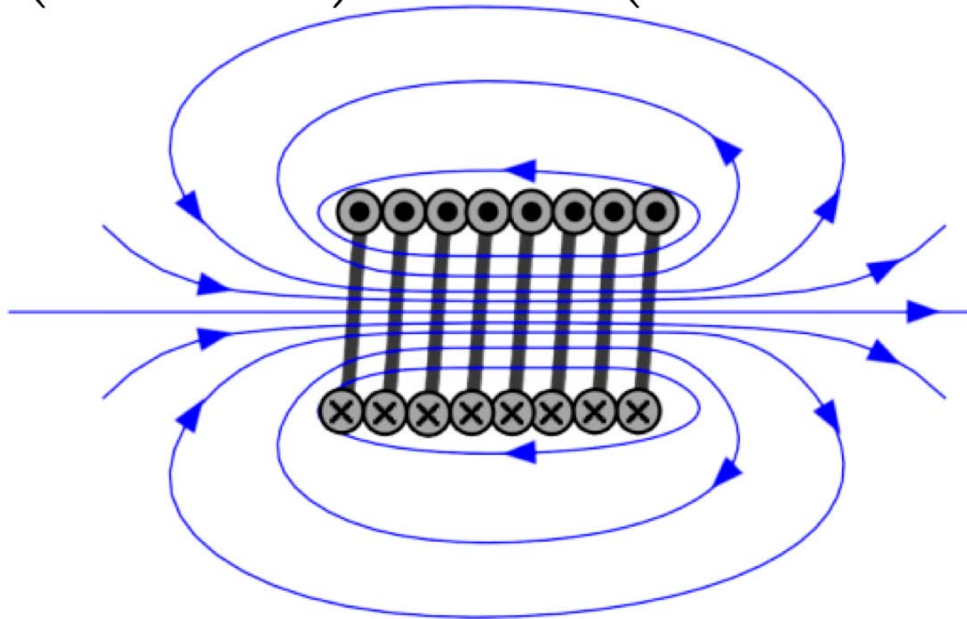
Magnetic Field Intensity

H [Amp-turn/m]

due to macroscopic
currents

$$\Phi = \int \bar{B} \cdot d\bar{A}$$

$$\bar{B} = \mu_0 (\bar{H} + \bar{M}) = \mu_0 (\bar{H} + \chi_m \bar{H}) = \mu_0 \mu_r \bar{H}$$



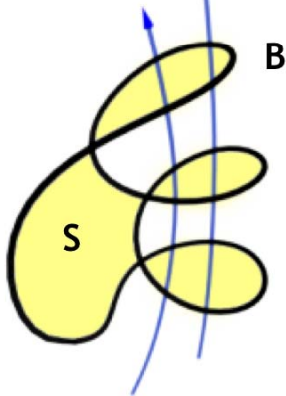
Flux Linkage of a Solenoids

FOR A SUFFICIENTLY LONG SOLENOID...

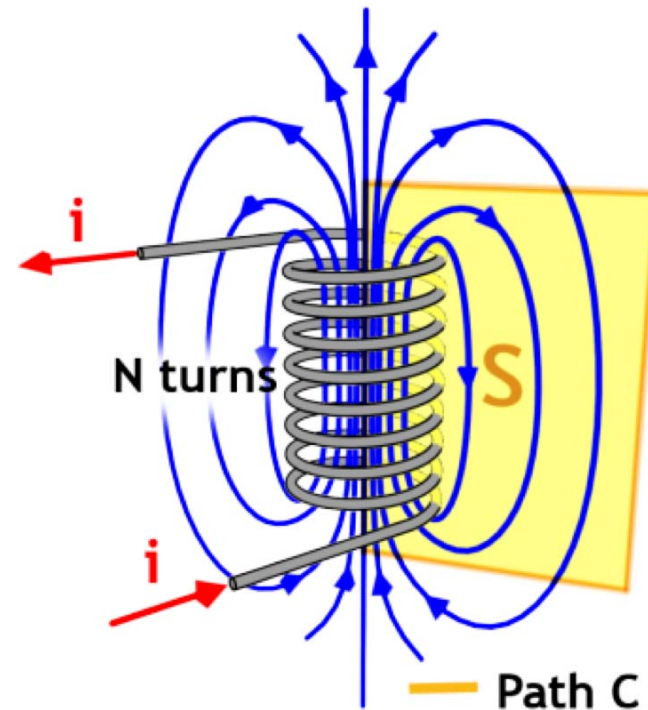
$$H_{inside} = \frac{Ni}{h}$$

$$B_{inside} = \mu_o \frac{Ni}{h}$$

$$\Phi_{inside} = \mu_o \frac{Ni}{h} A$$



In the solenoid the individual flux lines pass through the integrating surface S more than once



B = Magnetic flux density inside solenoid

A = Solenoid cross sectional area

N = Number of turns around solenoid

$$\lambda = N\Phi = NBA = \mu_o \frac{N^2 i}{h} A$$

... flux linked by solenoid

Inductors

... is a passive electrical component that stores energy in a magnetic field created by the electric current passing through it. (This is in equivalence to the energy stored in the electric field of capacitors.)

IN GENERAL:

$$\lambda = Li$$

FOR A LINEAR COIL: $\lambda = N\Phi = \mu_o \frac{N^2 i}{h} A = Li \Rightarrow L = \mu_o \frac{A}{h} N^2$

The magnetic permeability of the vacuum: $\mu_o = 4\pi \times 10^{-7} \text{ H/m}$ [henry per meter]

An inductor's ability to store magnetic energy is measured by its inductance, in units of henries. The henry (symbol: H) is named after Joseph Henry (1797-1878), the American scientist who discovered electromagnetic induction independently of and at about the same time as Michael Faraday (1791-1867) in England.

EQUIVALENCE OF UNITS:

$$H = \frac{\text{Wb}}{\text{A}} = \frac{\text{V} \cdot \text{s}}{\text{A}} = \Omega \cdot \text{s} = \frac{\text{J}}{\text{A}^2} = \frac{\text{J/C} \cdot \text{s}}{\text{C/s}} = \frac{\lambda = N\Phi = NBA = \mu_2 \frac{N^2 i}{h} A}{\text{J} \cdot \text{s}^2} = \frac{\text{m}^2 \cdot \text{kg}}{\text{C}^2} = \frac{\text{m}^2 \cdot \text{kg}}{\text{s}^2 \cdot \text{A}^2}$$

Stored Energy in an Inductor

FROM 8.02: $v(t) = L \frac{di(t)}{dt}$ ← voltage over an inductor

$$P_{elec} = v \cdot i = L \frac{di}{dt} \cdot i = \frac{1}{2} L \frac{d}{dt} i^2$$

If L is not a function of time ...

$$P_{elec} = \frac{d}{dt} \left(\frac{1}{2} L i^2 \right) = \frac{dW_s}{dt}$$

... where E is energy stored in the field of the inductor any instant in time

$$W_s(i, r) = \frac{1}{2} L i^2$$

FOR A LINEAR COIL: $L = \mu_o \frac{A}{h} N^2$

Calculation of energy stored in the inductor

$$W_s = \int v i dt = \int i \frac{d\lambda}{dt} dt = \int i d\lambda = \frac{1}{2} \frac{\lambda_0^2}{L}$$

Note that flux is $\lambda = N\phi = NA\mu_0 H_z$

And that inductance is $L = \mu_0 \frac{A}{h} N^2$

Energy stored is $W_s = \frac{1}{2} \frac{(NA\mu_0 H_z)^2}{\mu_0 \frac{A}{h} N^2} = \frac{1}{2} \mu_0 H_z^2 Ah$

Energy stored per volume $\frac{W_s}{Volume} = \frac{1}{2} \frac{(NA\mu_0 H_z)^2}{\mu_0 \frac{A}{h} N^2} = \frac{1}{2} \mu_0 H_z^2$

General: Stored Energy in the Coil

$$\text{FROM 8.02: } v(t) = L \frac{di(t)}{dt} \quad \leftarrow \text{voltage over an inductor}$$

$$\text{Since } \lambda = Li \text{ then } v(t) = \frac{d\lambda(t)}{dt}$$

→ Change in the magnetic flux within the inductor generates voltage

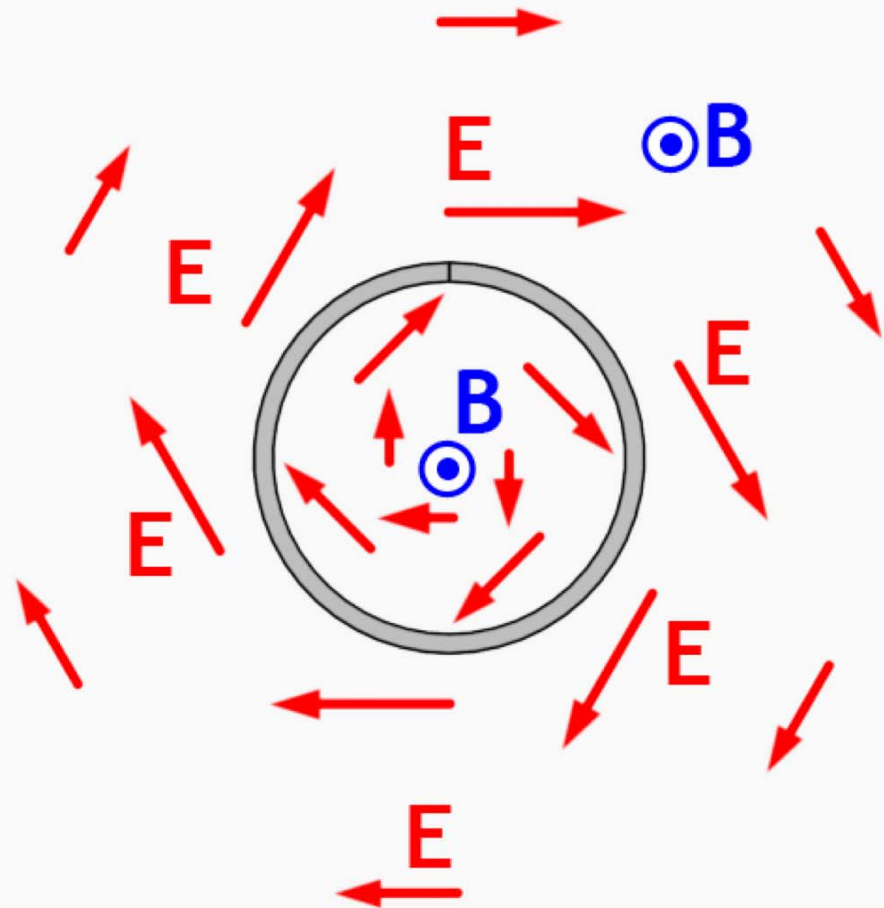
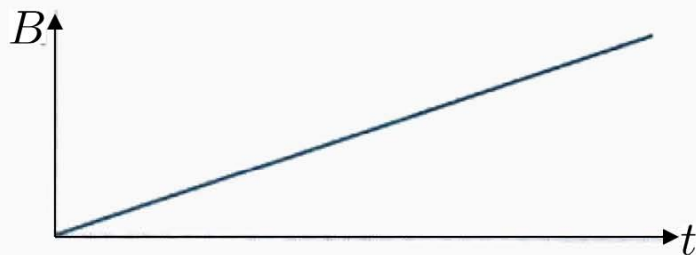
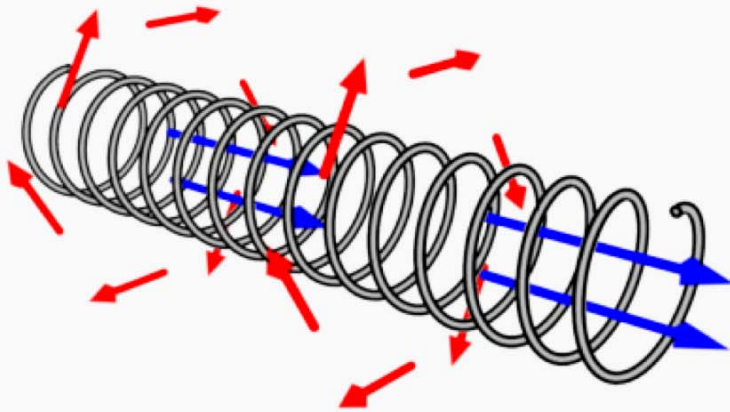
$$P_{elec} = v \cdot i = \frac{d\lambda}{dt} \cdot i$$

$$W_s = \int P_{elec} dt = \int i \frac{d\lambda}{dt} dt = \int i d\lambda$$

... where W_s is energy stored in the field of the inductor any instant in time

Induced electromotive force (emf)

Michael Faraday in 1831 noticed that time-varying magnetic field produces an emf in a solenoid



from Chabay and Sherwood, Ch 22

THEREFORE, THERE ARE TWO WAYS TO PRODUCE ELECTRIC FIELD

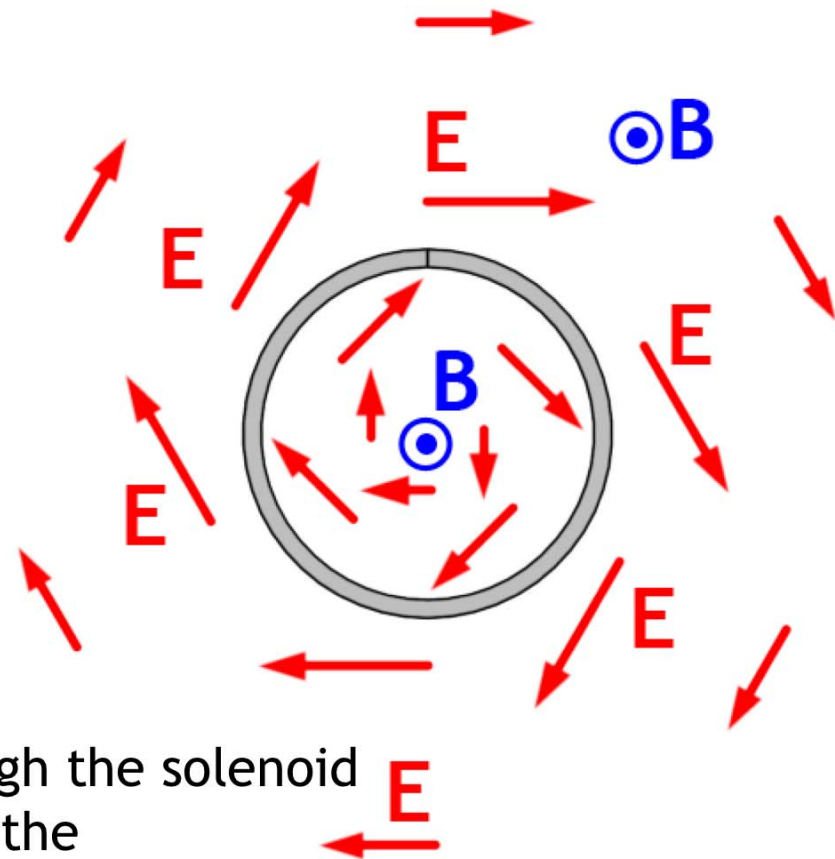
- (1) Coulomb electric field is produced by electric charges according to Coulomb's law:

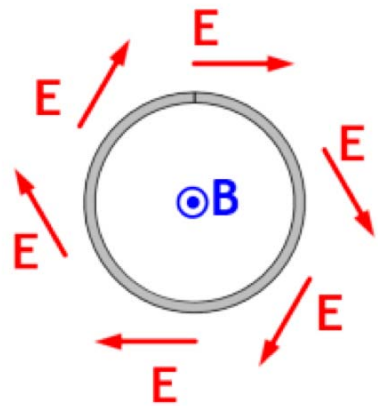
$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$

- (2) Non-Coulomb electric field E_{NC} is associated with time-varying magnetic flux density dB/dt

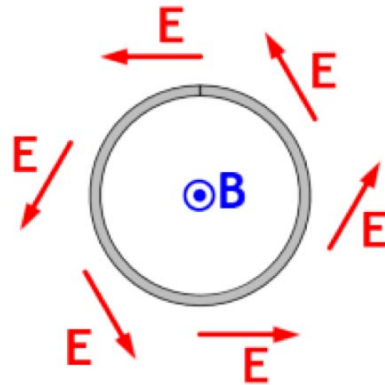
For a solenoid, E_{NC}

- curls around a solenoid
- is proportional to $-dB/dt$ through the solenoid
- decreases with $1/r$, where r is the radial distance from the solenoid axis

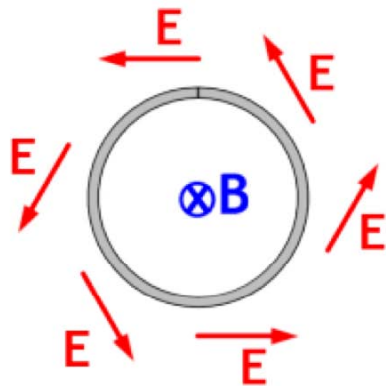




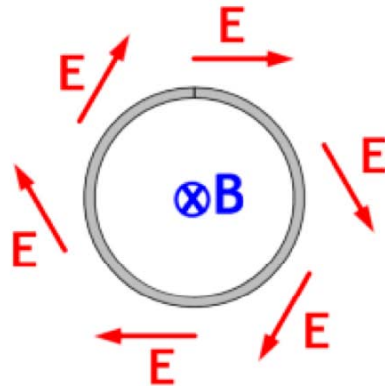
B out, *increasing*
 $-\frac{dB}{dt}$ into page



B out, *decreasing*
 $-\frac{dB}{dt}$ out of page



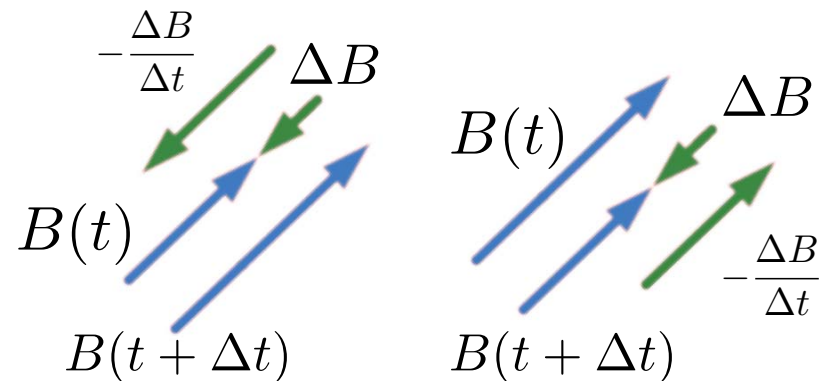
B in, *increasing*
 $-\frac{dB}{dt}$ out of page



B in, *decreasing*
 $-\frac{dB}{dt}$ into page

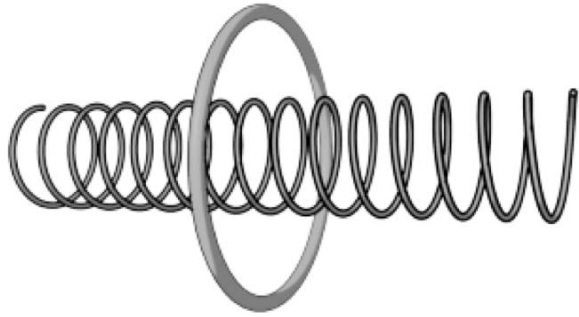
What is the Direction of Magnetically Induced (non-Coulomb) Field, E_{NC} ?

Find the change in the magnetic flux density as a basis for determining the direction of $-dB_1/dt$



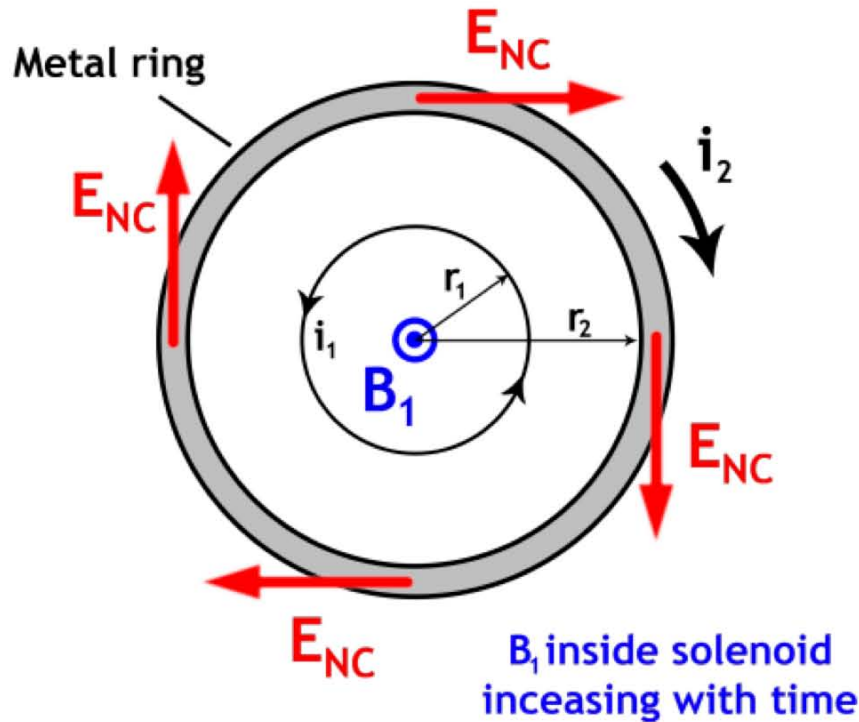
Lenz's Rule

The induced electric field would drive the current in the direction to make the magnetic field that attempts to keep the flux constant

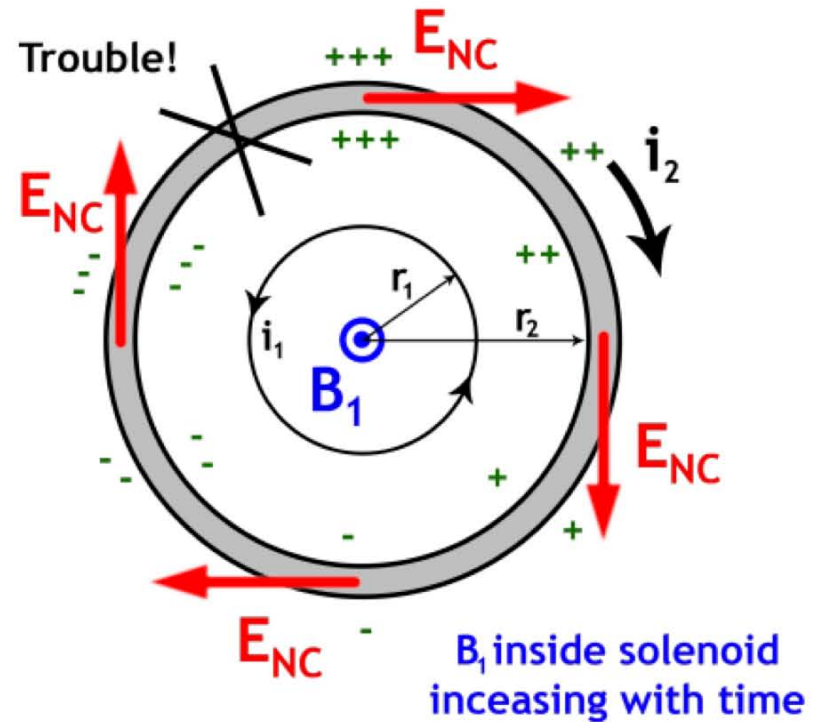


METAL RING IS PLACED AROUND A SOLENOID

Magnetically Induced
(non-Coulomb) E_{NC}
Drives Current in a Loop
Surrounding the Solenoid



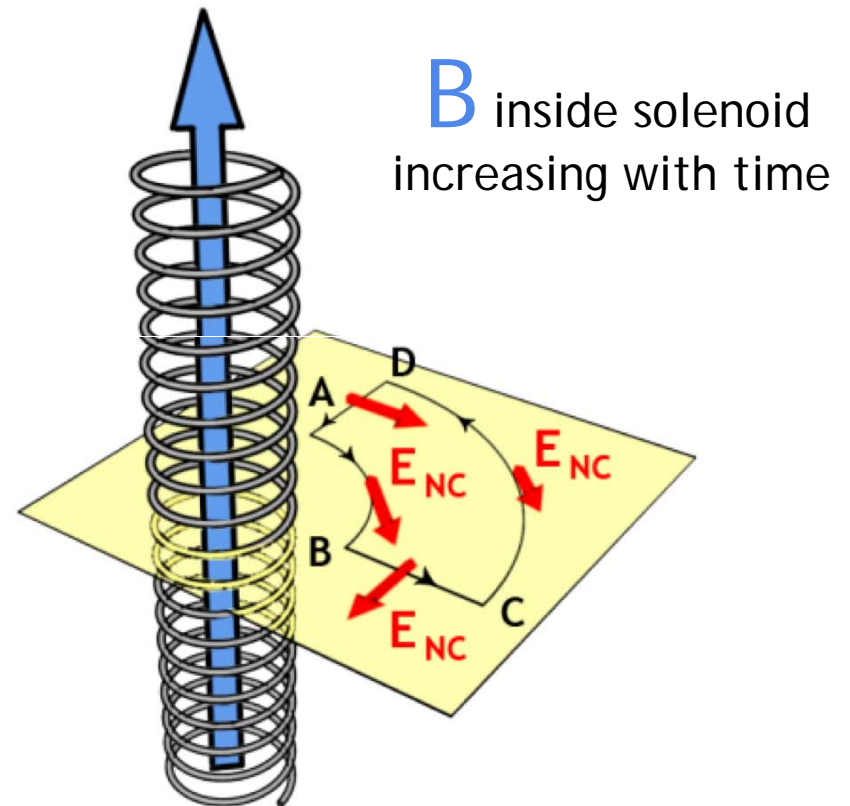
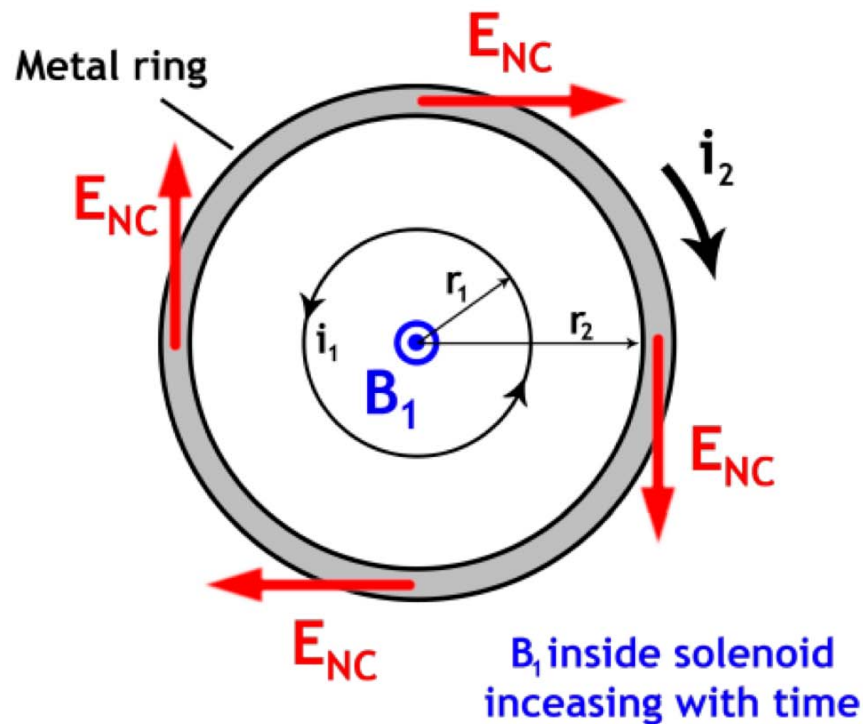
END VIEW: The non-Coulomb electric field drives a current i_2 in the ring



This pattern of surface charge is impossible, because it would imply a huge E at the marked location, and in the wrong direction !

$$emf = \oint \vec{E}_{NC} \cdot d\vec{l} = E_{NC} (2\pi r_2)$$

$$I_2 = emf / R \quad \text{where } R \text{ is the ring resistance}$$

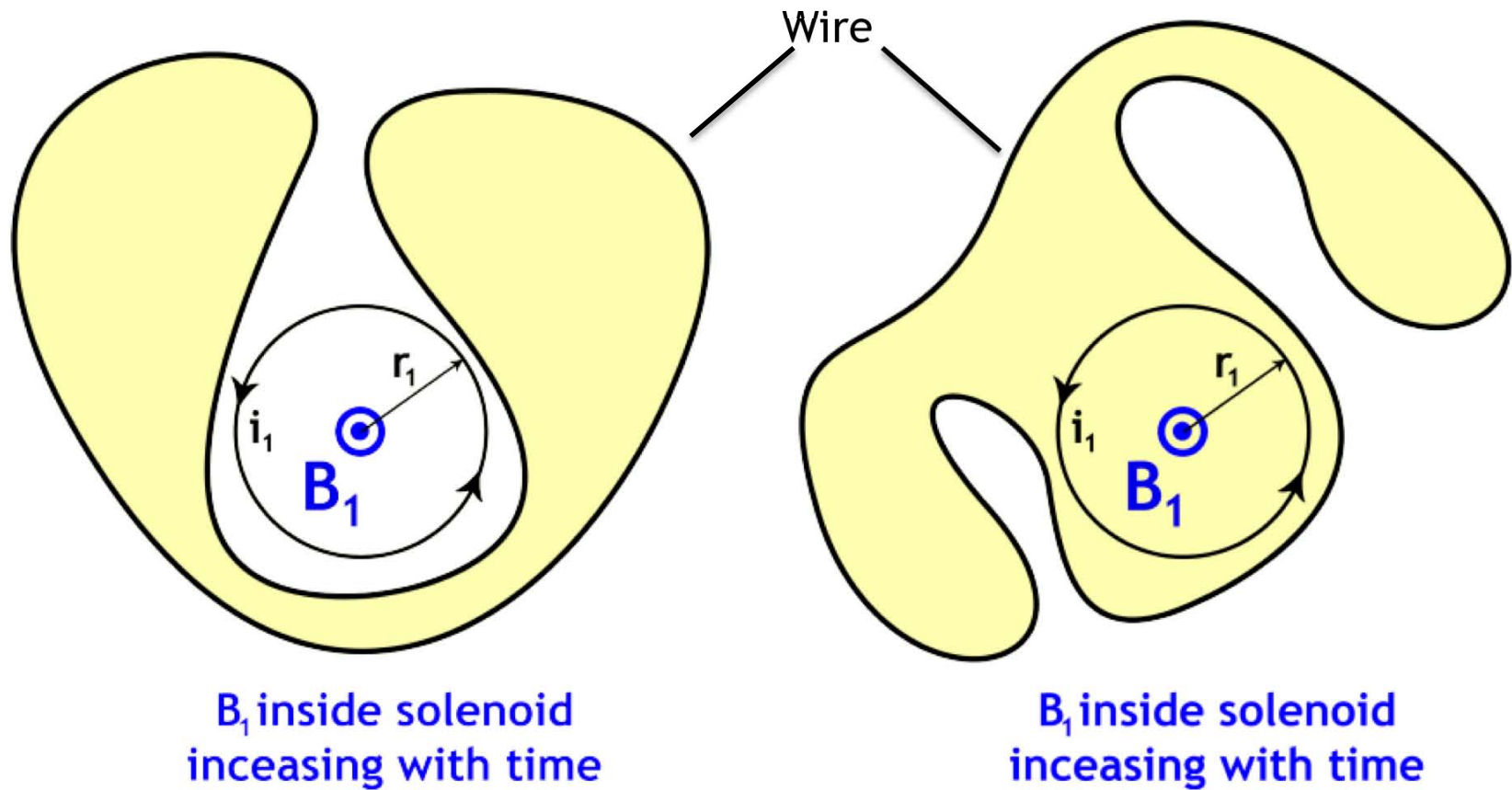


Integral of E_{NC} along a path that does not encircle the solenoid is zero since $E_{NC} \sim 1/r$

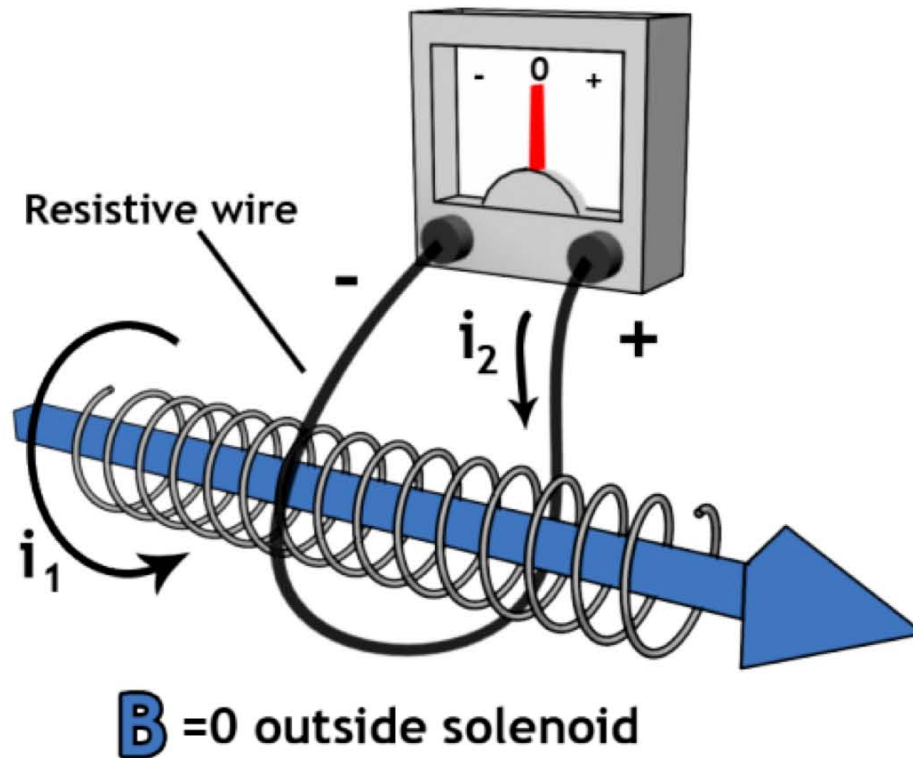
What if we double the value of r_2 ? \rightarrow Still get the same emf around the loop

emf in a ring encircling the solenoid is the same for any radius

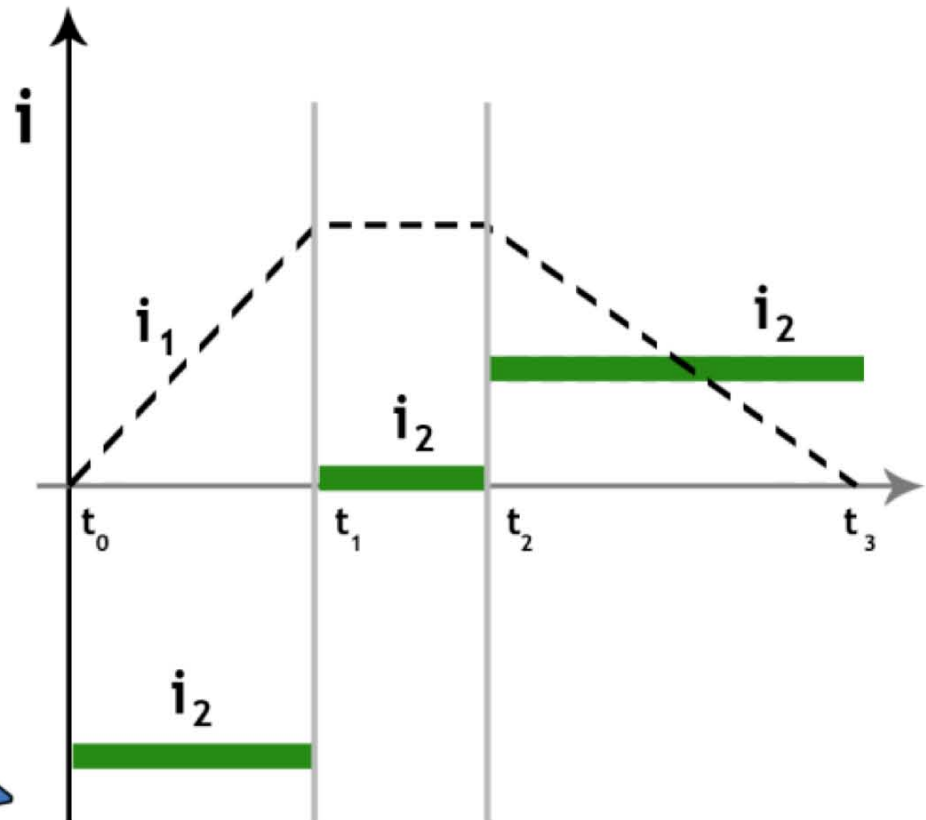
Will Current Run in these Wires ?



An ammeter measures current in a loop surrounding the solenoid. Initially I_1 is constant, so B_1 is constant, and no current runs through the ammeter.



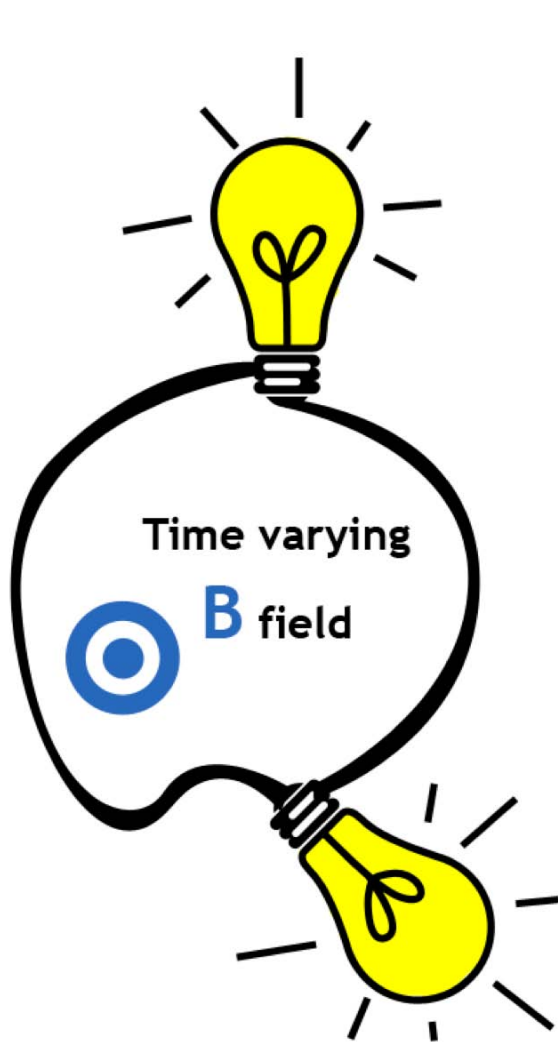
Example



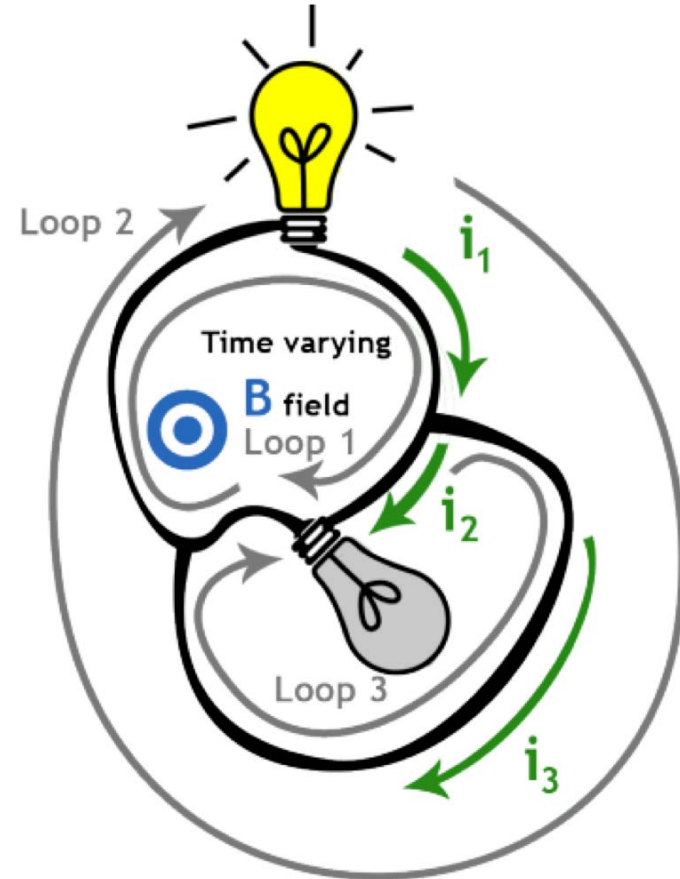
Vary the solenoid current I_1 and observe the current I_2 that runs in the outer wire, through the ammeter

from Chabay and Sherwood, Ch 22

Peculiar Circuit - Two Bulbs Near a Solenoid



Two light bulbs connected around a long solenoid with varying B.



... Add a thick copper wire.

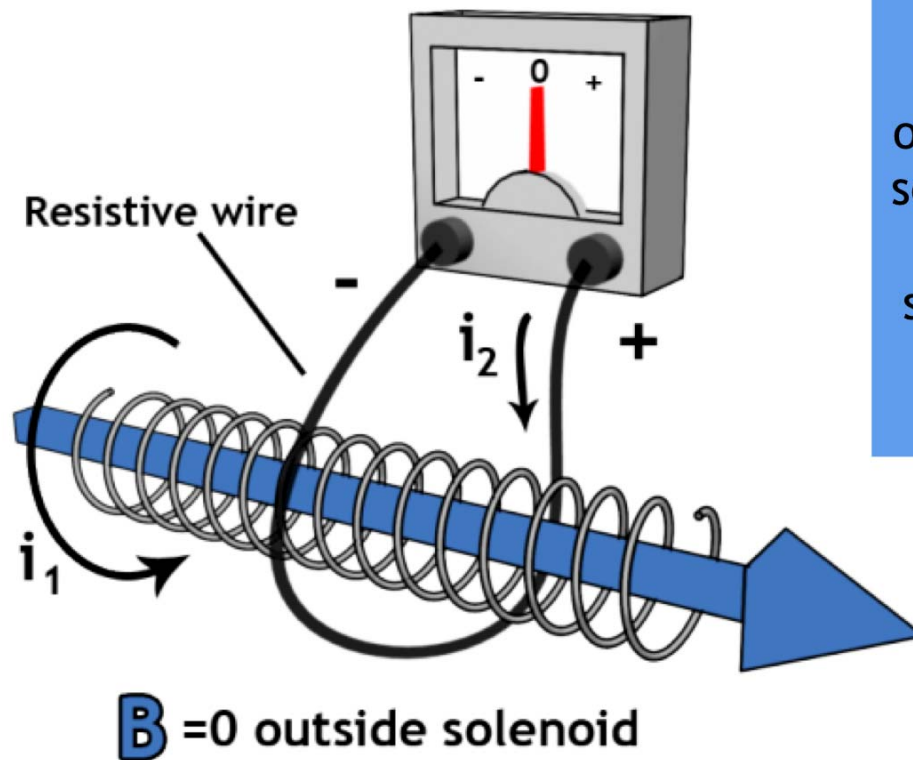
$$\text{Loop 1: } \mathcal{E} - R_1 I_1 - R_2 I_2 = 0$$

$$\text{Loop 2: } \mathcal{E} - R_1 I_1 = 0$$

$$\text{Loop 3: } R_2 I_2 = 0 \text{ (no flux enclosed)}$$

Question:


If we use a solenoid with twice the cross-sectional area, but the same magnetic flux density (same magnitude of I_1), what is the magnitude of I_2 ?



We can build the solenoid with the larger cross-sectional area, $2 \cdot A$, out of two solenoids with the initial cross-sectional area, A . Each of the smaller solenoids would induce current I_2 , so by superposition, for the twice-as-big solenoid the current would be twice-as-big !

$$|emf| = \left| \frac{d}{dt}(B_1 A) \right|$$

$$|emf| = \left| \frac{d}{dt}(B_{\perp} A) \right| = \left| \frac{d}{dt}(\Phi_{mag}) \right|$$


$$\Phi_{mag} = \int \vec{B} \cdot d\hat{a} = \int B_{\perp} da$$

Faraday's Law

The induced *emf* along a round-trip path is equal to the rate of change of the magnetic flux on the area encircled by the path.

$$emf = -\frac{d\Phi_{mag}}{dt}$$

$$\oint \vec{E}_{NC} \cdot d\hat{l} = -\frac{d}{dt} \int \vec{B} \cdot d\hat{a}$$

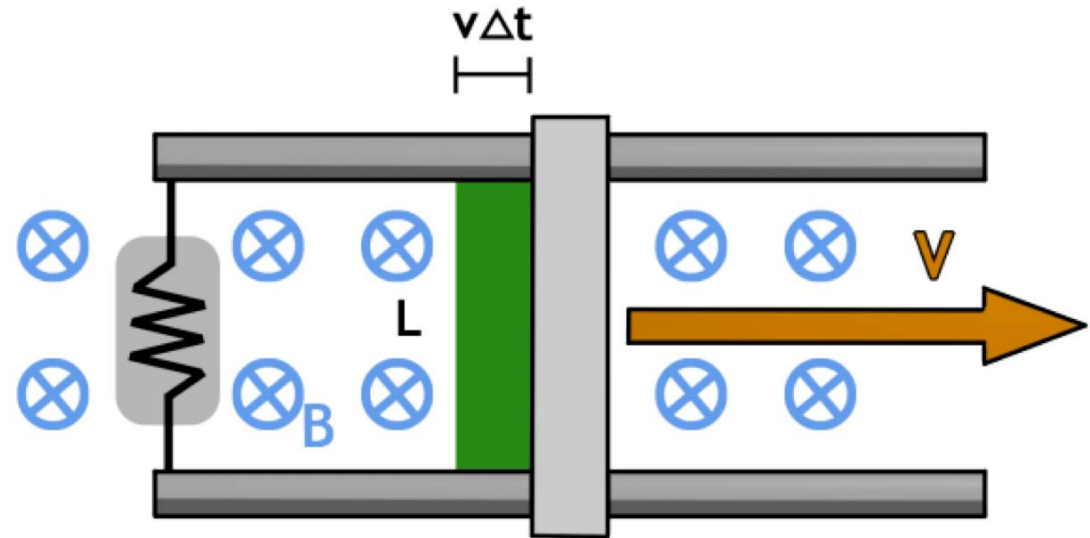
Faraday's Law and Motional emf

What is the emf over the resistor ?

$$emf = -\frac{d\Phi_{mag}}{dt}$$

In a short time Δt the bar moves a distance $\Delta x = v \cdot \Delta t$, and the flux increases by
 $\Delta\Phi_{mag} = B (L v \cdot \Delta t)$

$$emf = \frac{\Delta\Phi_{mag}}{\Delta t} = BLv$$

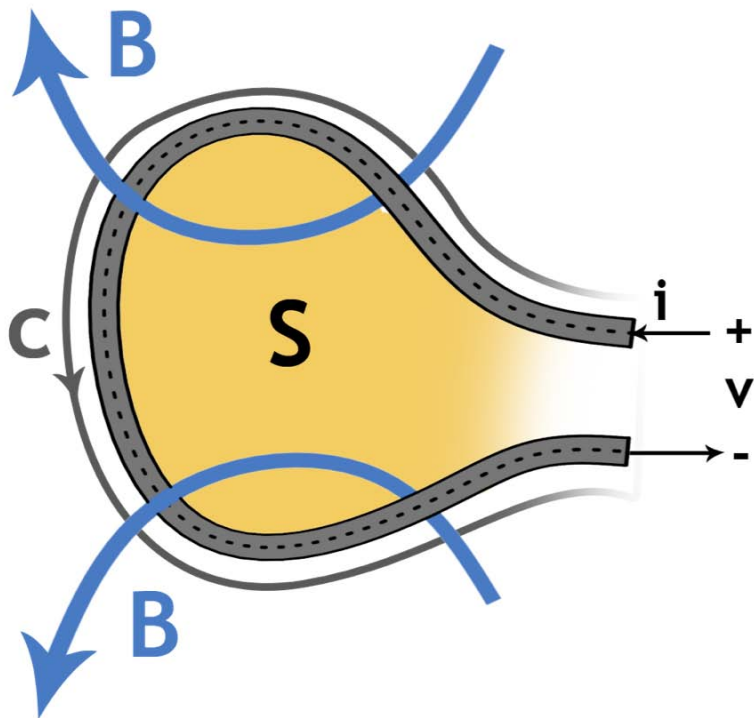


There is an increase in flux through the circuit as the bar of length L moves to the right (orthogonal to magnetic field H) at velocity, v .

Terminal Voltages & Inductance

Assume:

- Perfectly conducting wire
- Stationary contour C
- Negligible magnetic flux at the terminals



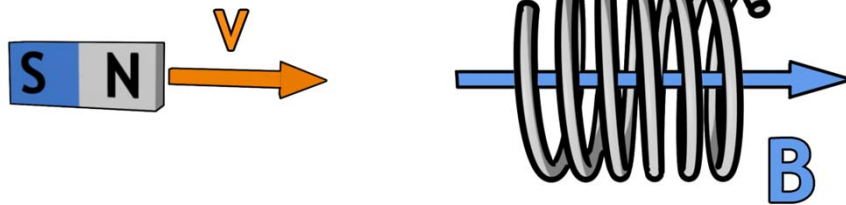
$$emf = d\lambda/dt$$

If the current i created the magnetic flux density B , then the flux linkage is given by $\lambda = Li$.
In this case, $emf = L di/dt$.
 L is the self inductance of the coil.

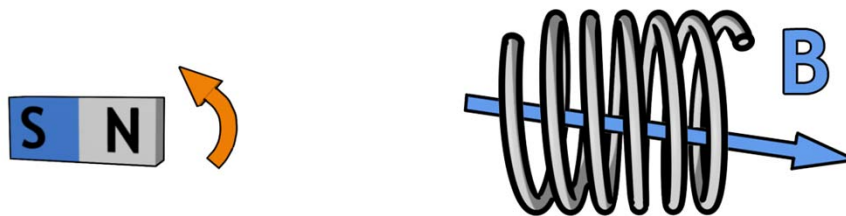
Faraday's Law for a Coil

The induced emf in a coil of N turns is equal to N times the rate of change of the magnetic flux on one loop of the coil.

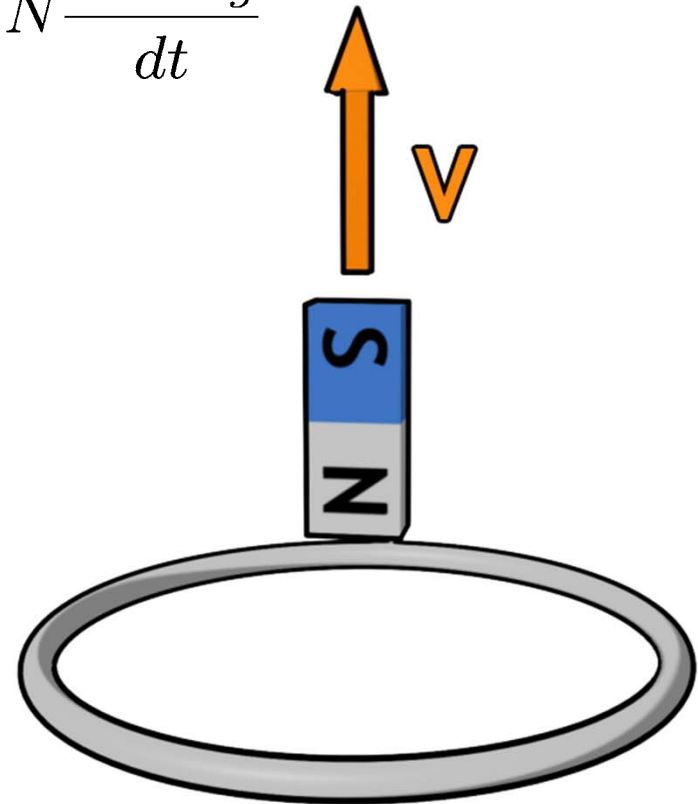
$$emf = -N \frac{d\Phi_{mag}}{dt}$$



Moving a magnet towards a coil produces a time-varying magnetic field inside the coil

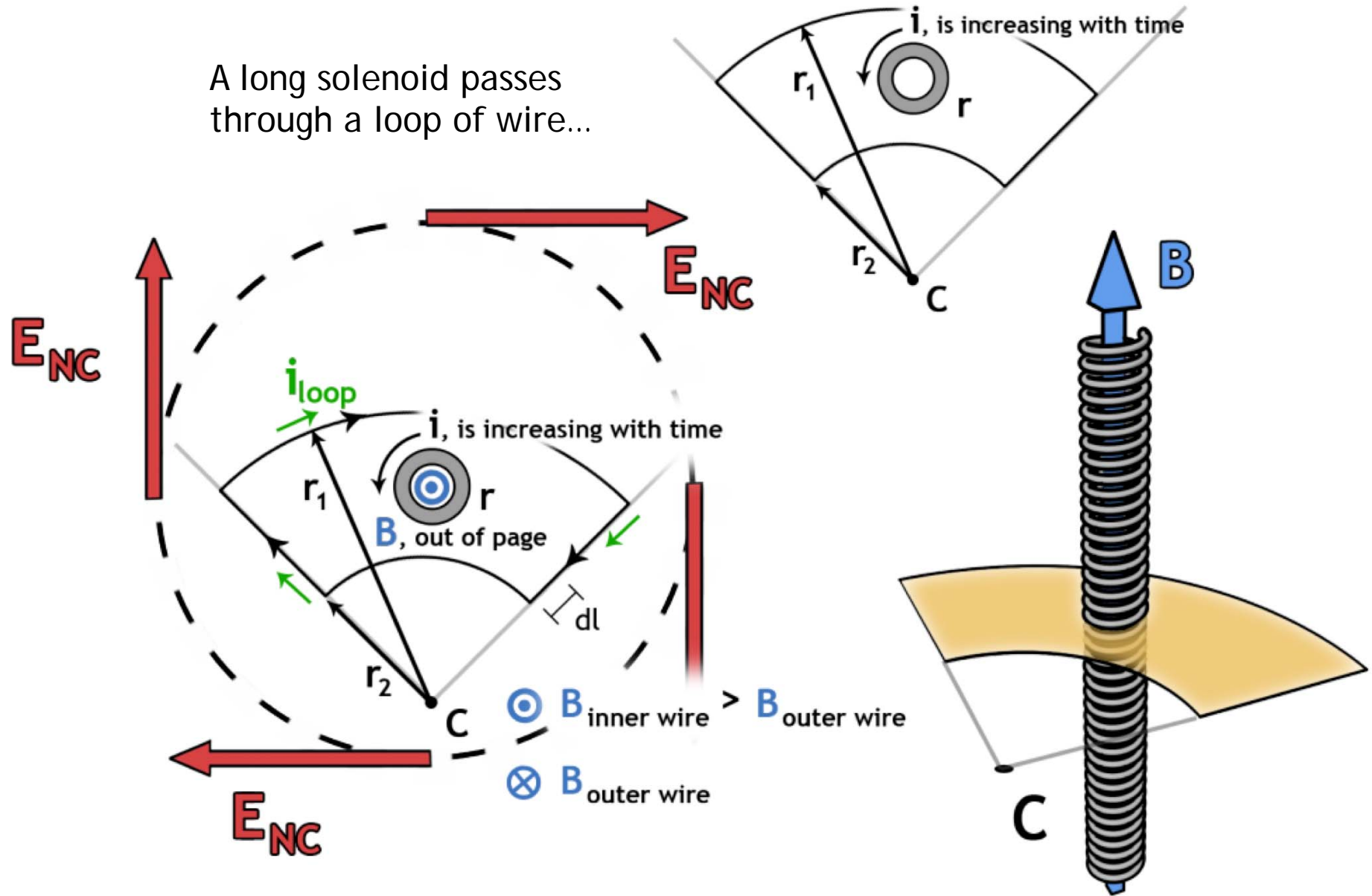


Rotating a bar of magnet (or the coil) produces a time-varying magnetic field inside the coil



Will the current run CLOCKWISE or ANTICLOCKWISE ?

A long solenoid passes through a loop of wire...



Electric Fields

$$\oint_S \epsilon_0 \bar{E} \cdot d\bar{A} = \int_V \rho dV$$
$$= Q_{\text{enclosed}}$$

GAUSS

FARADAY

$$\oint_C \bar{E} \cdot d\bar{l} = -\frac{d}{dt} \left(\int_S \bar{B} \cdot d\bar{A} \right)$$



$$emf = \frac{d\Phi}{dt}$$

Magnetic Fields

$$\oint_S \bar{B} \cdot d\bar{A} = 0$$

GAUSS

AMPERE

$$\oint_C \bar{H} \cdot d\bar{l}$$
$$= \int_S \bar{J} \cdot d\bar{A} + \frac{d}{dt} \int_S \epsilon E dA$$

Next ...

- *MAGNETIC MATERIALS*
- *MAGNETIC CIRCUITS*

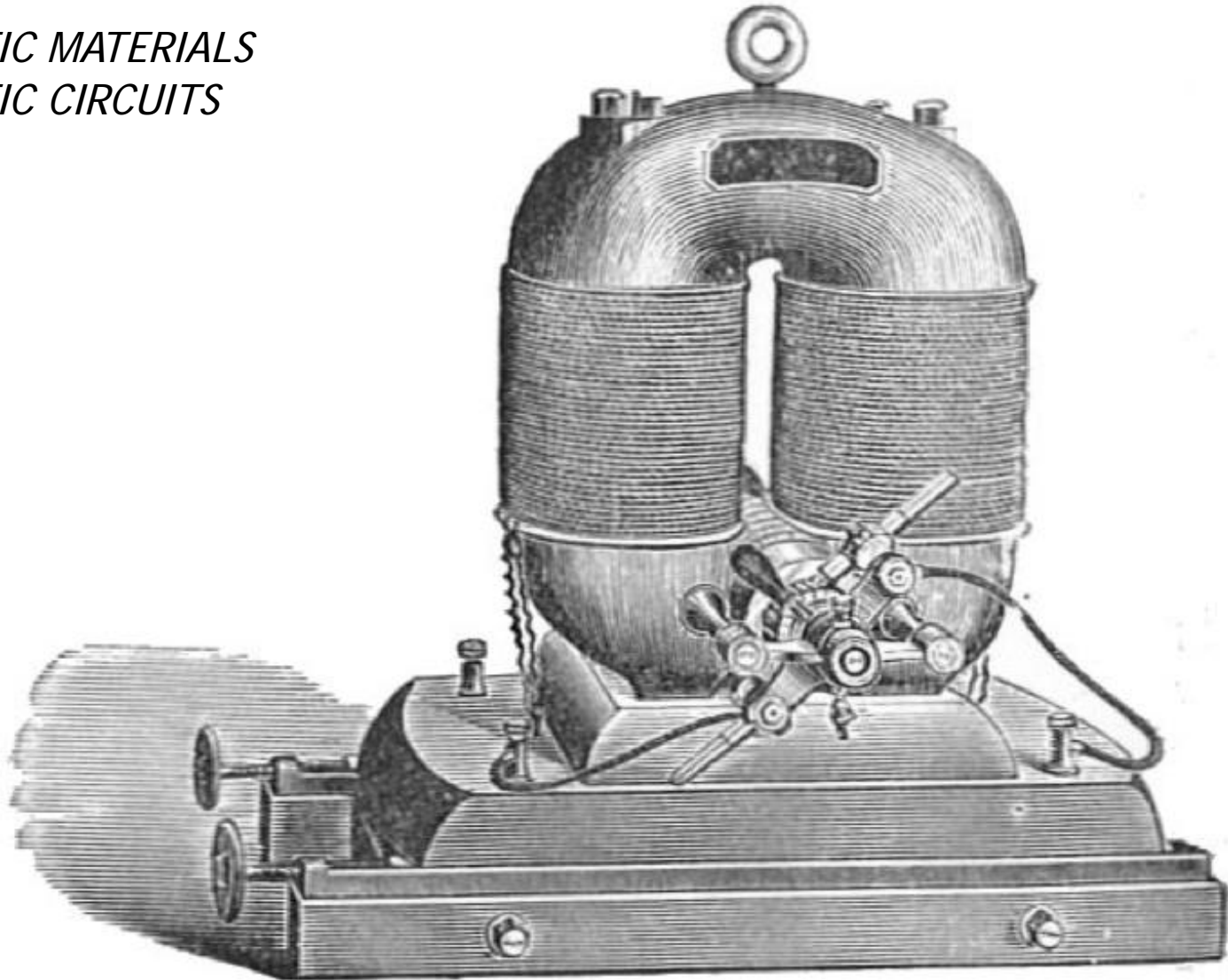


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KEY TAKEAWAYS

FOR A SUFFICIENTLY LONG SOLENOID...

INDUCTANCE:

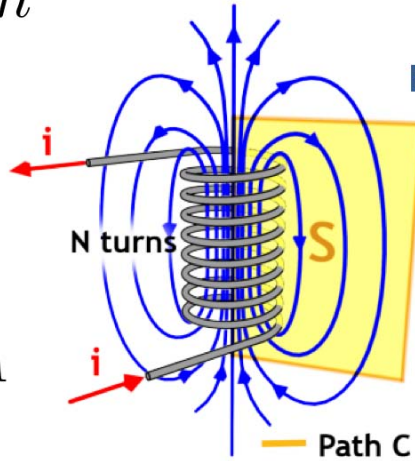
UNITS of INDUCTANCE:

$$\lambda = N\Phi = \mu_o \frac{N^2 i}{h} A = Li \Rightarrow L = \mu_o \frac{A}{h} N^2 \quad \text{H} = \frac{\text{Wb}}{\text{A}} = \frac{\text{V} \cdot \text{s}}{\text{A}}$$

$$H_{\text{inside}} = \frac{Ni}{h}$$

$$B_{\text{inside}} = \mu_o \frac{Ni}{h}$$

$$\Phi_{\text{inside}} = \mu_o \frac{Ni}{h} A$$



ENERGY STORED in an INDUCTOR:

$$W_s = \frac{1}{2} Li^2$$

$$\frac{W_s}{\text{Volume}} = \frac{1}{2} \mu_o H^2$$

THERE ARE TWO WAYS TO PRODUCE ELECTRIC FIELD

- (1) Coulomb electric field is produced by electric charges according to Coulomb's law
- (2) Non-Coulomb electric field E_{NC} is due to time-varying magnetic flux density dB/dt

Faraday's Law: The induced emf along a round-trip path is equal to the rate of change of the magnetic flux on the area encircled by the path.

Lenz's Rule:

The induced electric field would drive the current in the direction to make the magnetic field that attempts to keep the flux constant

$$\oint_C \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \left(\int_S \vec{B} \cdot d\vec{A} \right)$$

$\text{emf} = -\frac{d\Phi_{\text{mag}}}{dt}$

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