

## Honeycombs - In-plane behaviour

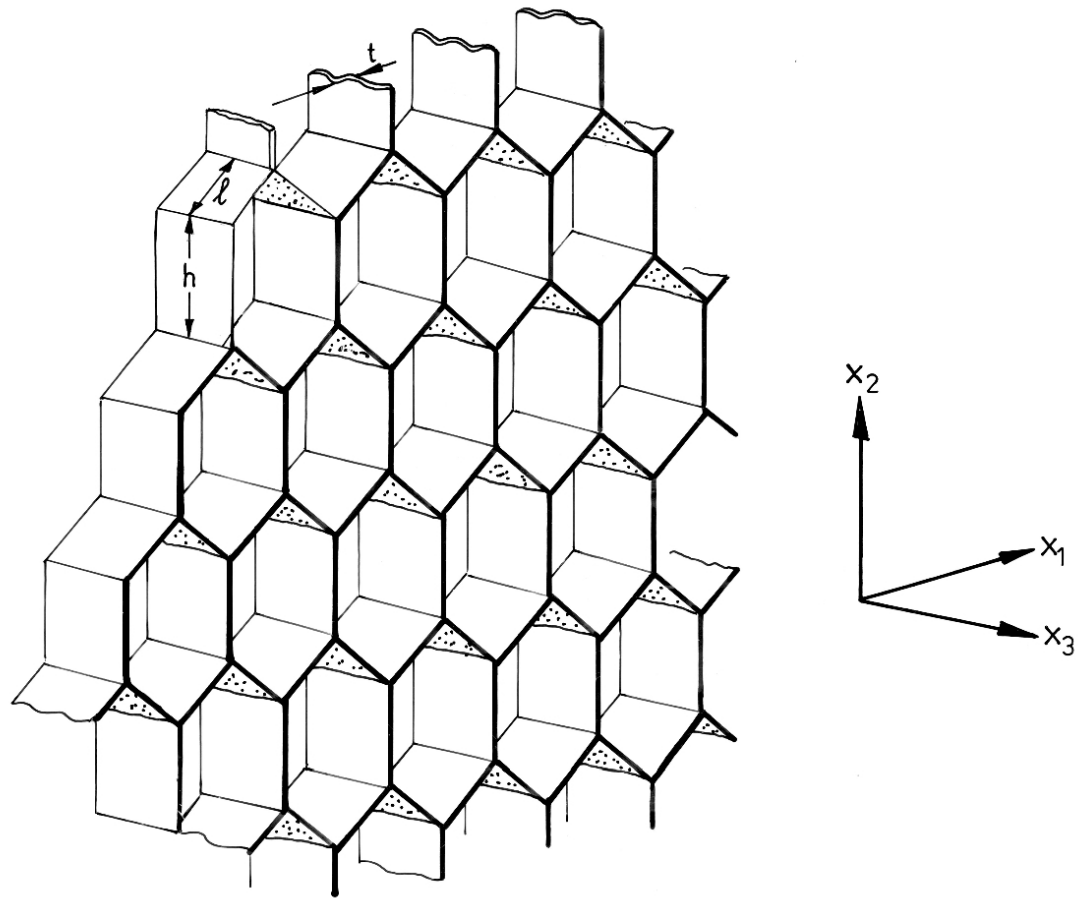
- prismatic cells
  - polymer, metal, ceramic honeycombs widely available
  - used for sandwich structure cores, energy absorption, carriers for catalysts
  - some natural materials (eg. wood, cork) can be idealized as honeycombs
  - mechanisms of deformation + failure in hexagonal honeycombs parallel those in foams
    - simpler geometry (unit cell) - easier to analyze
  - mechanisms of deformation in triangular honeycombs parallels those in 3D trusses (lattice materials)
- 

## Stress - Strain curves + Deformation behaviour : In - Plane

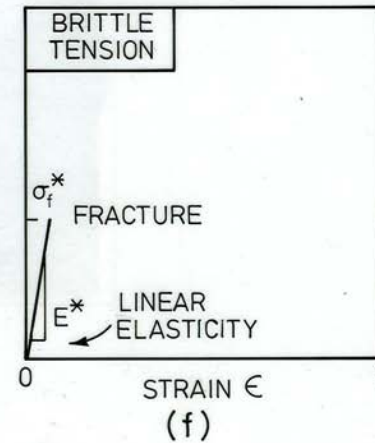
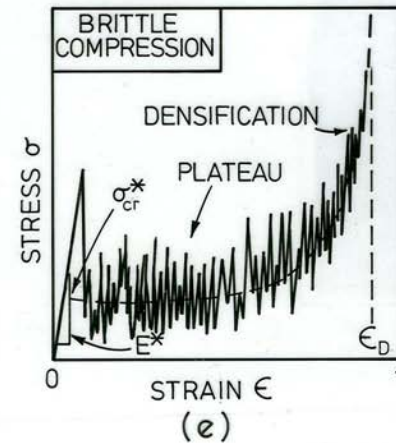
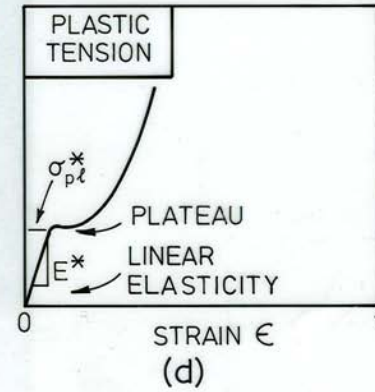
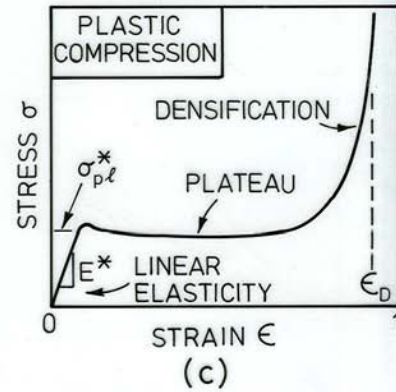
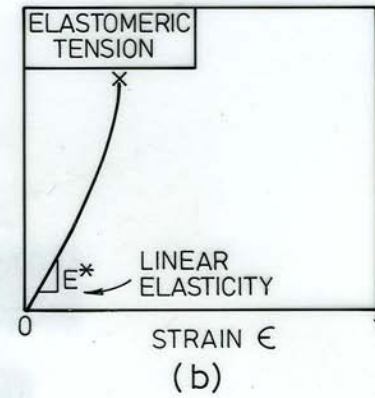
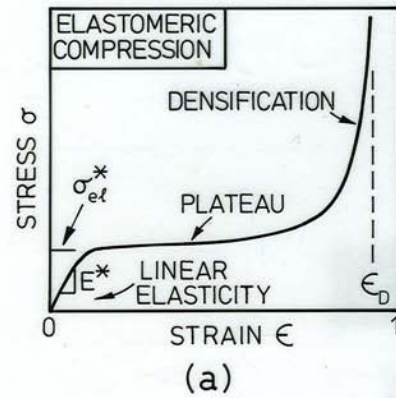
### Compression

- 3 regimes - linear elastic - bending
  - stress plateau - buckling
  - yielding
  - brittle crushing
- densification - cell walls touch
- increasing  $t/l \Rightarrow E^* \uparrow \quad \sigma^* \uparrow \quad \epsilon_D \downarrow$

# Honeycomb Geometry

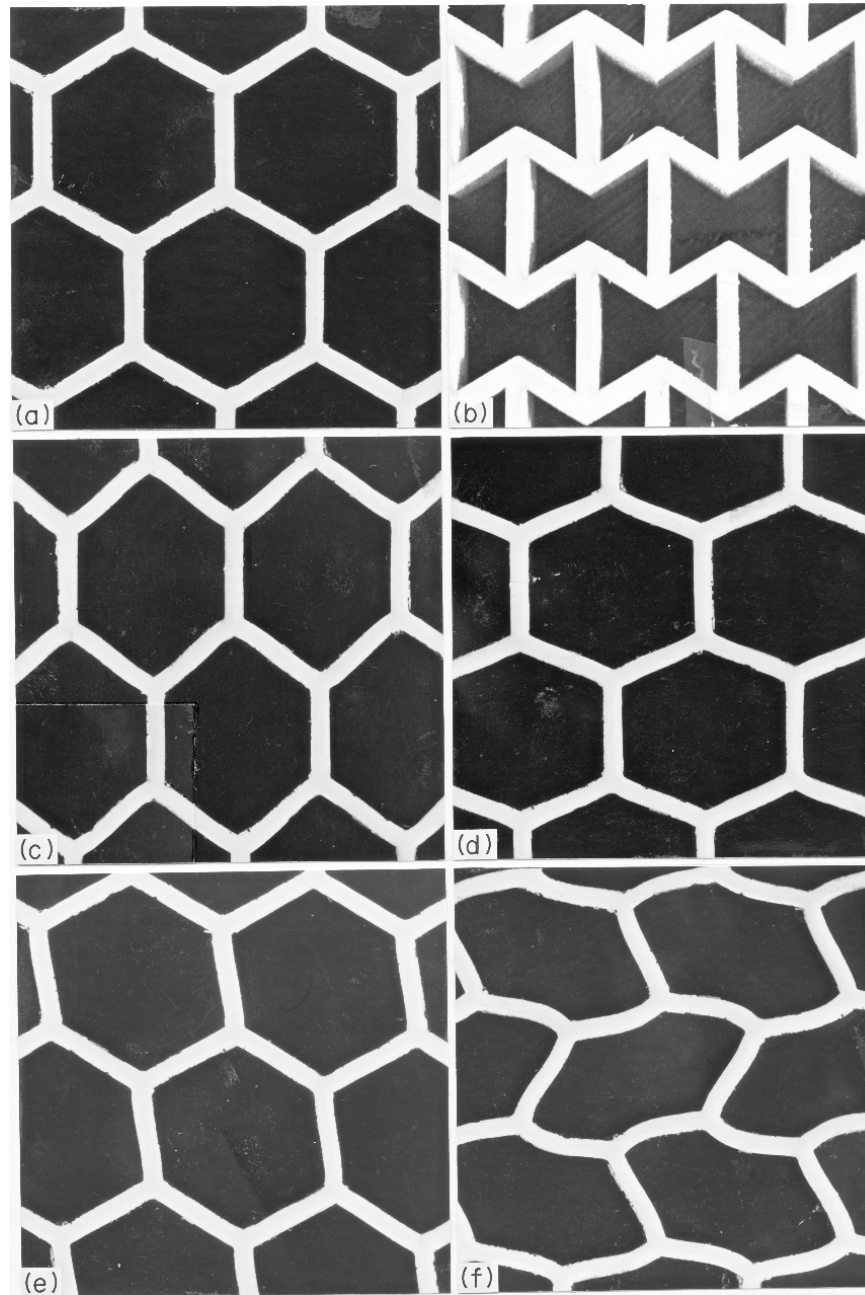


Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.



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# Deformation mechanisms

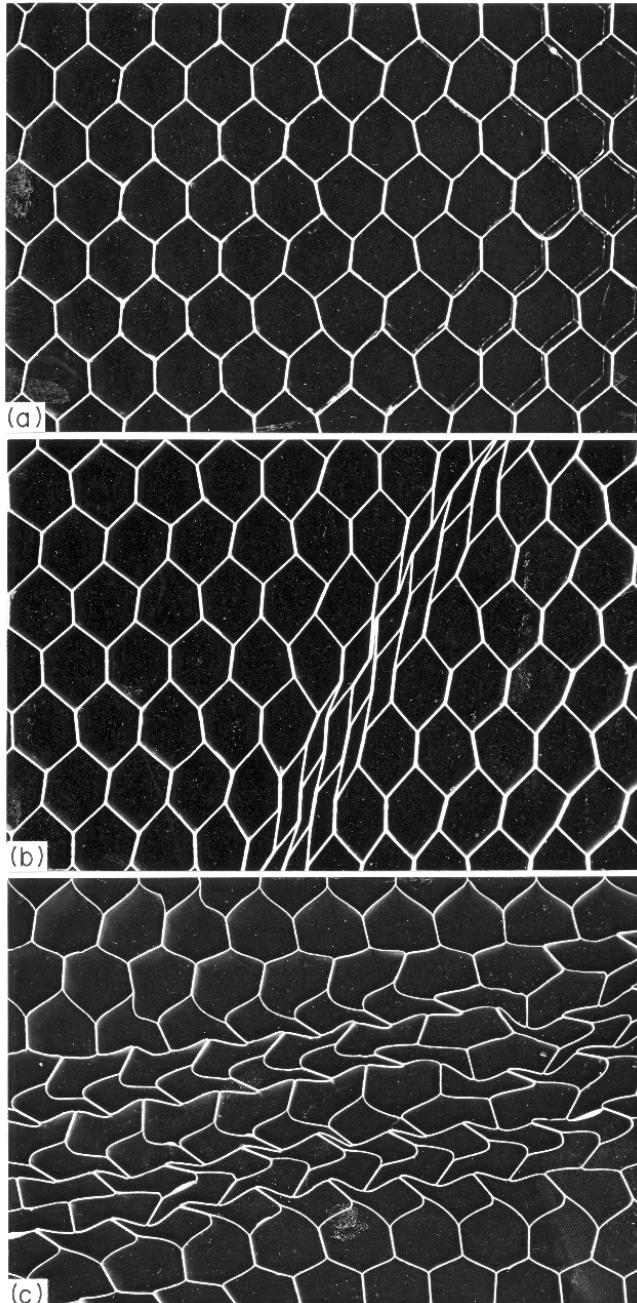


Bending  
 $X_1$  Loading

Bending  
 $X_2$  Loading

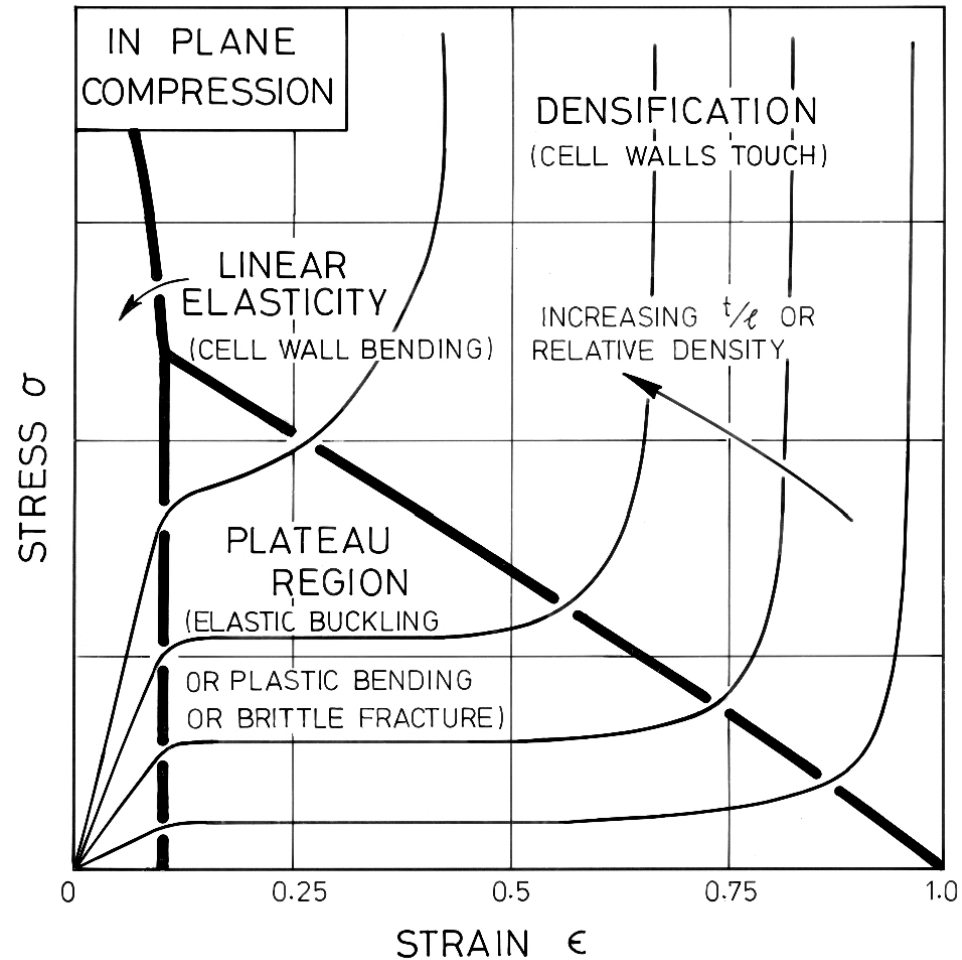
Bending  
Shear

Buckling



## Plastic collapse in an aluminum honeycomb

# Stress-Strain Curve



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

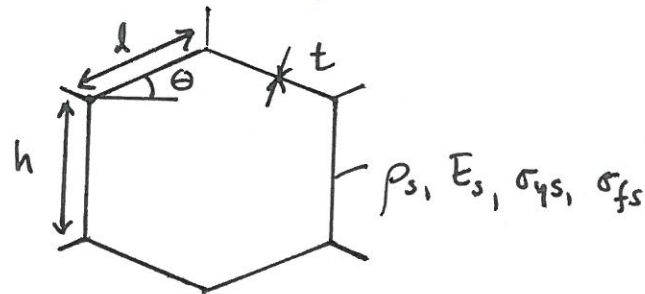


tension

- linear elastic - bending
- stress plateau - exists only if cell walls yield
  - no buckling in tension
  - brittle honeycombs fracture in tension

Variables affecting honeycomb properties

- relative density  $\frac{\rho^*}{\rho_s} = \frac{t/l (h/l + 2)}{2 \cos \theta (h/l + \sin \theta)} = \frac{2}{\sqrt{3}} \frac{t}{l}$  regular hexagons
- solid cell wall properties:  $\rho_s, E_s, \sigma_{ys}, \sigma_{fs}$
- cell geometry:  $h/l, \theta$



## In-plane properties

assumptions:

- $\rho^*/\rho_s$  small - neglect axial + shear contribution to def<sup>ion</sup>
- deformations small - neglect changes in geometry
- cell wall - linear elastic, isotropic

symmetry

- honeycombs are orthotropic - rotate 180° about each of three mutually perpendicular axes & structure is the same.

## Linear elastic deformation

$$\begin{bmatrix} \epsilon_1 \\ \epsilon_2 \\ \epsilon_3 \\ \epsilon_4 \\ \epsilon_5 \\ \epsilon_6 \end{bmatrix} = \begin{bmatrix} 1/E_1 & -\nu_{21}/E_2 & -\nu_{31}/E_3 & 0 & 0 & 0 \\ -\nu_{12}/E_1 & 1/E_2 & -\nu_{32}/E_3 & 0 & 0 & 0 \\ -\nu_{13}/E_1 & -\nu_{23}/E_2 & 1/E_3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1/G_{23} & 0 & 0 \\ 0 & 0 & 0 & 0 & 1/G_{13} & 0 \\ 0 & 0 & 0 & 0 & 0 & 1/G_{12} \end{bmatrix} \begin{bmatrix} \sigma_1 \\ \sigma_2 \\ \sigma_3 \\ \sigma_4 \\ \sigma_5 \\ \sigma_6 \end{bmatrix}$$

(Symmetric)



- Matrix notation:
 

$\epsilon_1 = \epsilon_{11}$	$\epsilon_4 = \gamma_{23}$	$\sigma_1 = \sigma_{11}$	$\sigma_4 = \sigma_{23}$
$\epsilon_2 = \epsilon_{22}$	$\epsilon_5 = \gamma_{13}$	$\sigma_2 = \sigma_{22}$	$\sigma_5 = \sigma_{13}$
$\epsilon_3 = \epsilon_{33}$	$\epsilon_6 = \gamma_{12}$	$\sigma_3 = \sigma_{33}$	$\sigma_6 = \sigma_{12}$

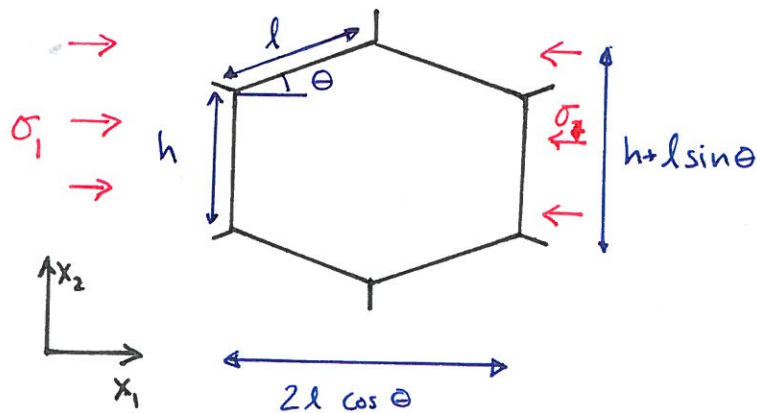
- in-plane ( $x_1-x_2$ ): 4 independent elastic constants:

$$E_1, E_2, \nu_{12}, G_{12}$$

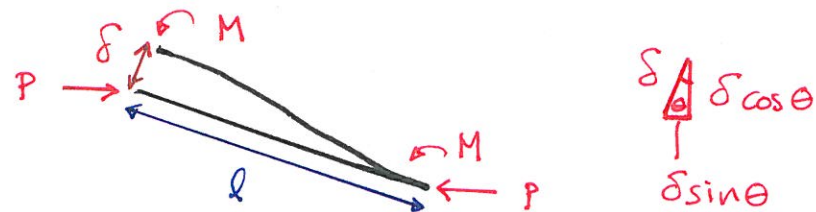
and compliance matrix symmetric -  $\frac{\nu_{12}}{E_1} = -\frac{\nu_{21}}{E_2}$  (reciprocal relation)

[ notation for Poisson's ratio:  $\nu_{ij} = -\frac{\epsilon_j}{\epsilon_i}$  ]

Young's modulus in  $x_1$  direction



unit cell in  $x_1$  direction:  $2l \cos \theta$   
 " " " " " "

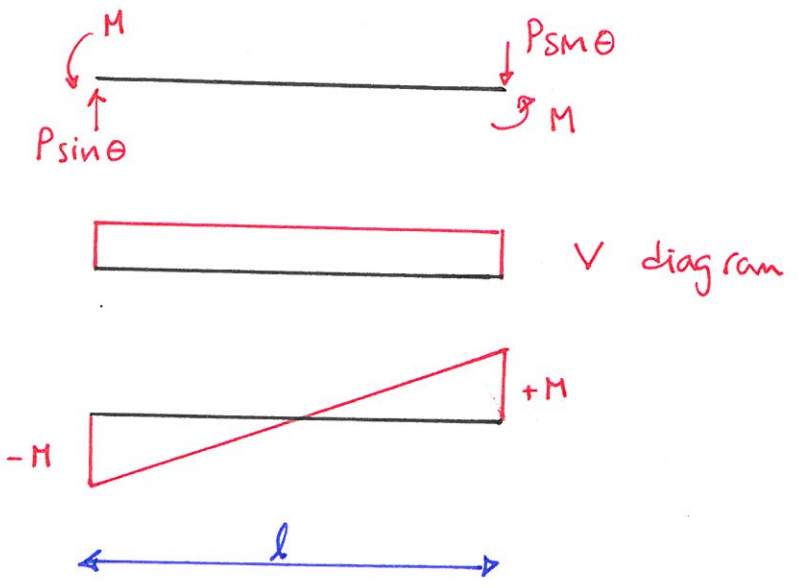


$$\sigma_1 = \frac{P}{(h + l \sin \theta) b}$$

$$\epsilon_1 = \frac{\delta \sin \theta}{l \cos \theta}$$

# In-Plane Deformation: Linear Elasticity

Figure removed due to copyright restrictions. See Figure 5: L. J. Gibson,  
M. F. Ashby, et al. "[The Mechanics of Two-Dimensional Cellular Materials.](#)"



M diagram: 2 cantilevers of length  $l/2$

$$\delta = 2 \cdot \frac{P \sin \theta (l/2)^3}{3 E_s I}$$

$$= \frac{2 P l^3 \sin \theta}{24 E_s I}$$

$$\delta = \frac{P l^3 \sin \theta}{12 E_s I}$$

$$I = \frac{b t^3}{12}$$

Combining:  $E_1^* = \frac{\sigma_1}{\epsilon_1} = \frac{P}{(h + l \sin \theta) b} \frac{l \cos \theta}{\delta \sin \theta}$

$$= \frac{P}{(h + l \sin \theta) b} \frac{l \cos \theta}{l^3 \sin^2 \theta} \cancel{1/2} E_s \cancel{1/2} \frac{t^3}{1/2}$$

$$E_1^* = E_s \left(\frac{t}{l}\right)^3 \frac{\cos \theta}{(h/l + \sin \theta) \sin^2 \theta}$$

↑ solid property    ↑ relative density    ↑ cell geometry

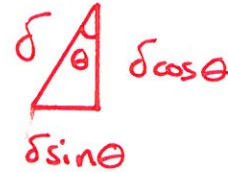
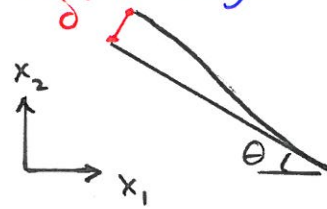
$$= \frac{4}{\sqrt{3}} \left(\frac{t}{l}\right)^3 E_s \text{ regular hexagons}$$

$h/l = 1 \quad \theta = 30^\circ$

6

Poisson's ratio for loading in  $x_1$  direction

$$\nu_{12}^* = - \frac{\epsilon_2}{\epsilon_1}$$



$$\epsilon_1 = - \frac{\delta \sin \theta}{l \cos \theta} \quad (\text{shortens})$$

$$\epsilon_2 = \frac{\delta \cos \theta}{h + l \sin \theta} \quad (\text{lengthens})$$

$$\nu_{12}^* = - \frac{\delta \cos \theta}{h + l \sin \theta} \left( - \frac{l \cos \theta}{\delta \sin \theta} \right) = \frac{\cos^2 \theta}{(h/l + \sin \theta) \sin \theta}$$

- $\nu_{12}^*$  depends only on cell geometry ( $h/l, \theta$ ), not on  $E_s, t/l$
- regular hexagonal cells:  $\nu_{12}^* = 1$
- $\nu$  can be negative for  $\theta < 0$

eg.  $h/l = 2 \quad \theta = -30^\circ \quad \nu_{12}^* = \frac{3/4}{(3/2)(-1/2)} = -1$

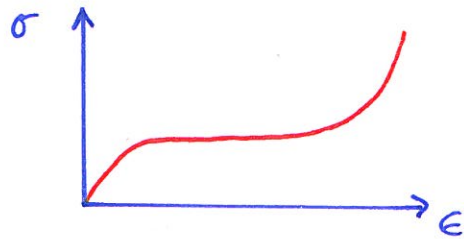
$E_2^* \quad \nu_{21}^* \quad G_{12}^*$

- can be found in similar way; results in book.

## Compressive strength (plateau stress)

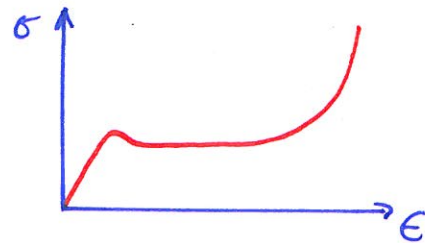
• cell collapse by:

(1) elastic buckling



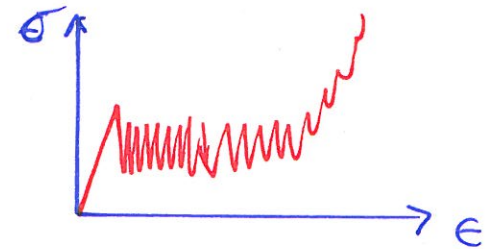
• buckling of vertical struts throughout honeycomb

(2) plastic yielding



• localization of yield  
• as def<sup>m</sup> progresses, propagation of failure band

(3) brittle crushing

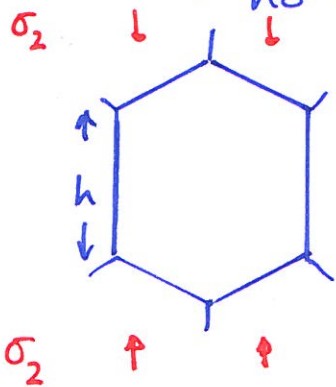


• peaks + valleys correspond to fracture of individual cell walls

## Plateau stress: elastic buckling, $\sigma_{el}^*$

• elastomeric honeycombs - cell collapse by elastic buckling of walls of length  $h$  when loaded in  $x_2$  direction

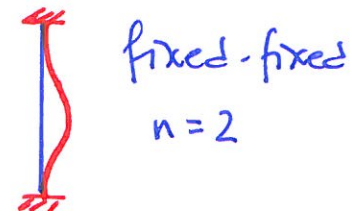
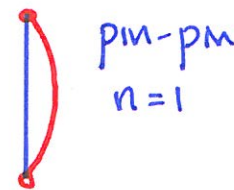
• no buckling for  $\sigma_1$ ; bending of inclined walls goes to densification.



Euler buckling load

$$P_{cr} = \frac{n^2 \pi^2 E_s I}{h^2}$$

$n$  = end constraint factor



# Elastic Buckling

Figure removed due to copyright restrictions. See Figure 7: L. J. Gibson, M. F. Ashby, et al. "[The Mechanics of Two-Dimensional Cellular Materials](#)."



- here, constraint  $n$  depends on stiffness of adjacent inclined members
- can find by elastic line analysis (see appendix if interested)
- rotational stiffness at ends of column,  $h$ , matched to rotational stiffness of inclined members
- find  $h/l = 1 \quad 1.5 \quad 2$   
 $n = 0.686 \quad 0.760 \quad 0.860$

and  $(\sigma_{el}^*)_2 = \frac{P_{cr}}{2l \cos \theta b} = \frac{n^2 \pi^2 E_s}{h^2 2l \cos \theta b} \frac{bt^3}{12}$

$$(\sigma_{el}^*)_2 = \frac{n^2 \pi^2}{24} E_s \frac{(t/l)^3}{(h/l)^2 \cos \theta}$$

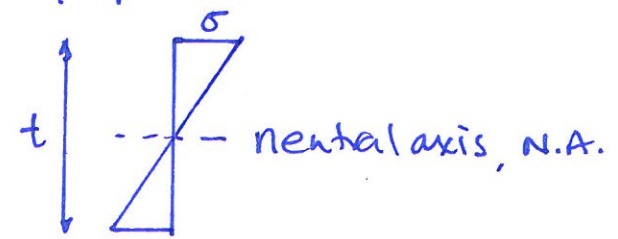
regular hexagons:  $(\sigma_{el}^*)_2 = 0.22 E_s (t/l)^3$

‡ since  $E_2^* = 4/\sqrt{3} E_s (t/l)^3 = 231 E_s (t/l)^3$

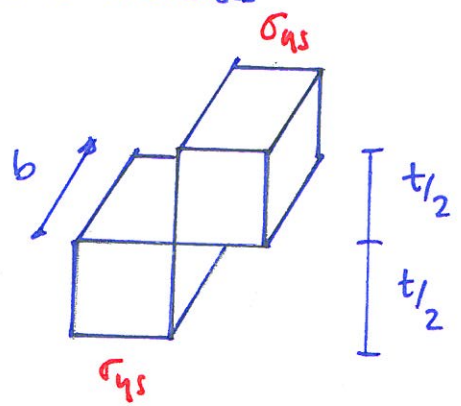
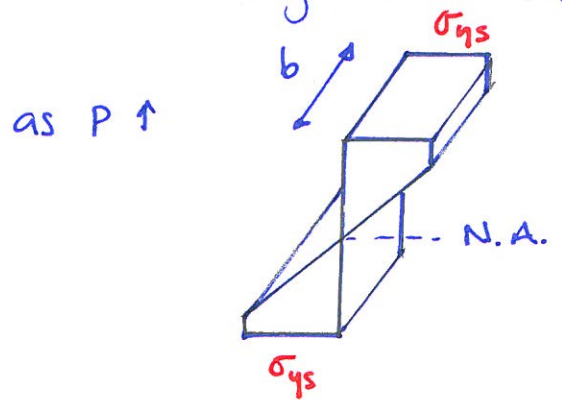
strain at buckling  $(\epsilon_{el}^*)_2 = 0.10$ , for regular hexagons independent of  $E_s, t/l$

Plateau stress: plastic yielding,  $\sigma_{pl}^*$

- failure by yielding in cell walls
- yield strength of cell walls =  $\sigma_{ys}$
- plastic hinge forms when cross-section fully yielded
- beam theory - linear elastic  $\sigma = \frac{My}{I}$



- once stress at outer fiber =  $\sigma_{ys}$ , yielding begins & then progresses through the section, as the load increases



- When section fully yielded (right fig.), form plastic "hinge"
- section rotates, like a pin

# Plastic Collapse

Figure removed due to copyright restrictions. See Figure 8: L. J. Gibson, M. F. Ashby, et al. "[The Mechanics of Two-Dimensional Cellular Materials](#)."

- moment at formation of plastic hinge (plastic moment,  $M_p$ ):

$$M_p = \left( \sigma_{ys} b \frac{t}{2} \right) \left( \frac{t}{2} \right) = \frac{\sigma_{ys} b t^2}{4}$$

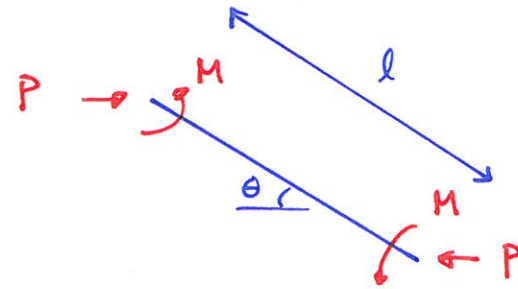
- applied moment, from applied stress

$$2 M_{app} - P l \sin \theta = 0$$

$$M_{app} = \frac{P l \sin \theta}{2}$$

$$\sigma_1 = \frac{P}{(h + l \sin \theta) b}$$

$$M_{app} = \sigma_1 (h + l \sin \theta) b \frac{l \sin \theta}{2}$$



- plastic collapse of honeycomb at  $(\sigma_{pl}^*)_1$ , when  $M_{app} = M_p$

$$(\sigma_{pl}^*)_1 (h + l \sin \theta) \frac{l \sin \theta}{2} = \sigma_{ys} \frac{b t^2}{4}$$

$$\boxed{(\sigma_{pl}^*)_1 = \sigma_{ys} \left( \frac{t}{l} \right)^2 \frac{1}{2 (h/l + \sin \theta) \sin \theta}}$$

regular hexagons:  $(\sigma_{pl}^*)_1 = \frac{2}{3} \sigma_{ys} \left( \frac{t}{l} \right)^2$

similarly,  $(\sigma_{pl}^*)_2 = \sigma_{ys} \left( \frac{t}{l} \right)^2 \frac{1}{2 \cos^2 \theta}$

• can do similar analysis for ...

- for thin-walled honeycombs, elastic buckling can precede plastic collapse (for  $\sigma_2$ )
- elastic buckling stress = plastic collapse stress  $(\sigma_{el}^*)_2 = (\sigma_{pl}^*)_2$

$$\frac{n^2 \pi^2 E_s}{24} \frac{(t/l)^3}{(h/l)^2 \cos \theta} = \frac{\sigma_{ys} (t/l)^2}{2 \cos^2 \theta}$$

$$(t/l)_{\text{critical}} = \frac{12 (h/l)^2}{n^2 \pi^2 \cos \theta} \left( \frac{\sigma_{ys}}{E_s} \right)$$

regular hexagons:  $(t/l)_{\text{critical}} = 3 \frac{\sigma_{ys}}{E_s}$

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- e.g. metals  $\sigma_{ys}/E_s \sim .002$   $(t/l)_{\text{crit}} \sim 0.6\%$ 
  - most metal honeycomb denser than this
- polymers  $\sigma_{ys}/E_s \sim 3-5\%$   $(t/l)_{\text{crit}} \sim 10-15\%$ 
  - low density polymers with yield point may buckle before yield.



Plateau stress: brittle crushing,  $(\sigma_c^*)$ ,

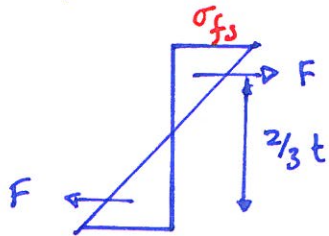
- ceramic honeycombs - fail in brittle manner
- cell wall bending - stress reaches modulus of rupture - wall fracture

loading in  $x_1$  direction:

$$P = \sigma_1 (h + l \sin \theta) b \quad \sigma_{fs} = \text{modulus of rupture of cell wall}$$

$$M_{\text{max applied}} = \frac{Pl \sin \theta}{2} = \frac{\sigma_1 (h + l \sin \theta) b l \sin \theta}{2}$$

Moment at fracture,  $M_f$



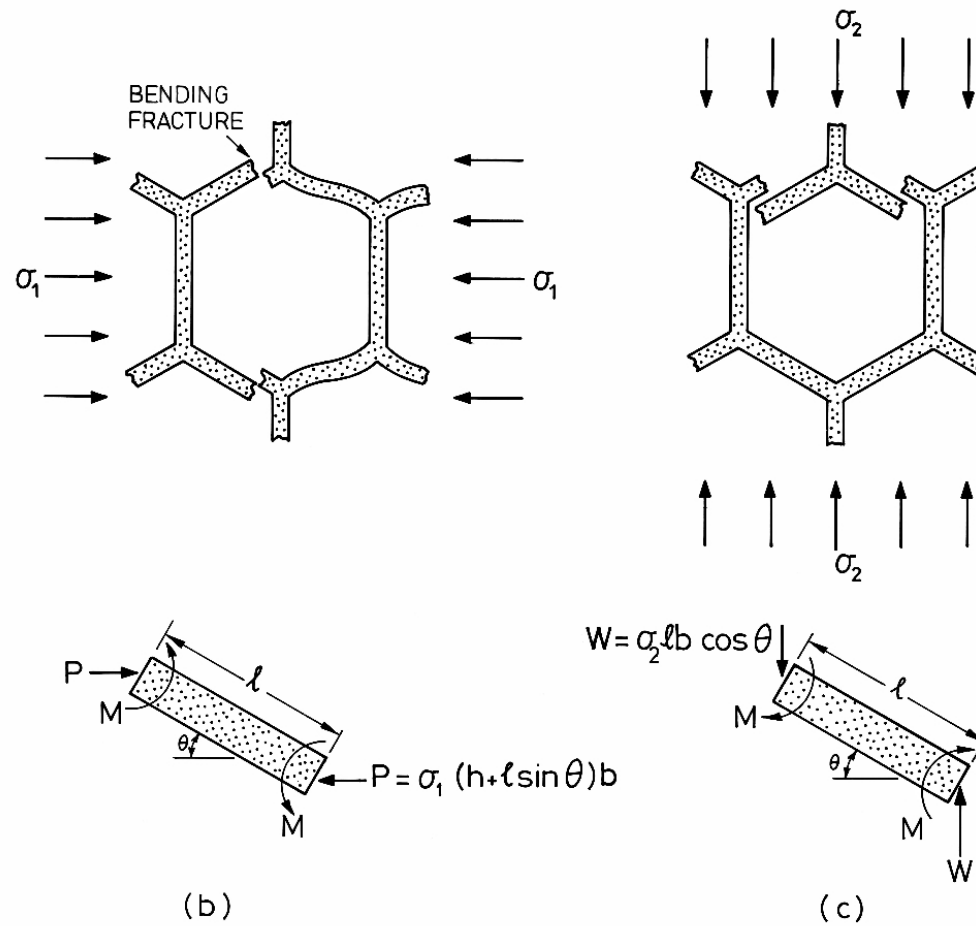
$$M_f = \left( \frac{1}{2} \sigma_{fs} b \frac{t}{2} \right) \left( \frac{2}{3} t \right) = \frac{\sigma_{fs} b t^2}{6}$$

$$\boxed{(\sigma_c^*)_1 = \sigma_{fs} \left( \frac{t}{l} \right)^2 \frac{1}{3 (h/l + \sin \theta) \sin \theta}}$$

regular hexagons:  $(\sigma_c^*)_1 = \frac{4}{9} \sigma_{fs} \left( \frac{t}{l} \right)^2$



# Brittle Crushing



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

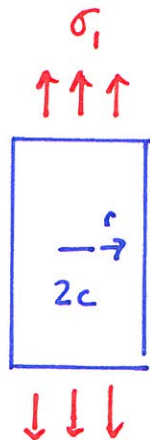
## Tension

- no elastic buckling
- plastic plateau stress approx same in tension + compression (small geometric difference due to deformation)
- brittle honeycombs: fast fracture

## Fracture toughness

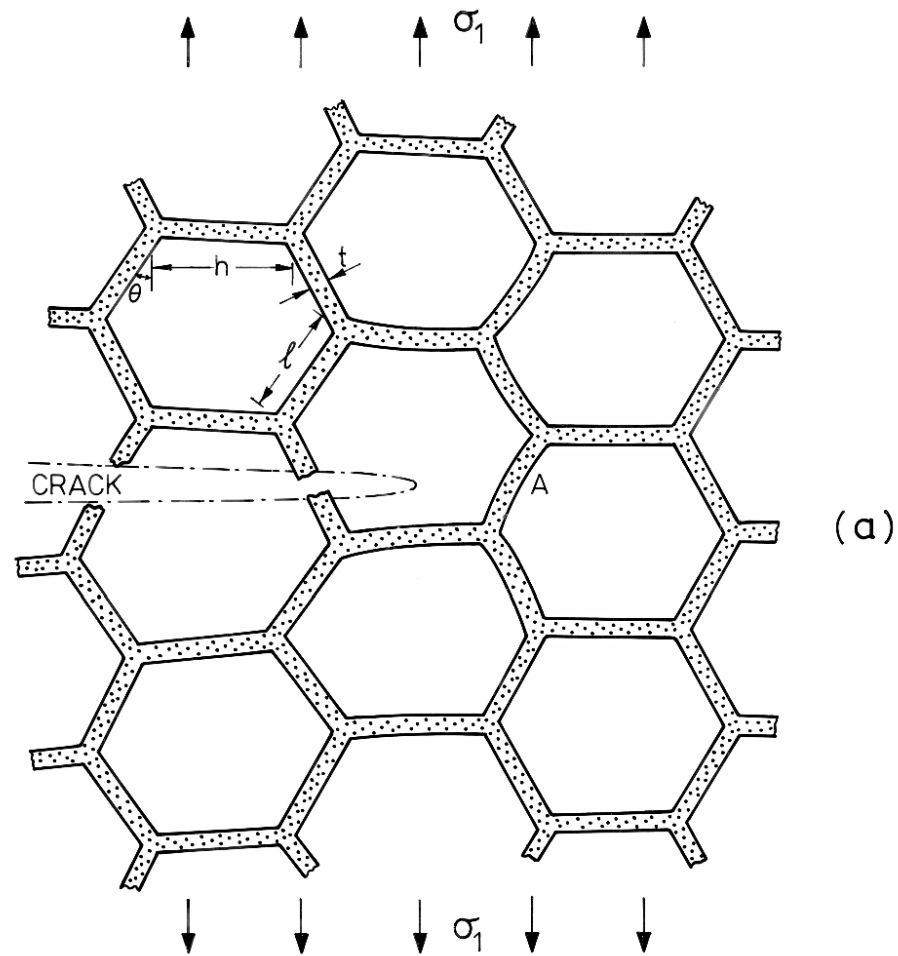
- assume:
- crack length large relative to cell size (continuum assumption)
  - axial forces can be neglected
  - cell wall material has constant modulus of rupture,  $\sigma_{fs}$

continuum: crack of length  $2c$  in a linear elastic solid material normal to a remote tensile stress  $\sigma_1$  creates a local stress field at the crack tip



$$\sigma_{\text{local}} = \sigma_{\text{I}} = \frac{\sigma_1 \sqrt{\pi c}}{\sqrt{2\pi r}}$$

# Fracture Toughness



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.

honey comb: cell walls bent - fail when applied moment = fracture moment

$$M_{app} \propto Pl \quad \text{on wall A} \quad M_f \propto \sigma_{fs} bt^2$$

$$M_{app} \propto Pl \propto \sigma_l l^2 b \propto \frac{\sigma_l \sqrt{c} l^2 b}{\sqrt{l}} \propto \sigma_{fs} bt^2$$

$$(\sigma_f^*)_1 \propto \sigma_{fs} \left(\frac{t}{l}\right)^2 \sqrt{\frac{l}{c}}$$

$$K_{Ic}^* = \sigma_f^* \sqrt{\pi c} = C \sigma_{fs} \left(\frac{t}{l}\right)^2 \sqrt{l}$$

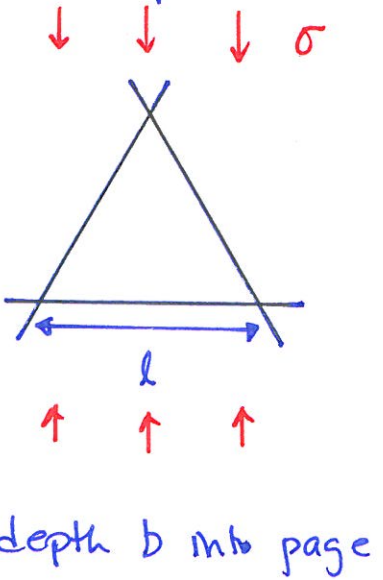
C = constant

depends on cell size,  $l$ !

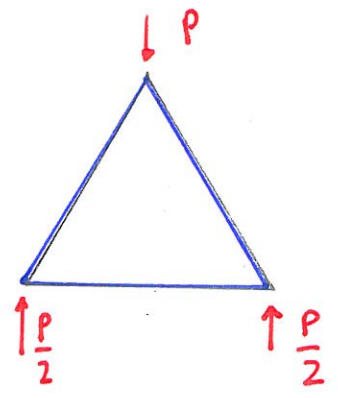
Summary: hexagonal honeycombs, in-plane properties

- linear elastic moduli:  $E_1^*$   $E_2^*$   $\nu_{12}^*$   $G_{12}^*$
- plateau stresses  $(\sigma_{el}^*)_2$  elastic buckling  
(compression)
- $\sigma_{pl}^*$  plastic yield
- $\sigma_{cr}^*$  brittle crushing
- fracture toughness  $K_{Ic}^*$  brittle fracture  
(tension)

Honeycombs: In-plane behaviour - triangular cells



- triangulated structures - trusses
- can analyze as pm-jointed (no moment @ joints)
- forces in members all axial (no bending)
- if joints fixed + include bending, difference ~ 2%
- force in each member proportional to P



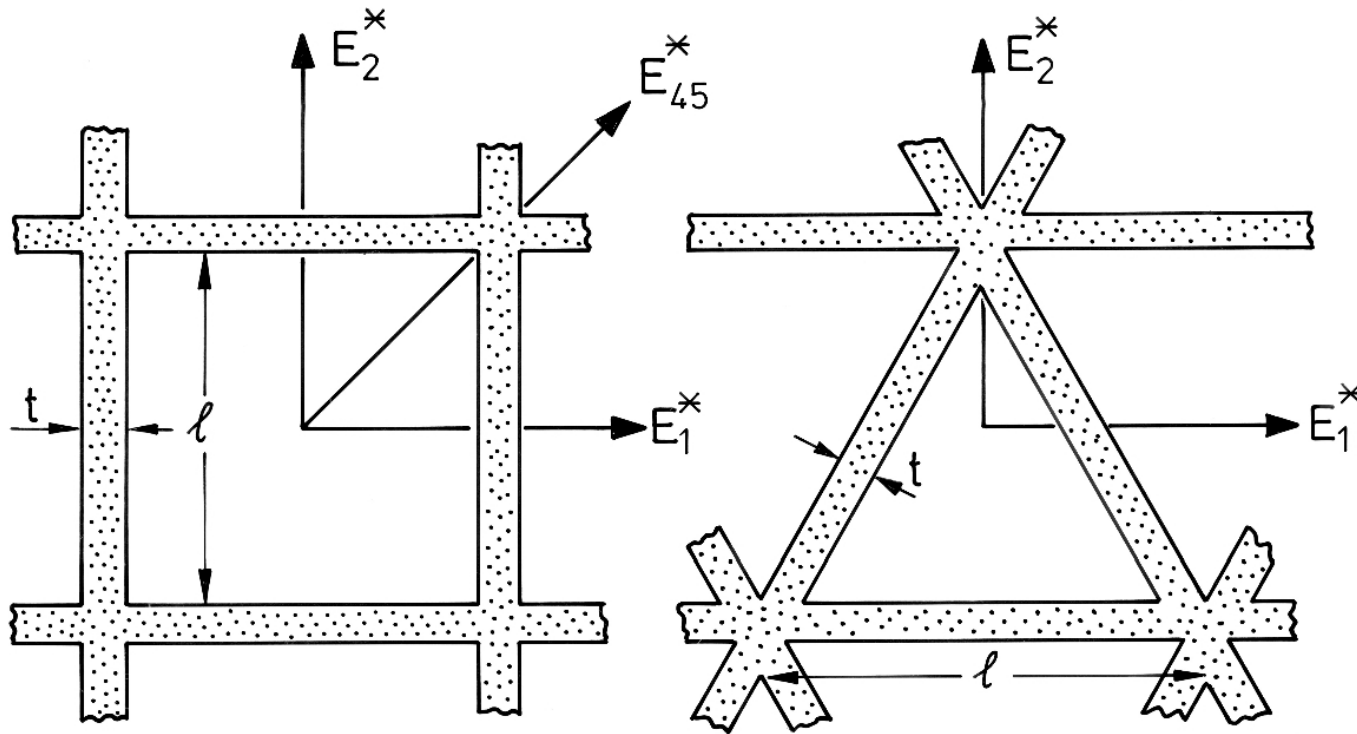
$$\sigma \propto \frac{P}{lb} \quad \epsilon \propto \frac{\delta}{l} \quad \delta \propto \frac{Pl}{AE_s} \quad (\text{axial shortening: Hooke's law})$$

$$E^* \propto \frac{\sigma}{\epsilon} \propto \frac{P}{lb} \frac{l}{\delta} \propto \frac{P}{b} \frac{btE_s}{Pl} \propto E_s \left( \frac{t}{l} \right)$$

$$E^* = c E_s \left( \frac{t}{l} \right)$$

exact calculation:  $E^* = 1.15 E_s \left( \frac{t}{l} \right)$  for equilateral triangle

# Square and Triangular Honeycombs



Gibson, L. J., and M. F. Ashby. *Cellular Solids: Structure and Properties*. 2nd ed. Cambridge University Press, 1997. Figure courtesy of Lorna Gibson and Cambridge University Press.



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