# Lecture 36 - Bipolar Junction Transistor (cont.)

May 4, 2007

# **Contents:**

- 1. Current-voltage characteristics of ideal BJT (cont.)
- 2. Charge-voltage characteristics of ideal BJT
- 3. Small-signal behavior of ideal BJT

# Reading material:

del Alamo, Ch. 11, §§11.2 (11.2.5), 11.3, 11.4 (11.4.1)

# Key questions

- How do the output characteristics of the ideal BJT look like?
- How do the charge-voltage characteristics of the ideal BJT look like?
- What is the topology of the small-signal equivalent circuit model of the ideal BJT in FAR?
- What are the key dependencies of its elements?

## 1. Current-voltage characteristics of ideal BJT (cont.)

Ideal BJT current equations (superposition of forward active + reverse):

$$\begin{split} I_C &= I_S(\exp\frac{qV_{BE}}{kT} - \exp\frac{qV_{BC}}{kT}) - \frac{I_S}{\beta_R}(\exp\frac{qV_{BC}}{kT} - 1) \\ I_B &= \frac{I_S}{\beta_F}(\exp\frac{qV_{BE}}{kT} - 1) + \frac{I_S}{\beta_R}(\exp\frac{qV_{BC}}{kT} - 1) \\ I_E &= -\frac{I_S}{\beta_F}(\exp\frac{qV_{BE}}{kT} - 1) - I_S(\exp\frac{qV_{BE}}{kT} - \exp\frac{qV_{BC}}{kT}) \end{split}$$

Equivalent circuit model representation:



Complete model has only three parameters:  $I_S$ ,  $\beta_F$ , and  $\beta_R$ .

### □ Common-emitter output I-V characteristics

vs.  $V_{CB}$ :



vs.  $V_{CE}$ :



$$V_{CEsat} = -V_{BCon} + V_{BEon}$$

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## $I_C$ vs. $V_{CB}$ with $I_B$ as parameter:



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#### Where is the reverse regime?

![](_page_5_Figure_3.jpeg)

#### Common-emitter output characteristics:

![](_page_6_Figure_3.jpeg)

#### Zoom into inverse regime:

![](_page_7_Figure_3.jpeg)

## 2. Charge-voltage characteristics of ideal BJT

In BJT, two types of stored charge:

- depletion layer charge
- minority carrier charge

In forward-active regime:

![](_page_8_Figure_7.jpeg)

#### $\Box$ Depletion layer charge

In B-E and B-C SCR's, respectively:

$$Q_{jE} = A_E \sqrt{\frac{2\epsilon q N_E N_B (\phi_{biE} - V_{BE})}{N_E + N_B}}$$
$$Q_{jC} = A_C \sqrt{\frac{2\epsilon q N_B N_C (\phi_{biC} - V_{BC})}{N_B + N_C}}$$

 $\phi_{biE}$  and  $\phi_{biC}$  are respective built-in potentials.

Since  $N_E \gg N_B \gg N_C$ ,

$$Q_{jE} \simeq A_E \sqrt{2\epsilon q N_B(\phi_{biE} - V_{BE})}$$
  
 $Q_{jC} \simeq A_C \sqrt{2\epsilon q N_C(\phi_{biC} - V_{BC})}$ 

Depletion capacitance:

$$C_{je} = \frac{\partial Q_{jE}}{\partial V_{BE}} \simeq A_E \sqrt{\frac{\epsilon q N_B}{2(\phi_{biE} - V_{BE})}} = \frac{C_{jeo}}{\sqrt{1 - \frac{V_{BE}}{\phi_{biE}}}}$$
$$C_{jc} = \frac{\partial Q_{jC}}{\partial V_{BC}} \simeq A_C \sqrt{\frac{\epsilon q N_C}{2(\phi_{biC} - V_{BC})}} = \frac{C_{jco}}{\sqrt{1 - \frac{V_{BC}}{\phi_{biC}}}}$$

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## $\Box$ Minority carrier charge

Excess minority carriers in QNR's  $\Rightarrow$  excess majority carriers to keep quasi-neutrality  $\Rightarrow$  diffusion capacitance.

Key result from pn diode: in "short" or "transparent" QNR:

stored charge = minority carrier transit time × injected minority carrier current

• For **emitter** in FAR:

![](_page_10_Figure_7.jpeg)

$$Q_E = \tau_{tE} I_B$$

with hole transit time:

$$\tau_{tE} = \frac{W_E^2}{2D_E}$$

#### • For **base** in FAR:

![](_page_11_Figure_3.jpeg)

$$Q_B = \tau_{tB} I_C$$

with electron transit time:

$$\tau_{tB} = \frac{W_B^2}{2D_B}$$

Comments:

- Units of  $Q_E$  and  $Q_B$  are C.
- $Q_E$  and  $Q_B$  scale with  $A_E$ .

Total minority carrier charge in FAR:

$$Q_F = Q_E + Q_B = \tau_{tE}I_B + \tau_{tB}I_C = (\frac{\tau_{tE}}{\beta_F} + \tau_{tB})I_C = \tau_F I_C$$

$$\tau_F \equiv intrinsic \ delay \ [s]$$

 $\tau_F$  is overall time constant for minority carrier storage in BJT in FAR:

$$\tau_F = \frac{\tau_{tE}}{\beta_F} + \tau_{tB}$$

Note: emitter contribution to  $\tau_F$  is  $\tau_{tE}/\beta_F$  because  $I_B$  is  $\beta_F$  times smaller than  $I_C$ .

If  $V_{BE}$  changes,  $Q_E$  and  $Q_B$  change  $\Rightarrow$  capacitive effect:

$$C_F = \frac{dQ_F}{dV_{BE}} = \tau_F \frac{qI_C}{kT}$$

Location of this capacitance? Think of which terminals supply stored charge (minority and majority carriers):

For  $Q_E$ :

- minority carriers (holes) injected from *base*
- $\bullet$  majority carriers (electrons) come from emitter contact

For  $Q_B$ :

- minority carriers (electrons) injected from *emitter*
- $\bullet$  majority carriers (holes) come from base contact

Equivalent-circuit model:

![](_page_13_Figure_10.jpeg)

Similar picture in reverse regime: charge storage in base and collector

$$Q_R = \tau_R I_E$$

 $\tau_R$  a bit complicated because it accounts for charge storage in *in-trinsic* and *extrinsic* base and collector regions.

![](_page_14_Figure_5.jpeg)

Diffusion capacitance:

$$C_R = \frac{dQ_R}{dV_{BC}} = \tau_R \frac{qI_E}{kT}$$

Located between base and collector terminals.

By superposition, complete equivalent circuit model valid in all four regimes:

![](_page_15_Figure_3.jpeg)

## 3. Small-signal behavior of ideal BJT

In analog (and digital) applications, interest in behavior of BJT to *small-signal* applied on top of bias

 $\Rightarrow$  small-signal equivalent circuit model.

# $\square$ Small-signal equivalent circuit model in FAR

Must *linearize* hybrid- $\pi$  model in FAR:

![](_page_16_Figure_7.jpeg)

-Non-linear voltage-controlled current source linearized to *linear voltage-controlled current source*.

-Diode linearized to *resistor*.

-Charge storage elements linearized to *capacitors*.

• Linearized voltage-controlled current source

Apply small signal  $v_{be}$  on top of bias  $V_{BE}$ .

![](_page_17_Figure_4.jpeg)

Collector current:

$$I_{C} + i_{c} = I_{S} \exp \frac{q(V_{BE} + v_{be})}{kT} \simeq I_{S} \exp \frac{qV_{BE}}{kT} (1 + \frac{qv_{be}}{kT}) = I_{C} (1 + \frac{qv_{be}}{kT})$$

Small-signal collector current:

$$i_c = \frac{qI_C}{kT} v_{be}$$

Define *transconductance*:

$$g_m = \frac{qI_C}{kT}$$

 $g_m$  depends only on absolute value of  $I_C$  and T (unlike MOSFET, where  $g_m$  depends on device geometry)

• Linearized diode

![](_page_18_Figure_3.jpeg)

Base current:

$$I_B + i_b = I_S \exp \frac{q(V_{BE} + v_{be})}{kT} \simeq \frac{I_S}{\beta_F} \exp \frac{qV_{BE}}{kT} (1 + \frac{qv_{be}}{kT})$$

Small-signal base current:

$$i_b = \frac{qI_B}{kT} v_{be}$$

Define *conductance*:

$$g_{\pi} = \frac{qI_B}{kT} = \frac{q}{kT}\frac{I_C}{\beta_F} = \frac{g_m}{\beta_F}$$

Then, in general,

$$g_{\pi} \ll g_m$$

• Capacitors

$$Q_{jE} \to C_{je}$$

$$Q_{jC} \to C_{jc}$$

$$Q_F \to C_\pi = \tau_F \frac{qI_C}{kT}$$

Two components in  $C_{\pi}$ :

![](_page_19_Figure_7.jpeg)

Note:

 $C_{\pi} = \tau_F g_m$ 

Cite as: Jesús del Alamo, course materials for 6.720J Integrated Microelectronic Devices, Spring 2007. MIT OpenCourseWare (http://ocw.mit.edu/), Massachusetts Institute of Technology. Downloaded on [DD Month YYYY]. • Small-signal equivalent circuit model for ideal BJR in FAR:

![](_page_20_Figure_3.jpeg)

## Key conclusions

- In BJT, two types of stored charge: depletion layer charge and minority carrier charge.
- Depletion layer charge accounted through depletion capacitances.
- Minority carrier charge accounted through time constant  $\tau_F$  (*in-trinsic delay*):

$$\tau_F = \frac{\tau_{tE}}{\beta_F} + \tau_{tB}$$

Emitter contribution to  $\tau_F$  is  $\beta_F$  times smaller than  $\tau_{tE}$  because  $I_B$  is  $\beta_F$  times smaller than  $I_C$ .

• Non-linear hybrid- $\pi$  model for ideal BJT including charge storage elements:

![](_page_21_Figure_9.jpeg)

• Small-signal equivalent circuit model of ideal BJT in FAR:

![](_page_22_Figure_3.jpeg)

with:

$$g_m = \frac{qI_C}{kT}$$
  $g_\pi = \frac{qI_B}{kT} = \frac{g_m}{\beta_F}$   $C_\pi = \tau_F g_m$