#### 6.012 - Microelectronic Devices and Circuits

# Lecture 11 - MOSFETs II; Large Signal Models - Outline

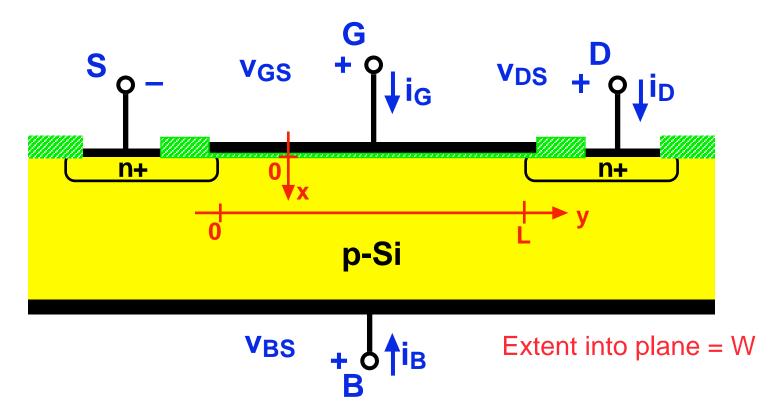
- Announcements On Stellar - 2 write-ups on MOSFET models
- The Gradual Channel Approximation (review and more) MOSFET model: gradual channel approximation (Example: n-MOS)  $i_D \approx \begin{cases} 0 & \text{for } (v_{GS} - V_T)/\alpha \leq 0 \leq v_{DS} & \text{(cutoff)} \\ K (v_{GS} - V_T)^2/2\alpha & \text{for } 0 \leq (v_{GS} - V_T)/\alpha \leq v_{DS} & \text{(saturation)} \\ K (v_{GS} - V_T - \alpha v_{DS}/2)v_{DS} & \text{for } 0 \leq v_{DS} \leq (v_{GS} - V_T)/\alpha & \text{(linear)} \end{cases}$ with  $K \equiv (W/L)\mu_e C_{ox}^{*}$ ,  $V_T = V_{FB} - 2\phi_{p-Si} + [2\epsilon_{Si} qN_A(|2\phi_{p-Si}| - v_{BS})]^{1/2}/C_{ox}^{*}$ and  $\alpha = 1 + [(\epsilon_{Si} qN_A/2(|2\phi_{p-Si}| - v_{BS})]^{1/2}/C_{ox} & \text{(frequently } \alpha \approx 1)$
- Refined device models for transistors (MOS and BJT) Other flavors of MOSFETS: p-channel, depletion mode The Early Effect:
  - 1. Base-width modulation in BJTs:  $w_B(v_{CE})$
  - 2. Channel-length modulation in MOSFETs:  $L(v_{DS})$

#### **Charge stores:**

- 1. Junction diodes
- 2. BJTs
- 3. MOSFETs

**Extrinsic parasitics:** Lead resistances, capacitances, and inductances

# An n-channel MOSFET showing gradual channel axes



### **Gradual Channel Approximation:**

- A <u>one-dimensional electrostatics problem</u> in the x direction is solved to find the channel charge, q<sub>N</sub><sup>\*</sup>(y); this charge depends on v<sub>GS</sub>, v<sub>CS</sub>(y) and v<sub>BS</sub>.
- A <u>one-dimensional drift problem</u> in the y direction then gives the channel current,  $i_D$ , as a function of  $v_{GS}$ ,  $v_{DS}$ , and  $v_{BS}$ .

# Gradual Channel Approximation i-v Modeling

(n-channel MOS used as the example)

The Gradual Channel Approximation is the approach typically used to model the drain current in field effect transistors.\*

It assumes that the drain current,  $i_D$ , consists entirely of carriers flowing in the channel of the device, and is thus proportional to the sheet density of carriers at any point and their net average velocity. It is not a function of y, but its components in general are:

$$i_D = -W \cdot -q \cdot n_{ch}^*(y) \cdot \bar{s}_{ey}(y)$$

In this expression, W is the width of the device, -q is the charge on each electron,  $n_{ch}^*(y)$  is sheet electron concentration in the channel (i.e. electrons/cm<sup>2</sup>) at y, and  $\overline{s}_{ey}(y)$  is the net electron velocity in the y-direction.

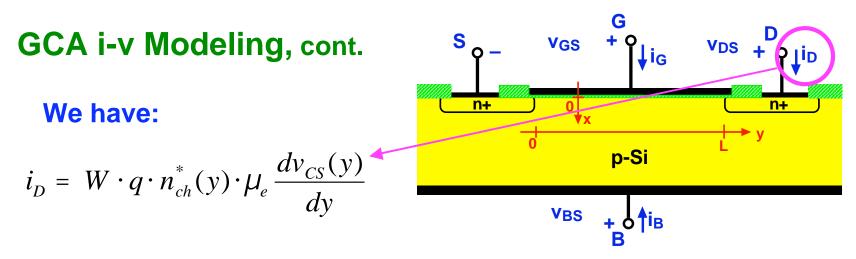
If the electric field is not too large,  $\overline{s}_{ey}(y) = -\mu_e E_y(y)$ , and

$$i_D = -W \cdot q \cdot n_{ch}^*(y) \cdot \mu_e E_y(y) = W \cdot q \cdot n_{ch}^*(y) \cdot \mu_e \frac{dv_{CS}(y)}{dy}$$

\* Junction FETs (JFETs), MEtal Semiconductor FETs (MESFETs<sup>1</sup>), and Heterojunction FETs Clif Fonstad, 3/18/08 (HJFETs<sup>2</sup>), as well as Metal Oxide Semiconductor FETs (MOSFETs). Lecture 11 - Slide 3

Cont.

1. Also called Shottky Barrrier FETs (SBFETs). 2. Includes HEMTs, TEGFETs, MODFETs, SDFETs, HFETs, PHEMTs, MHEMTs, etc.



To eliminate the derivative from this equation we integrate both sides with respect to y from the source (y = 0) to the drain (y = L). This corresponds to integrating the right hand side with respect to  $v_{CS}$  from 0 to  $v_{DS}$ , because  $v_{CS}(0) = 0$  to  $v_{CS}(L) = v_{DS}$ :

$$\int_{0}^{L} i_{D} \, dy = W \cdot \mu_{e} \cdot q \cdot \int_{0}^{L} n_{ch}^{*}(y) \frac{dv_{CS}(y)}{dy} \, dy = W \cdot \mu_{e} \cdot q \cdot \int_{0}^{v_{DS}} n_{ch}^{*}(v_{CS}) \, dv_{CS}$$

The left hand integral is easy to evaluate; it is simply  $i_D L$ . Thus we have:

$$\int_{0}^{L} i_{D} dy = i_{D}L \quad \Rightarrow \left( i_{D} = \frac{W}{L} \cdot \mu_{e} \cdot q \cdot \int_{0}^{v_{DS}} n_{ch}^{*}(v_{CS}) dv_{CS} \right)$$

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Cont.

### GCA i-v Modeling, cont.

- The various FETs differ primarily in the nature of their channels and thereby, the expressions for  $n_{ch}^*(y)$ .
- For a <u>MOSFET</u> we speak in terms of the inversion layer charge,  $q_n^*(y)$ , which is equivalent to  $-q \cdot n_{ch}^*(y)$ . Thus we have:

$$i_{D} = -\frac{W}{L} \mu_{e} \int_{0}^{v_{DS}} q_{n}^{*}(v_{GS}, v_{CS}, v_{BS}) dv_{CS}$$

We derived  $q_n^*$  earlier by solving the vertical electrostatics problem, and found:

$$q_{n}^{*}(v_{GS}, v_{CS}, v_{BS}) = -C_{ox}^{*} \left[ v_{GS} - v_{CS} - V_{T}(v_{CS}, v_{BS}) \right]$$
  
with  $V_{T}(v_{CS}, v_{BS}) = V_{FB} - 2\phi_{p-Si} + \left\{ 2\varepsilon_{Si}qN_{A} \left[ \left| 2\phi_{p-Si} \right| - v_{BS} + v_{CS} \right] \right\}^{1/2} / C_{ox}^{*}$ 

Using this in the equation for  $i_D$ , we obtain:

$$i_D(v_{GS}, v_{DS}, v_{BS}) = \frac{W}{L} \mu_e \int_0^{v_{DS}} \left\{ C_{ox}^* \left[ v_{GS} - v_{CS} - V_T (v_{CS}, v_{BS}) \right] \right\} dv_{CS}$$

At this point we can do the integral, but it is common to simplify the expression of  $V_T(v_{CS}, v_{BS})$  first, to get a more useful result.

### Approach #1 - Live with it

Even though  $V_T(v_{CS}, v_{BS})$  is a non-linear function of  $v_{CS}$ , we can still put it in this last equation for  $i_D$ :

$$i_{D} = \frac{W}{L} \mu_{e} \int_{0}^{v_{DS}} \left\{ C_{ox}^{*} \left[ v_{GS} - v_{CS} - V_{FB} + 2\phi_{p-Si} - \frac{t_{ox}}{\varepsilon_{ox}} \sqrt{2\varepsilon_{Si}qN_{A} \left[ \left| 2\phi_{p-Si} \right| - v_{BS} + v_{CS} \right]} \right] \right\} dv_{CS}$$

and do the integral, obtaining:

$$i_{D}(v_{DS}, v_{GS}, v_{BS}) = \frac{W}{L} \mu_{e} C_{ox}^{*} \left\{ \left( v_{GS} - \left| 2\phi_{p} \right| - V_{FB} - \frac{v_{DS}}{2} \right) v_{DS} + \frac{3}{2} \sqrt{2\varepsilon_{Si} q N_{A}} \left[ \left( \left| 2\phi_{p} \right| + v_{DS} - v_{BS} \right)^{3/2} - \left( \left| 2\phi_{p} \right| - v_{BS} \right)^{3/2} \right] \right\}$$

The problem is that this result is very unwieldy, and difficult to work with. More to the point, we don't have to live with it because it is easy to get very good, approximate solutions that are much simpler to work with.

# <u>Approach #2</u> - Ignore it

Early on researchers noticed that the difference between  $V_T$  at 0 and at y, i.e.  $V_T(0, v_{BS})$  and  $V_T(v_{DS}, v_{BS})$ , is small, and that using  $V_T(0, v_{BS})$  alone gives a result that is still quite accurate and is very easy to use:

$$i_{D}(v_{GS}, v_{DS}, v_{BS}) = \frac{W}{L} \mu_{e} \int_{0}^{v_{DS}} \left\{ C_{ox}^{*} \left[ v_{GS} - v_{CS} - V_{T}(0, v_{BS}) \right] \right\} dv_{CS}$$

$$= \frac{W}{L} \mu_{e} C_{ox}^{*} \left\{ \left[ v_{GS} - V_{T}(0, v_{BS}) \right] v_{DS} - \frac{v_{DS}^{2}}{2} \right\}$$
The variable,  $v_{CS}$ , is set to 0 in  $V_{T}$ .

This result looks much simpler than the result of Approach #1, and it is much easier to use in hand calculations. It is, in fact, the one most commonly used by the vast majority of engineers. At the same time, the fact that it was obtained by ignoring the dependence of  $V_T$  on  $v_{CS}$  is cause for concern, unless we have a way to judge the validity of our approximation. We can get the necessary metric through Approach #3.

Approach #3 - Linearize it (i.e. expand it, keep first order term) In this approach we leave the variation of  $V_T$  with  $v_{CS}$  in, but linearize it by doing a Taylor's series expansion about  $v_{CS} = 0$ :

$$V_T[v_{CS}, v_{BS}] \approx V_T(0, v_{BS}) + \left. \frac{\partial V_T}{\partial v_{CS}} \right|_{v_{CS}=0} \cdot v_{CS}$$

Taking the derivative and evaluating it at  $v_{cs} = 0$  yields:

$$V_{T}[v_{CS}, v_{BS}] \approx V_{T}(0, v_{BS}) + \frac{t_{ox}}{\varepsilon_{ox}} \sqrt{\frac{\varepsilon_{Si}qN_{A}}{2(|2\phi_{p}| - v_{BS})}} \cdot v_{CS}$$

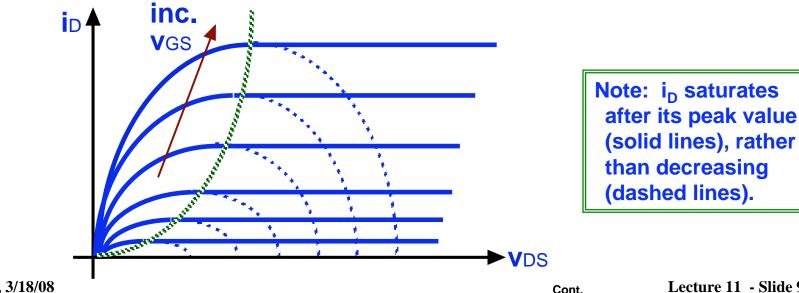
With this  $q_n^*$  is  $q_n^*(v_{GS}, v_{CS}, v_{BS}) \approx -C_{ox}^* \left[ v_{GS} - v_{CS} + V_T(0, v_{BS}) - \frac{t_{ox}}{\varepsilon_{ox}} \sqrt{\frac{\varepsilon_{Si} q N_A}{2(|2\phi_p| - v_{BS})}} \cdot v_{CS} \right]$  $= -C_{\alpha x}^{*} [v_{GS} - \alpha v_{CS} + V_{T} (v_{BS})]$ 

where  $\alpha = \left[1 + \frac{t_{ox}}{\varepsilon_{ox}} \sqrt{\varepsilon_{Si} q N_A / 2 \left( \left| 2\phi_p \right| - v_{BS} \right)} \right] \text{ and } V_T(v_{BS}) = V_T(0, v_{BS})$ Lecture 11 -Lecture 11 - Slide 8

Using this result in the integral in the expression for  $i_D$  gives:

 $i_{D}(v_{GS}, v_{DS}, v_{BS}) = \frac{W}{L} \mu_{e} \int_{0}^{v_{DS}} \left\{ C_{ox}^{*} \left[ v_{GS} - \alpha v_{CS} - V_{T}(0, v_{BS}) \right] \right\} dv_{CS}$  $= \frac{W}{L} \mu_e C_{ox}^* \left\{ \left[ v_{GS} - V_T(v_{BS}) \right] v_{DS} - \alpha \frac{v_{DS}^2}{2} \right\}$ Except for  $\alpha$  this is the Approach 2 result.

Plotting this equation for increasing values of  $v_{GS}$  we see that it traces inverted parabolas as shown below.



**Gradual Channel Approximation, cont.** 

The drain current expression, cont:

The point at which i<sub>D</sub> reaches its peak value and saturates is easily found. Taking the derivative and setting it equal to zero we find:  $\frac{\partial i_D}{\partial v_{DS}} = 0$  when  $v_{DS} = \frac{1}{\alpha} [v_{GS} - V_T (v_{BS})]$ 

What happens physically at this voltage is that the channel (inversion) at the drain end of the channel disappears:

$$q_{n}^{*}(L) \approx -C_{ox}^{*} \{ v_{GS} - V_{T}(v_{BS}) - \alpha v_{DS} \}$$
  
= 0 when  $v_{DS} = \frac{1}{\alpha} [v_{GS} - V_{T}(v_{BS})]$ 

For  $v_{DS} > [v_{GS}-V_T(v_{BS})]/\alpha$ , all the additional drain-to-source voltage appears across the high resistance region at the drain end of the channel where the mobile charge density is very small, and i<sub>D</sub> remains constant independent of v<sub>DS</sub>:

$$i_{D}(v_{GS}, v_{DS}, v_{BS}) = \frac{1}{2\alpha} \frac{W}{L} \mu_{e} C_{ox}^{*} [v_{GS} - V_{T}(v_{BS})]^{2} \text{ for } v_{DS} > \frac{1}{\alpha} [v_{GS} - V_{T}(v_{BS})]$$

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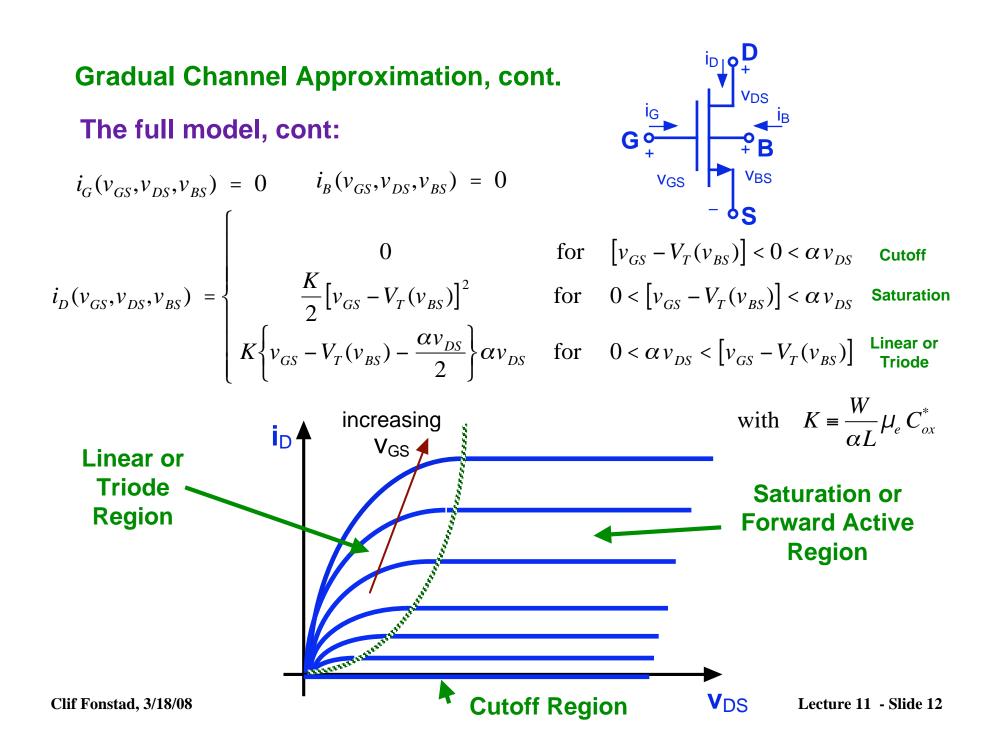
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#### **Gradual Channel Approximation, cont.**

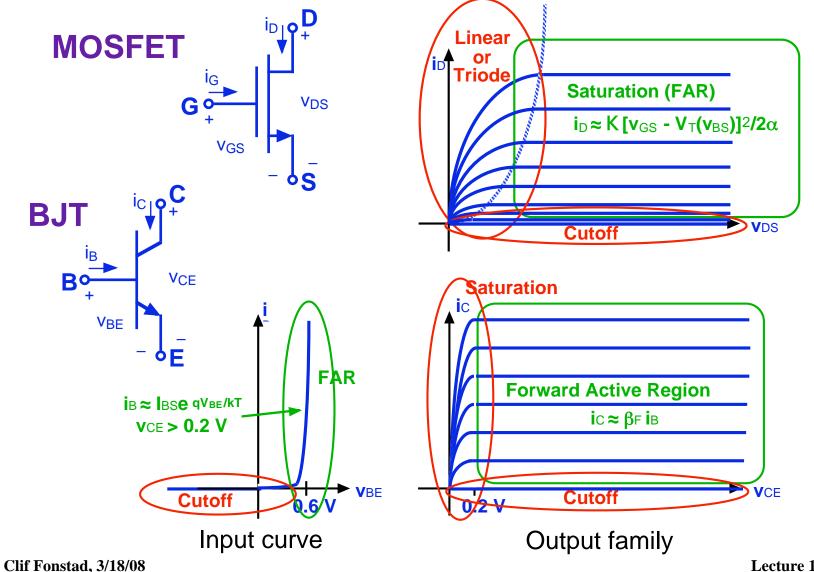
The full model:

### With this drain current expression, we now have the complete set of Gradual Channel Model expressions for the MOSFET terminal characteristics in the three regions of operation:

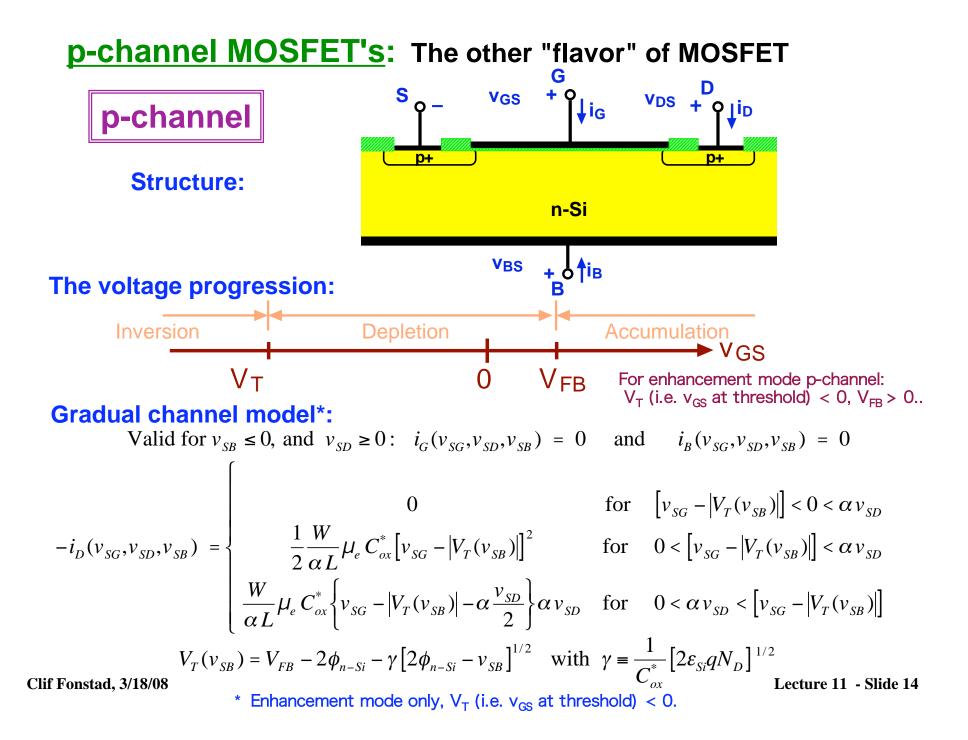
$$\begin{aligned} \text{Valid for } v_{BS} &\leq 0, \text{ and } v_{DS} \geq 0: \\ i_{G}(v_{GS}, v_{DS}, v_{BS}) &= 0 \quad \text{and} \quad i_{B}(v_{GS}, v_{DS}, v_{BS}) = 0 \\ i_{D}(v_{GS}, v_{DS}, v_{BS}) &= \begin{cases} 0 & \text{for } \left[v_{GS} - V_{T}(v_{BS})\right] < 0 < \alpha v_{DS} \\ \frac{1}{2} \frac{W}{\alpha L} \mu_{e} C_{\alpha x}^{*} \left[v_{GS} - V_{T}(v_{BS})\right]^{2} & \text{for } 0 < \left[v_{GS} - V_{T}(v_{BS})\right] < \alpha v_{DS} \\ \frac{W}{\alpha L} \mu_{e} C_{\alpha x}^{*} \left\{v_{GS} - V_{T}(v_{BS}) - \alpha \frac{v_{DS}}{2}\right\} \alpha v_{DS} & \text{for } 0 < \alpha v_{DS} < \left[v_{GS} - V_{T}(v_{BS})\right] \\ \text{with } V_{T}(v_{BS}) &= V_{FB} - 2\phi_{p-Si} + \frac{1}{C_{\alpha x}^{*}} \left\{2\varepsilon_{Si}qN_{A} \left[\left|2\phi_{p-Si}\right| - v_{BS}\right]\right\}^{1/2} \\ \alpha &\equiv 1 + \frac{1}{C_{\alpha x}^{*}} \left\{\frac{\varepsilon_{Si}qN_{A}}{2\left[\left|2\phi_{p-Si}\right| - v_{BS}\right]}\right\}^{1/2} & C_{\alpha x}^{*} &\equiv \frac{\varepsilon_{\alpha x}}{t_{\alpha x}} \\ \text{Leture 11 - Slide 11} \end{aligned}$$

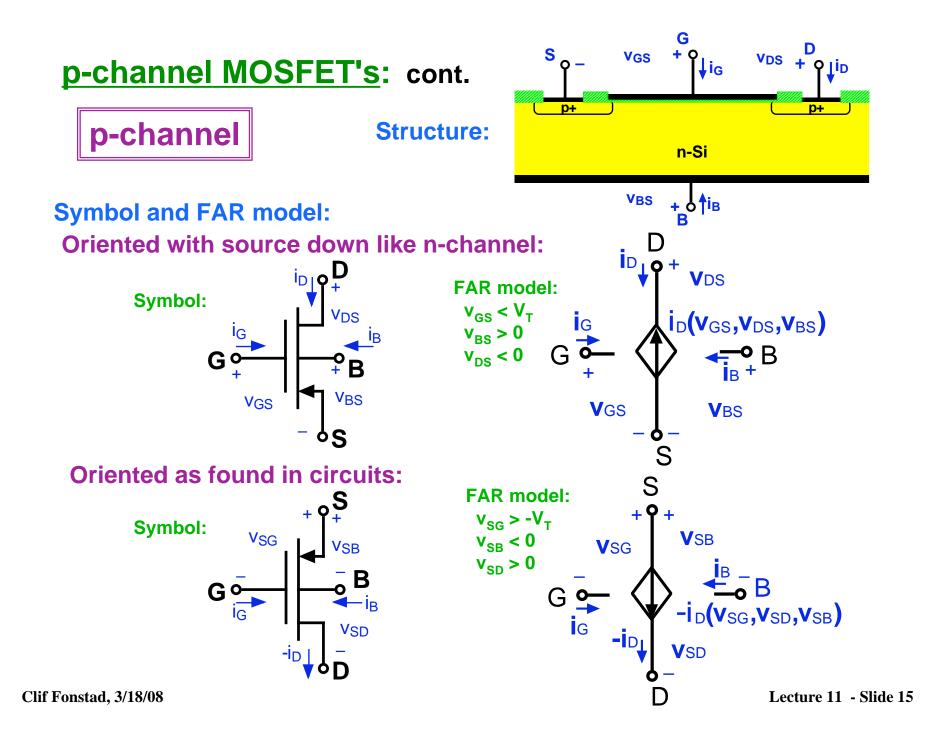


# The operating regions of MOSFETs and BJTs: Comparing an n-channel MOSFET and an npn BJT



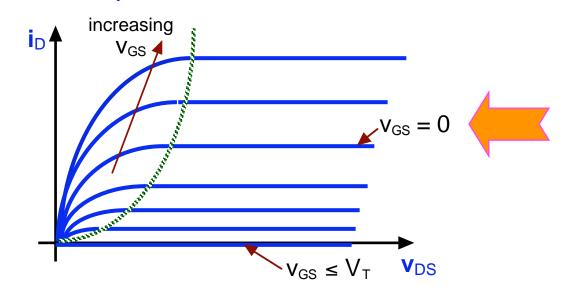
Lecture 11 - Slide 13





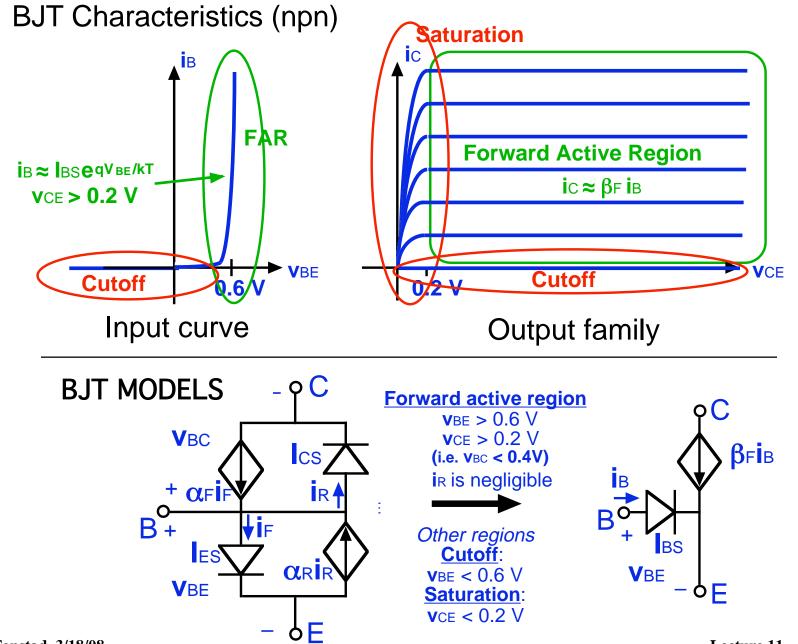
## **Depletion mode MOSFET's:** The very last MOSFET variant

- It is possible to have n-channel MOSFETs with  $V_T < 0$ .
- In this situation the channel exists with  $v_{GS} = 0$ , and a negative bias must be applied to turn it off.
- This type of device is called a "depletion mode" MOSFET. Devices with  $V_{\tau} > 0$  are "enhancement mode."



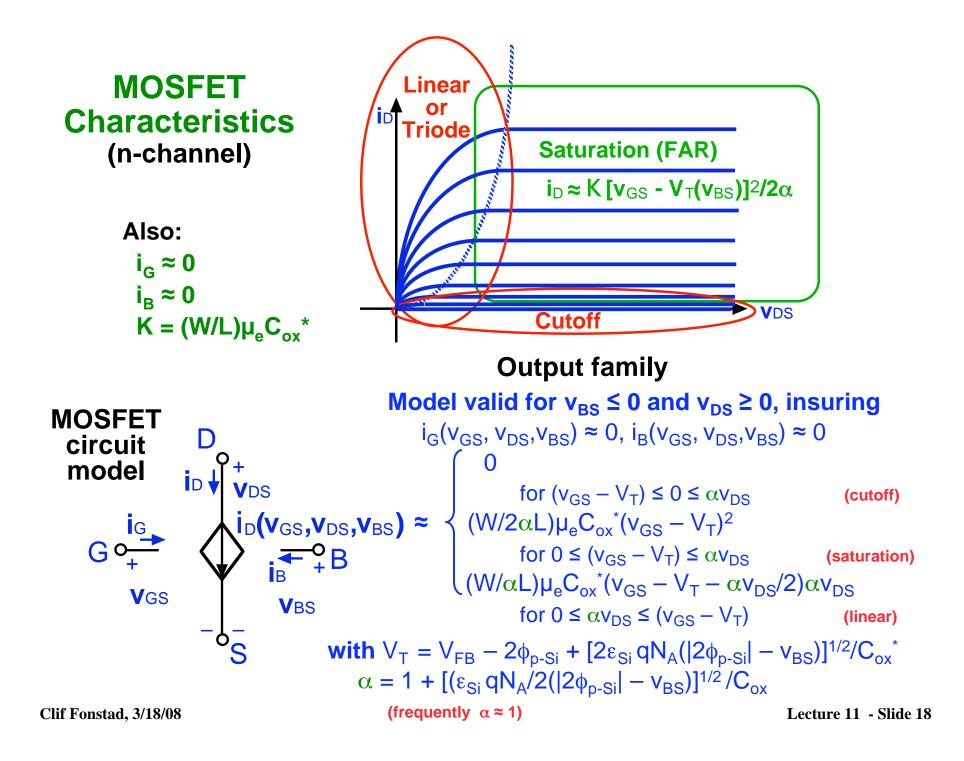
For a p-channel depletion mode MOSFET,  $V_T > 0$ .

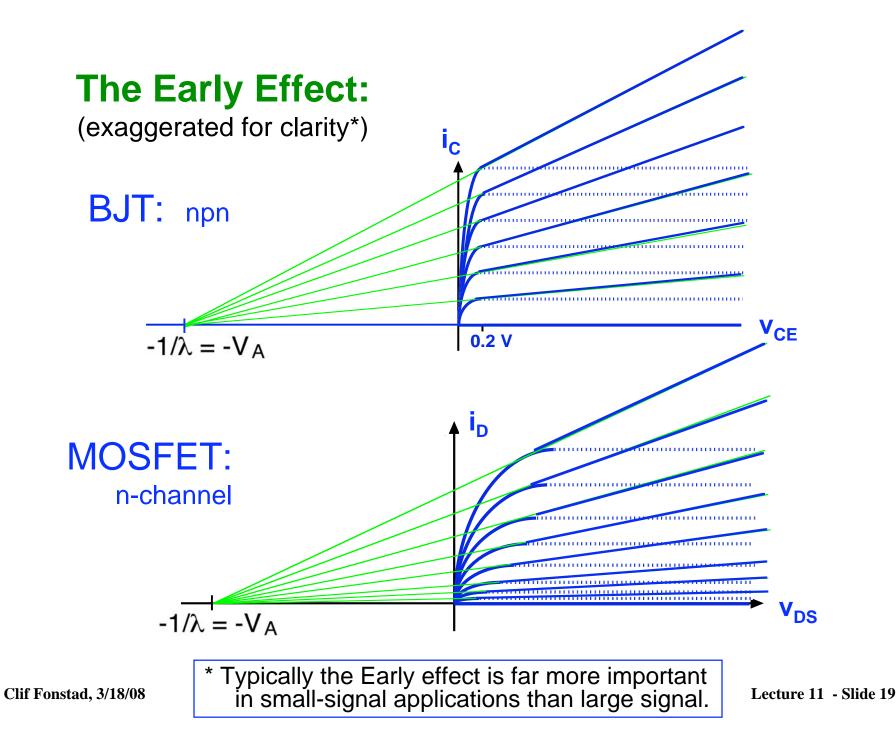
The expressions for  $i_D(v_{GS}, v_{DS}, v_{BS})$  are <u>exactly the same</u> for enhancement mode and depletion mode MOSFETs.



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#### Active Length Modulation - the Early Effect: MOSFET "Channel length modulation"

# **MOSFET**:

We begin by recognizing that the channel length decreases with increasing  $v_{DS}$  and writing this dependence to first order in  $v_{DS}$ :

$$L \approx L_o \left[ 1 - \lambda \left( v_{DS} - V_{DSat} \right) \right] \quad \text{and} \quad \frac{1}{L} \approx \frac{\left[ 1 + \lambda \left( v_{DS} - V_{DSat} \right) \right]}{L_o}$$
$$K = \frac{W}{\alpha L} \mu_e C_{ox}^*$$

**Inserting the channel length variation with v<sub>DS</sub> into K we have:** 

$$K \approx K_o [1 + \lambda (v_{DS} - V_{DSat})]$$
 where  $K_o \equiv \frac{W}{\alpha L_o} \mu_e C_{ox}^*$ 

Thus, in saturation:

$$i_D \approx \frac{K_o}{2} \left( v_{GS} - V_T \right)^2 \left[ 1 + \lambda \left( v_{DS} - V_{DSat} \right) \right]$$

**Note:**  $\lambda$  is the inverse of the Early Voltage, V<sub>A</sub> (i.e.,  $\lambda = 1/V_A$ ).

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#### Active Length Modulation - the Early Effect: BJT "Base width modulation" BJT:

We begin by recognizing that the base width decreases with increasing  $v_{CE}$  and writing this dependence to first order in  $v_{CE}$ :  $w_B^* \approx w_{Bo}^* (1 - \lambda v_{CE})$  and  $\frac{1}{w_B^*} \approx \frac{1}{w_{Bo}^*} (1 + \lambda v_{CE})$ 

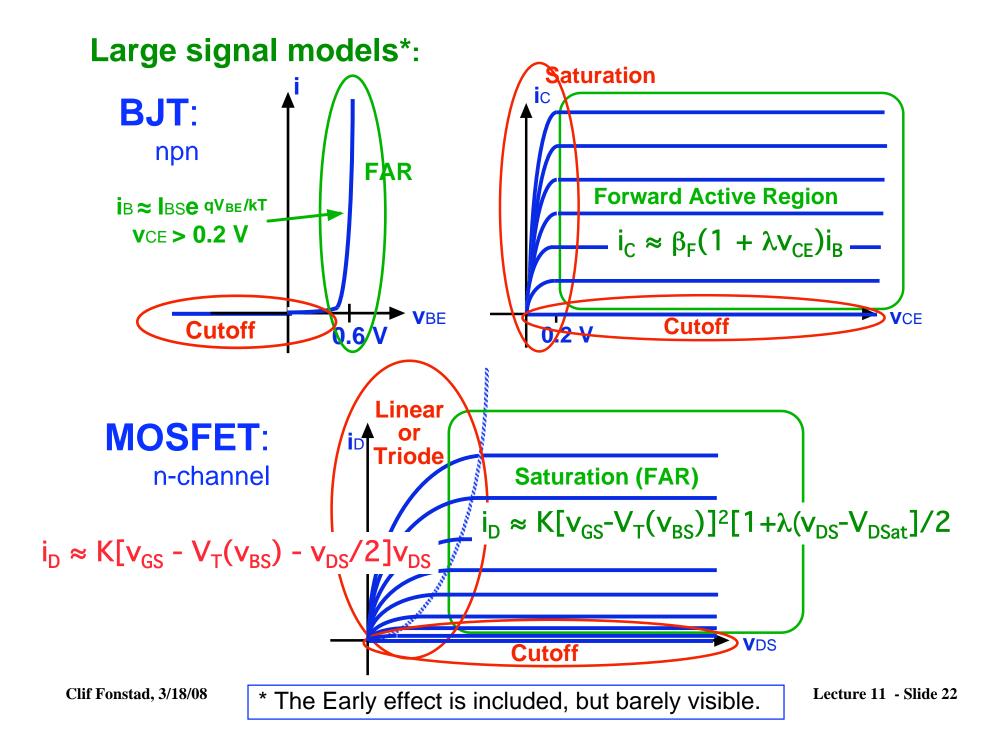
Then we recall that in a modern BJT the base defect,  $\delta_B$ , is negligible and  $\beta_F$  depends primarily on the emitter defect,  $\delta_E$ , and can be written:

$$\beta_F = \frac{\left(1 + \delta_B\right)}{\left(\delta_E + \delta_B\right)} \approx \frac{1}{\delta_E} = \frac{D_e}{D_h} \frac{N_{DE}}{N_{AB}} \frac{w_E^*}{w_B^*}$$

Inserting the base width variation with  $v_{CE}$  into  $\beta_{F}$  we have:

$$\beta_{F} \approx \beta_{Fo} (1 + \lambda v_{CE}) \quad \text{where} \quad \beta_{Fo} \equiv \frac{D_{e}}{D_{h}} \frac{N_{DE}}{N_{AB}} \frac{w_{E}}{w_{Bo}^{*}}$$
  
Thus, in the F.A.R.:  
 $i_{C} \approx \beta_{Fo} (1 + \lambda v_{CE}) i_{B}$ 

Note:  $\lambda$  is the inverse of the Early Voltage,  $V_A$  (i.e.,  $\lambda = 1/V_A$ ). Clif Fonstad, 3/18/08

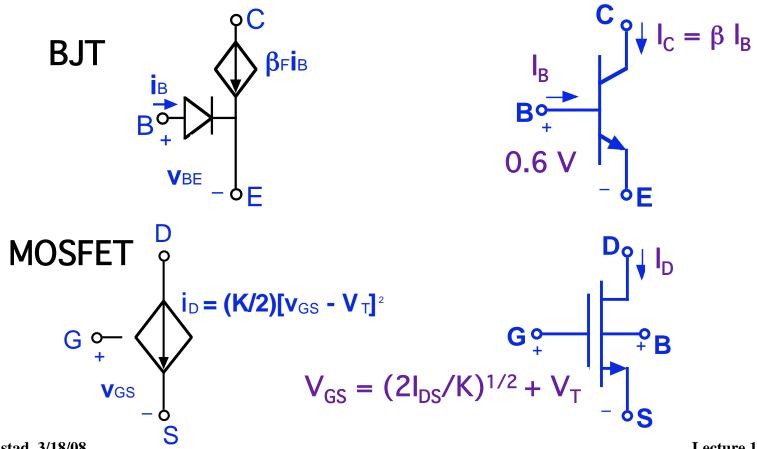


# Large signal models: when will we use them? Digital circuit analysis/design:

This requires use of the entire circuit, and will be the topic of Lectures 14, 15, and 16.

**Bias point analysis/design:** 

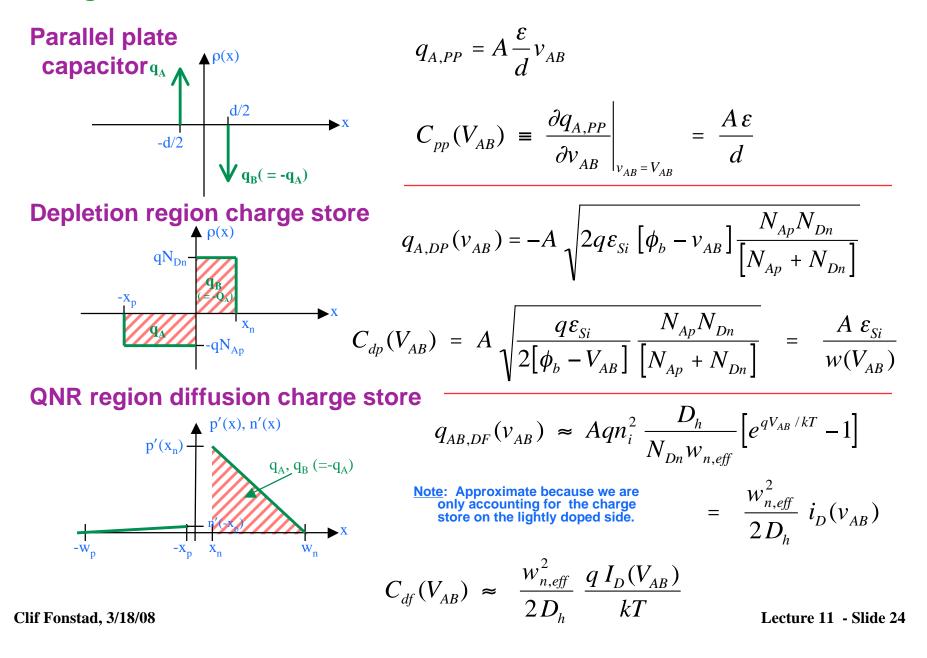
This uses the FAR models (below and Lec. 17ff).



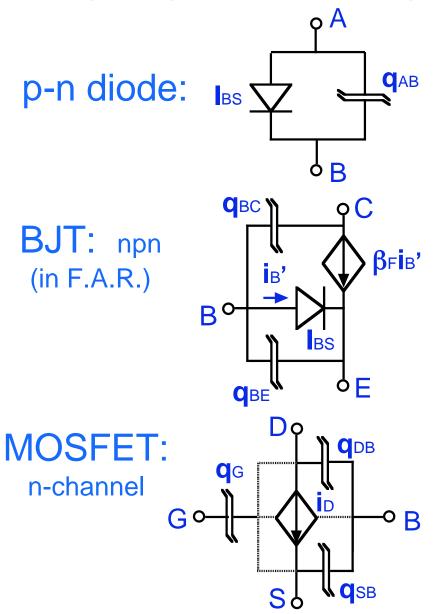
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#### Charge stores in devices: we must add them to our device models



#### Adding charge stores to the large signal models:



q<sub>AB</sub>: Excess carriers on p-side plus excess carriers on n-side plus junction depletion charge.

q<sub>BE</sub>: Excess carriers in base plus E-B junction depletion charge
 q<sub>BC</sub>: C-B junction depletion charge

 $q_{\rm G}$ : Gate charge; a function of  $v_{\rm gs},\,v_{\rm ps},$  and  $v_{\rm BS}.$ 

q<sub>DB</sub>: D-B junction depletion charge

q<sub>sB</sub>: S-B junction depletion charge

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# Lecture 11 - MOSFETS II; Large-Signal Models - Summary

- Gradual channel approximation for FETs **General approach MOSFETS** in strong inversion

  - **1. Ignore variation of V**<sub>T</sub> along channel **2. Linearize variation of V**<sub>T</sub> along channel: introduces  $\alpha$  factor
- Additional device model issues

The Early Effect:

- **1. Base-width modulation in BJTs:**  $W_B(V_{CE})$ In the F.A.R.:  $i_C \approx \beta_{Fo}(1 + Iv_{CF})i_B$
- **2.** Channel-length modulation in MOSFETs:  $L(v_{DS})$ In saturation:  $i_D \approx K_o (v_{GS} V_T)^2 [1 + \lambda (v_{DS} V_{DSat})]/2\alpha$

**Charge stores:** 

**1. Junction diodes - depletion and diffusion charge** 2. BJTs - at EB junction: depletion and diffusion charge at CB junction: depletion charge (focus on FAR) 3. MOSFETs - between B and S, D: depletion charge of n+-p junctions between G and S, D, B: gate charge (the dominant store) in cut-off:  $C_{gs} \approx C_{gd} \approx 0$ ; all is  $C_{gb}$ linear region:  $C_{gs} = C_{gd} = W \perp C_{ox}^*/2$ in saturation region:  $C_{gs} = (2/3) W \perp C_{ox}^*$  $C_{ad}^{gs} = 0$  (only parasitic overlap)

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