#### Remaining Issues

- Electron wave picture has fixed some thermal/electrical properties and electron velocity issues
- Still can not explain:
  - Hall coefficients
  - Colors of metals
  - Insulators, Semiconductors
- Can not ignore the ions (i.e. everything else but the valence electrons that we have been dealing with so far) any longer!
- Whatever we modify, can not change the electron wave picture that is now working well for many materials properties
- HOW DO THE VALENCE ELECTRON WAVES INTERACT WITH THE IONS AND THEIR POTENTIALS?

#### Improvements? What are ion cores doing...

- Scattering idea seems to work
- any effect of crystal (periodic) lattice?
- Diffraction
  - proves periodicity of lattice
  - proves electrons are waves
  - proves strong interaction between crystal and electrons (leads to band structures=semiconductors and insulators)
  - useful characterization technique
- Course: bias toward crystalline materials: many applications: materials related to either end of spectrum (atomic/molecular or crystalline)

localized -		extended
Point defects, atoms, molecules	Polymers, α Si	Bands; properties of many solids with or without extended defects

Diffraction is a useful characterization in all these materials

# Electrons in a Periodic Potential

- Rigorous path:  $H\Psi = E\Psi$
- We already know effect: DeBroglie and electron diffraction
- Unit cells in crystal lattice are 10<sup>-8</sup> cm in size
- Electron waves in solid are  $\lambda = h/p \sim 10^{-8}$  cm in size
- Certain wavelengths of valence electrons will diffract!

#### Diffraction Picture of the Origin of Band Gaps

Start with 1-D crystal again

λ~a Take lowest order, n=1, and consider an incident valence 1-D electron moving to the right a  $k_i = \frac{\pi}{a}; \psi_i = e^{i\frac{\pi}{a}x}$ d=a,  $n\lambda = 2d\sin\theta$  $\sin\theta = 1$ Reflected wave to left:  $k_o = -\frac{\pi}{a}; \psi_o = e^{-i\frac{\pi}{a}x}$  $n\lambda = 2a$  $\Delta k = k_i - k_o = \frac{2\pi}{2\pi}$  $k = \frac{2\pi}{2\pi}$ λ Total wave for electrons with diffracted wavelengths:  $\psi = \psi_i \pm \psi_o$  $k = \frac{\pi n}{k}$  $\psi_s = \psi_i + \psi_o = 2\cos\frac{\pi}{a}x$ a  $\psi_a = \psi_i - \psi_o = i2\sin\frac{\pi}{a}x$ 4

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Diffraction Picture of the Origin of Band Gaps Probability Density=probability/volume of finding electron= $|\psi|^2$ 

 $|\psi_a|^2 = 4\sin^2\frac{\pi}{a}x$   $|\psi_s|^2 = 4\cos^2\frac{\pi}{a}x$ a

•Only two solutions for a diffracted wave

•Electron density on atoms

- •Electron density off atoms
- •No other solutions possible at this wavelength: no free traveling wave

#### Nearly-Free Electron Model

- Assume electrons with wave vectors (k's) far from diffraction condition are still free and look like traveling waves and see ion potential, U, as a weak background potential
- Electrons near diffraction condition have only two possible solutions
  - electron densities between ions,  $E=E_{free}$ -U
  - electron densities on ions,  $E = E_{free} + U$
- Exact solution using  $H\Psi = E\Psi$  shows that E near diffraction conditions is also parabolic in k,  $E \sim k^2$

#### Nearly-Free Electron Model (still 1-D crystal)

• states



#### Consequences of Diffraction on E vs. k curves

• At  $k=\pi/a$ , there must be also a  $k=-\pi/a$  wave, since there is absolute diffraction at this k



#### Extended-Zone Scheme

- Bands form, separated by band gaps
- Note redundancy: no need for defining k outside  $+-\pi/a$  region



#### **Reduced-Zone Scheme**

• Only show  $k=+-\pi/a$  since all solutions represented there



# **Real Band Structures**

- GaAs: Very close to what we have derived in the nearly free electron model
- Conduction band minimum at k=0: Direct Band Gap



Figure by MIT OpenCourseWare.

# Real Band Structures

- Ge: Very close to GaAs, except conduction band minimum is in <111> direction, not at k=0
- Indirect Band Gap



Figure by MIT OpenCourseWare.

### Trends in III-V and II-VI Compounds



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Figure by MIT OpenCourseWare.

#### Lattice Constant (A)