

MIT OpenCourseWare  
<http://ocw.mit.edu>

6.013/ESD.013J Electromagnetics and Applications, Fall 2005

Please use the following citation format:

Markus Zahn, Erich Ippen, and David Staelin, *6.013/ESD.013J Electromagnetics and Applications, Fall 2005*. (Massachusetts Institute of Technology: MIT OpenCourseWare). <http://ocw.mit.edu> (accessed MM DD, YYYY). License: Creative Commons Attribution-Noncommercial-Share Alike.

Note: Please use the actual date you accessed this material in your citation.

For more information about citing these materials or our Terms of Use, visit:  
<http://ocw.mit.edu/terms>

6.013, Electromagnetic Fields, Forces, and Motion  
 Prof. Markus Zahn, Sept. 13, 2005  
**Lecture 2: Electromagnetic Field Boundary Conditions**

I. Boundary Conditions

1. Gauss' Continuity Condition

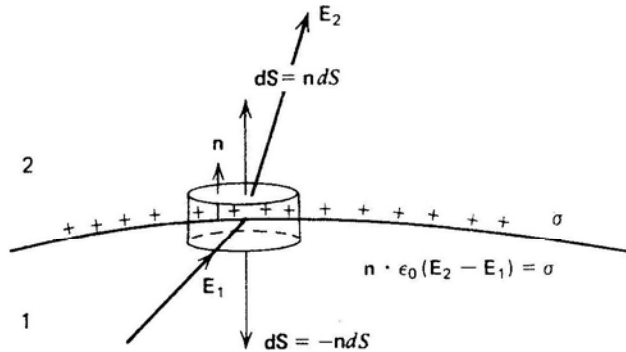


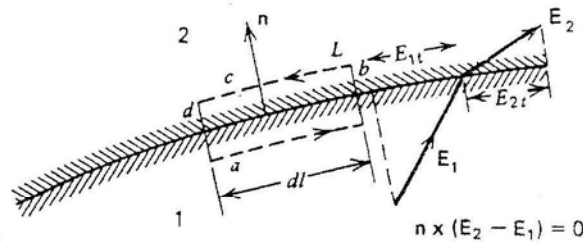
Figure 2-19 Gauss's law applied to a differential sized pill-box surface enclosing some surface charge shows that the normal component of  $\epsilon_0 \mathbf{E}$  is discontinuous in the surface charge density.

Zahn, Markus. Figs. 1.13-1.1.17, 1.19 (a) and (b), 1.23, 1.20, 2.19, 3.12 (a).  
 Electromagnetic Field Theory: A Problem Solving Approach. Robert E. Krieger  
 Publishing Company, Florida, 1987. Used with permission.

$$\oint_S \epsilon_0 \bar{\mathbf{E}} \cdot \bar{d\mathbf{a}} = \int_S \sigma_s dS \Rightarrow \epsilon_0 (E_{2n} - E_{1n}) dS = \sigma_s dS$$

$$\epsilon_0 (E_{2n} - E_{1n}) = \sigma_s \Rightarrow \bar{\mathbf{n}} \cdot [\epsilon_0 (\bar{\mathbf{E}}_2 - \bar{\mathbf{E}}_1)] = \sigma_s$$

2. Continuity of Tangential  $\bar{\mathbf{E}}$



(a)

Figure 3-12 (a) Stokes' law applied to a line integral about an interface of discontinuity shows that the tangential component of electric field is continuous across the boundary.

Zahn, Markus. Figs. 1.13-1.1.17, 1.19 (a) and (b), 1.23, 1.20, 2.19, 3.12 (a).  
 Electromagnetic Field Theory: A Problem Solving Approach. Robert E. Krieger  
 Publishing Company, Florida, 1987. Used with permission.

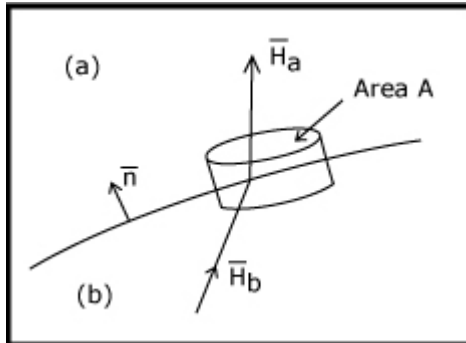
$$\oint_C \bar{\mathbf{E}} \cdot \bar{d\mathbf{s}} = (E_{1t} - E_{2t}) dl = 0 \Rightarrow E_{1t} - E_{2t} = 0$$

$$\bar{\mathbf{n}} \times (\bar{\mathbf{E}}_1 - \bar{\mathbf{E}}_2) = 0$$

Equivalent to  $\Phi_1 = \Phi_2$  along boundary

### 3. Normal $\bar{H}$

$$\oint_S \mu_0 \bar{H} \cdot d\bar{a} = 0$$



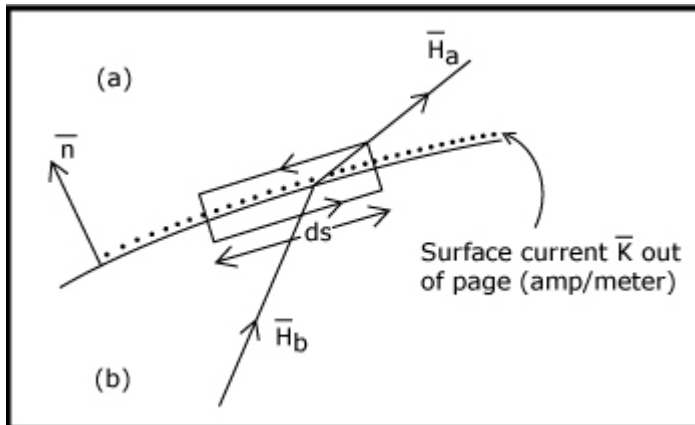
$$\mu_0 (H_{an} - H_{bn}) A = 0$$

$$H_{an} = H_{bn}$$

$$\bar{n} \cdot [\bar{H}_a - \bar{H}_b] = 0$$

### 4. Tangential $\bar{H}$

$$\oint_C \bar{H} \cdot d\bar{s} = \int_S \bar{J} \cdot d\bar{a} + \frac{d}{dt} \int_S \epsilon_0 \bar{E} \cdot d\bar{a}$$

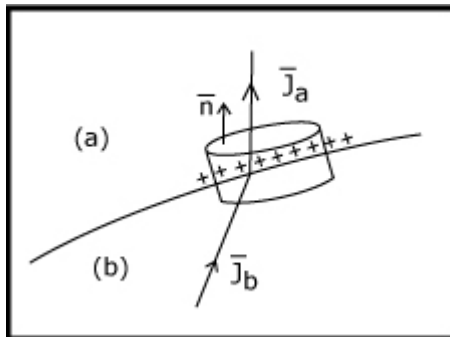


$$H_{bt} ds - H_{at} ds = K ds$$

$$H_{bt} - H_{at} = K$$

$$\bar{n} \times [\bar{H}_a - \bar{H}_b] = \bar{K}$$

### 5. Conservation of Charge Boundary Condition



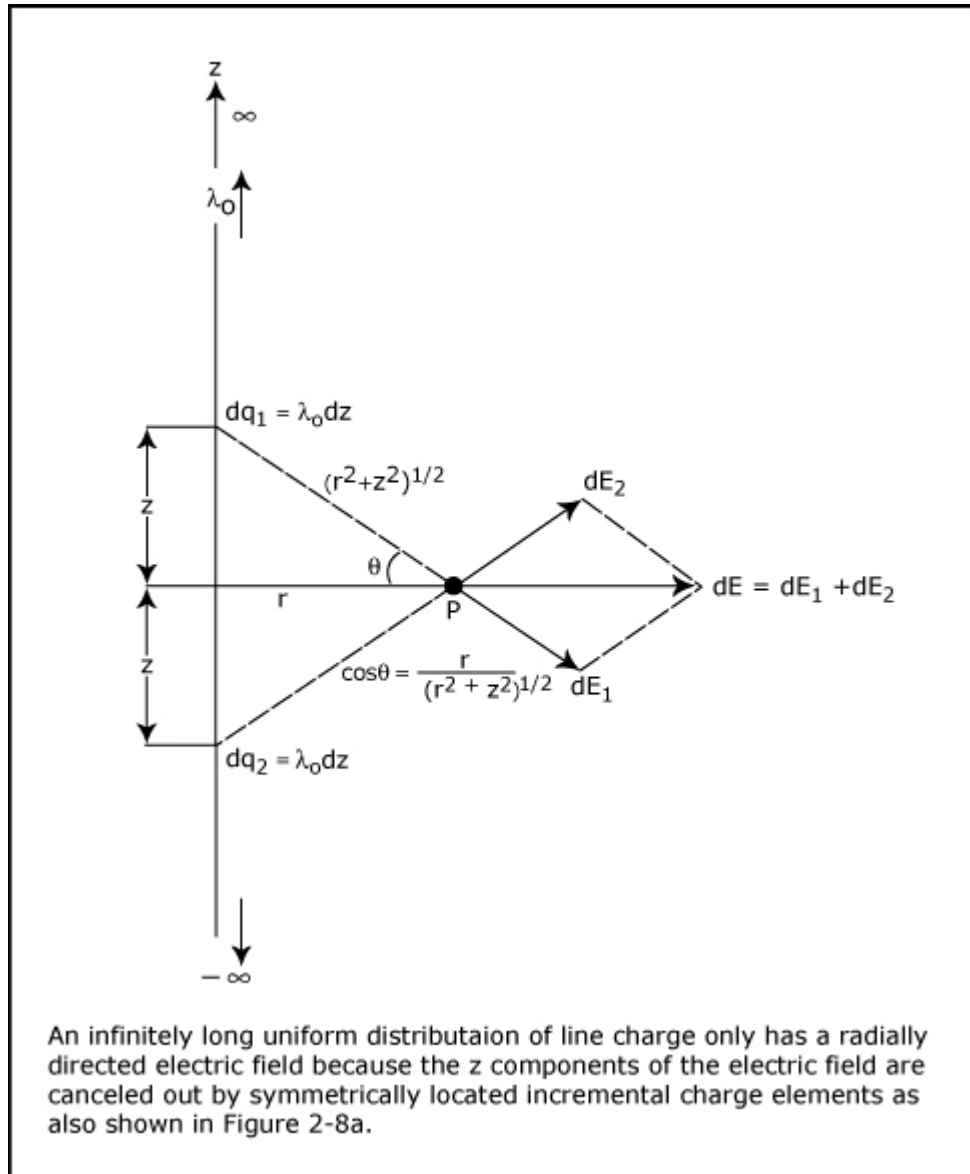
$$\oint_S \bar{J} \cdot d\bar{a} + \frac{d}{dt} \int_V \rho dV = 0$$

$$\bar{n} \cdot [\bar{J}_a - \bar{J}_b] + \frac{\partial}{\partial t} \sigma_s = 0$$

## II. Boundary Condition Problems

### 1. Electric Field from a Sheet of Surface Charge

#### a. Electric Field from a Line Charge



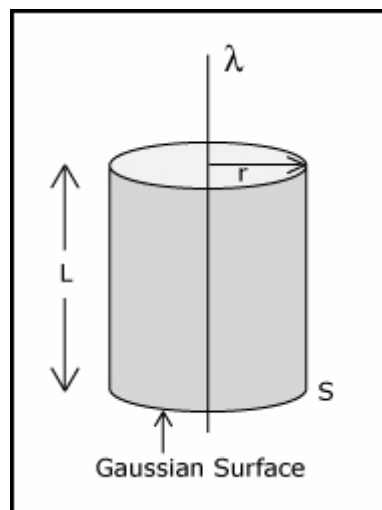
$$dE_r = \frac{dq}{4\pi\epsilon_0(r^2 + z^2)} \cos\theta = \frac{\lambda_0 r dz}{4\pi\epsilon_0(r^2 + z^2)^{3/2}}$$

$$E_r = \int_{z=-\infty}^{+\infty} dE_r = \frac{\lambda_0 r}{4\pi\epsilon_0} \int_{z=-\infty}^{+\infty} \frac{dz}{(r^2 + z^2)^{3/2}}$$

$$= \frac{\lambda_0 r}{4\pi\epsilon_0} \frac{z}{r^2(z^2 + r^2)^{1/2}} \Bigg|_{z=-\infty}^{+\infty}$$

$$= \frac{\lambda_0}{2\pi\epsilon_0 r}$$

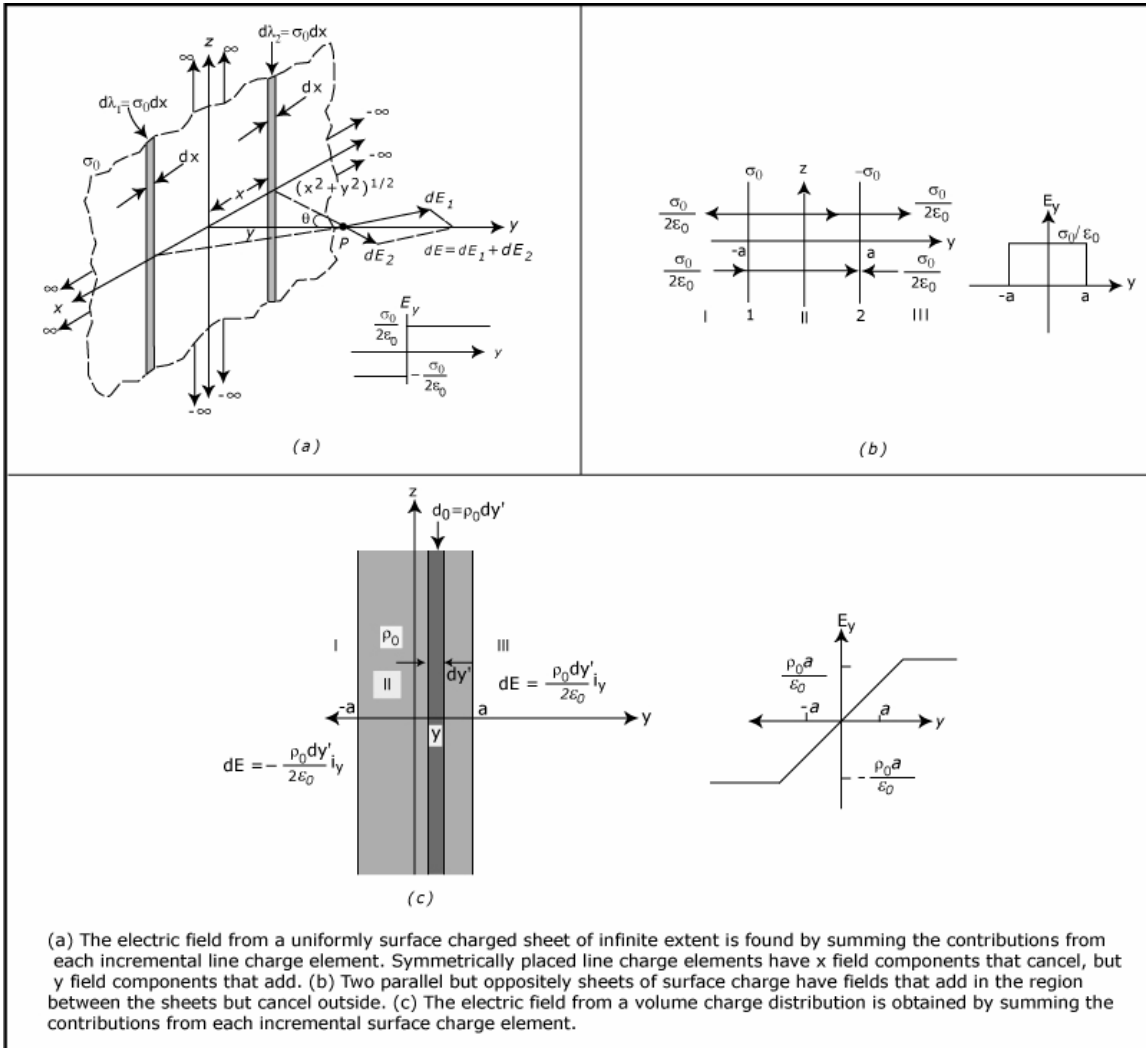
Another way: Gauss' Law



$$\int_S \epsilon_0 \vec{E} \cdot \vec{da} = \epsilon_0 E_r 2\pi r L = \lambda_0 L$$

$$E_r = \frac{\lambda_0}{2\pi\epsilon_0 r}$$

b. Electric Field from a Sheet Charge



$$dE_y = \frac{d\lambda}{2\pi\epsilon_0 (x^2 + y^2)^{3/2}} \cos \theta = \frac{\sigma_0 y dx}{2\pi\epsilon_0 (x^2 + y^2)^{3/2}}$$

$$E_y = \int_{x=-\infty}^{+\infty} dE_y = \frac{\sigma_0 y}{2\pi\epsilon_0} \int_{x=-\infty}^{+\infty} \frac{dx}{x^2 + y^2}$$

$$= \frac{\sigma_0 y}{2\pi\epsilon_0} \frac{1}{y} \tan^{-1} \frac{x}{y} \Big|_{-\infty}^{+\infty}$$

$$= \begin{cases} \frac{\sigma_0}{2\epsilon_0} & y > 0 \\ -\frac{\sigma_0}{2\epsilon_0} & y < 0 \end{cases}$$

Checking Boundary condition at  $y=0$

$$E_y(y = 0_+) - E_y(y = 0_-) = \frac{\sigma_0}{\epsilon_0}$$

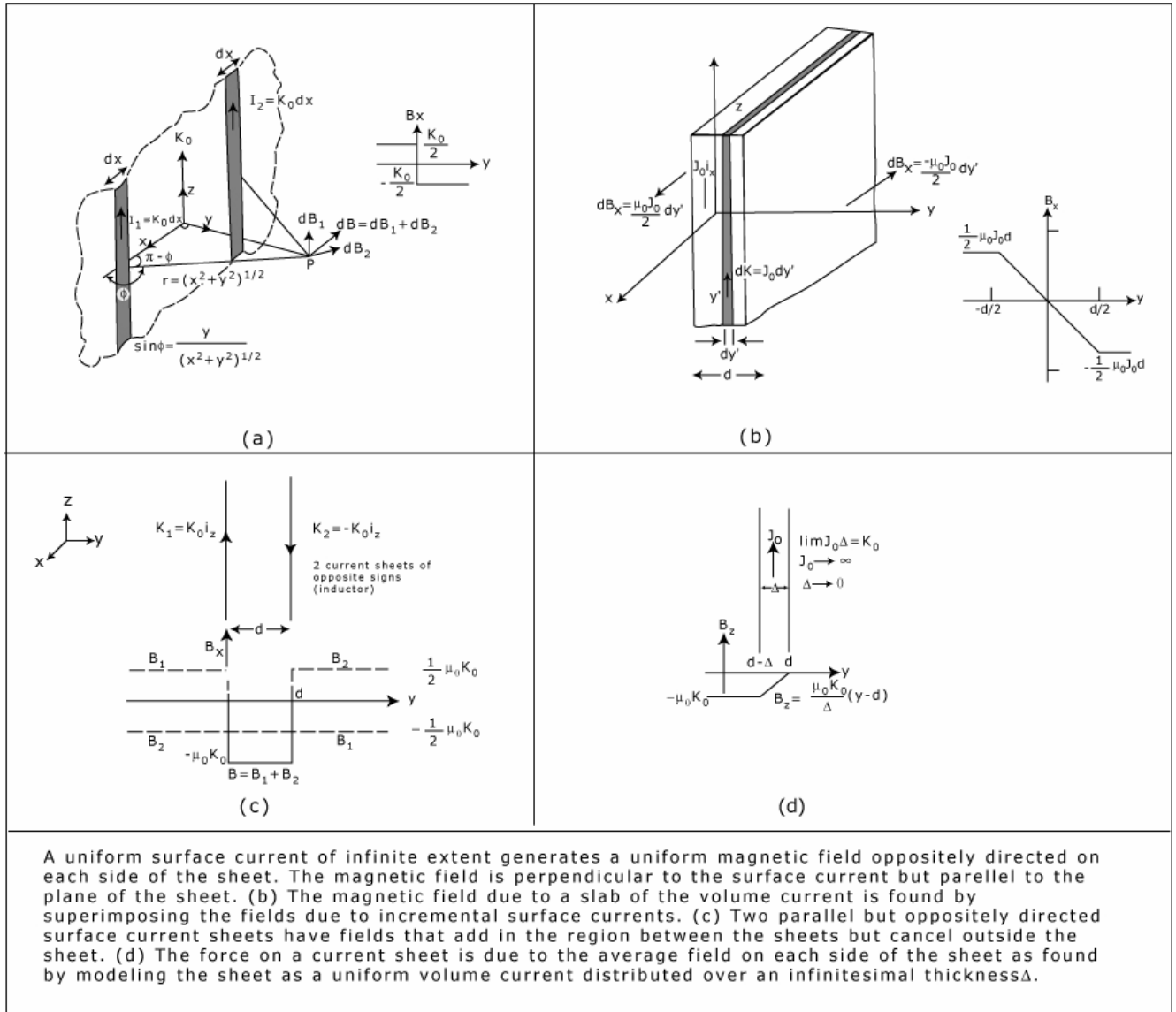
$$\frac{\sigma_0}{2\epsilon_0} - \left(-\frac{\sigma_0}{2\epsilon_0}\right) = \frac{\sigma_0}{\epsilon_0}$$

c. Two sheets of Surface Charge (Capacitor)

$$\bar{E}_1 = \begin{cases} \frac{\sigma_0}{2\epsilon_0} \bar{i}_y & y > -a \\ -\frac{\sigma_0}{2\epsilon_0} \bar{i}_y & y < -a \end{cases}, \bar{E}_2 = \begin{cases} -\frac{\sigma_0}{2\epsilon_0} \bar{i}_y & y > a \\ \frac{\sigma_0}{2\epsilon_0} \bar{i}_y & y < a \end{cases}$$

$$\bar{E} = \bar{E}_1 + \bar{E}_2 = \begin{cases} \frac{\sigma_0}{\epsilon_0} \bar{i}_y & |y| < a \\ 0 & |y| > a \end{cases}$$

## 2. Magnetic Field from a Sheet of Surface Current



From a line current  $I$

$$H_\phi = \frac{I}{2\pi r}$$

$$\bar{i}_\phi = -\sin\phi \bar{i}_x + \cos\phi \bar{i}_y$$

Thus from 2 symmetrically located line currents

$$dH_x = \frac{dl}{2\pi(x^2 + y^2)^{1/2}} (-\sin\phi)$$



$$= -\frac{K_0 dx}{2\pi} \frac{y}{x^2 + y^2}$$

$$H_x = -\frac{K_0}{2\pi} y \int_{x=-\infty}^{+\infty} \frac{dx}{x^2 + y^2}$$

$$= -\frac{K_0 y}{2\pi} \frac{1}{y} \tan^{-1} \frac{x}{y} \Big|_{x=-\infty}^{+\infty}$$

$$= \begin{cases} -\frac{K_0}{2} & y > 0 \\ +\frac{K_0}{2} & y < 0 \end{cases}$$

Check boundary condition at  $y=0$ :

$$H_x(y = 0_+) - H_x(y = 0_-) = -K_0$$

$$-\frac{K_0}{2} - \left(\frac{K_0}{2}\right) = -K_0$$